Upper Limit to the $A_2 \rightarrow \eta' \pi$ Decay Mode*

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An upper limit to the branching ratio $(A_2 \rightarrow \eta' \pi)/(A_2 \rightarrow \rho \pi)$ is determined at the 90% confidence level to be 0.011.

In a previous publication¹ we reported a measurement of the branching ratio

$$R' = (A_2 \rightarrow \eta' \pi) / (A_2 \rightarrow \rho \pi) . \tag{1}$$

The result, $R' = 0.07 \pm 0.03$, was obtained by observing the $\rho^{0}\pi^{-}$ decay in the channel

$$\pi^- p \to p \pi^+ \pi^- \pi^- \tag{2}$$

and by searching for the $\eta' \pi^-$ decay mode in the channel

$$\pi^{-}p \rightarrow p\pi^{-}\pi^{+}\pi^{-}(\eta^{0})$$
neutrals. (3)

In channel (3) the separation of $p\pi^-\eta'$ events from the background is limited by the measurement errors on the η^0 (missing neutrals) and on the η' $(\eta^0\pi^+\pi^-)$ masses. These errors vary with the " $\eta'\pi^{-n}$ " mass, resulting in a rapid increase of the background with this mass. The result for R' given in Ref. 1 was obtained, in a somewhat indirect way, by subtracting from the numbers of events near the η^0 , η' masses, the numbers of events in a surrounding control region. Other published determinations²⁻⁶ of the $\eta'\pi/\rho\pi$ branching ratio are shown in Table I. These results are also based, almost exclusively, on a study of channel (3). It is clear that the value of R' is not well determined from these studies.

In this note we report a new study, based on a larger data sample, in which we searched for the $A_2 \rightarrow \eta' \pi$ decay mode in $p \pi \neg \eta'$ events with the η^0 decaying by a charged mode. These events may contribute to the channel⁷

$$\pi^- p \to p \pi^- \pi^+ \pi^- \pi^+ \pi^- \pi^0 \text{ (or } \gamma). \tag{4}$$

In channel (4) the $p\pi^{-}\eta'$ state should be (and is) clearly separated from background. The data for reaction (4) were obtained in a 5-GeV/c π^{-} ex-

posure of the 82-in. hydrogen bubble chamber at SLAC. The sample of events analyzed for reaction (4) corresponds to $0.086 \pm 0.007 \ \mu b/event$.

By kinematic fitting and by visual observation of track ionization, 3379 events $(290 \pm 25 \ \mu b)$ were assigned to channel (4). Assuming the branching ratios⁸

$$r_1 = \frac{\eta' - \eta^0 \pi^+ \pi^-}{\eta' - \text{all}} = 0.454 \pm 0.015$$

and

$$r_2 = \frac{\eta - \pi^+ \pi^- \pi^0 \text{ (or } \gamma)}{\eta - \text{all}} = 0.289 \pm 0.006^7$$

the observation of one $p\pi^-\eta'$ event in channel (4) would correspond to a cross section of $0.086/r_1r_2$ = $0.65 \pm 0.06 \ \mu$ b. We observe $28 \pm 8 \ p\pi^-\eta'$ events (after subtraction of background under both η' and η) in channel (4) for a cross section of 18.3 ± 5.4 μ b.⁹ In the same data the cross section for π^-p $-pA_2^- \rightarrow p\rho^0\pi^- \rightarrow p\pi^+\pi^-\pi^-$ in the region $M(3\pi) = 1.2 - 1.4$ GeV, $t' = \Delta^2_{pp} - \Delta^2_{min} < 0.8 \text{ GeV}^2$ has been found to be $80 \pm 5 \ \mu$ b.¹⁰ The observation of one $p\pi^-\eta'$ event in channel (4) in the same mass and t' region would thus correspond to

$$R' = \frac{(0.65 \pm 0.06)}{2 \times (80 \pm 5)} = 0.004 \pm 0.0005$$

We have found *no* event with $M(6\pi) = 1.2-1.4$ GeV and t' < 0.8 GeV², without making any cuts on either the η^0 or the η' masses. Taking into account the error on the expected number of events and Poisson statistics, our result is therefore:

 $R' < (0.004 + 1.64 \times 0.0005) \times \ln 10 = 0.011$

with 90% confidence.

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Incident beam	Experiment	$\eta \operatorname{decay}^{a}$ observed	R′ ^b	Type of upper limit
 3.2 GeV/ $c \pi^{-1}$	Ref. 2	η _n	80,0±0,08	• • •
5.0 GeV/c π^{-}	Ref. 1	η_{m}	0.07 ± 0.03	•••
5.0 GeV/c π^+	Ref. 3	$\eta_n + \eta_c$	$0.04^{+0.03}_{-0.04}$	• • • •
8.0 GeV/ $c \pi^{-}$	Ref. 4	η_n	0.15 ± 0.09	• • •
7.1 GeV/ $c \pi^{+}$	Ref. 5	η_n	<0.04	not stated
3.7 GeV/ $c \pi^+$	Ref. 6	n,	<0.025	2 std. dev.
4.9 GeV/c π^{-}	This work	η_{r}	<0.011	90% confidence

TABLE I. The branching ratio $R' = (A_2 \rightarrow \eta' \pi)/(A_2 \rightarrow \rho \pi)$.

^a η_n means $\eta \rightarrow$ neutrals, η_c means $\eta \rightarrow \pi^+ \pi^- \pi^0$ (or γ).

^b When possible, the published results have been updated by using the appropriate branching fractions for η' and η from Ref. 8.

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²S. U. Chung *et al.*, Phys. Rev. Letters <u>15</u>, 325 (1965).

³K. Böckmann et al., Nucl. Phys. <u>B16</u>, 221 (1970).

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⁵M. Alston-Garnjost *et al.*, Phys. Letters <u>34B</u>, 156 (1971).

⁶K. W. J. Barnham *et al.*, Phys. Rev. Letters <u>26</u>, 1494 (1971).

⁷We have not fitted events to the hypothesis $\pi^- p$ $\rightarrow p \pi^- \pi^+ \pi^- \pi^+ \pi^- \gamma$. The majority of such events are expected to be fitted also by the hypothesis $\pi^- p$ $\rightarrow p \pi^- \pi^+ \pi^- \pi^+ \pi^- \pi^0$.

⁸Particle Data Group, Phys. Letters <u>39B</u>, 1 (1970). ⁹In this study the η mass region is taken to be 520– 580 MeV and the η' mass region is taken to be 927.5– 987.5 MeV.

 10 G. Ascoli *et al*. (to be published).

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Effects of a Neutral Intermediate Boson on Leptonic Atoms

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The effect of a neutral vector boson on the hyperfine interval of the μ^+ - e^- bound system is calculated using Weinberg's SU(2) × U(1) lepton model. A lower limit on the mixing parameter θ is obtained, and some consequences of the additional neutral bosons predicted by the SU(3) × SU(3) version of the model are examined.

Recently, there has been renewed interest in a unified model of leptonic weak and electromagnetic interactions due to Weinberg.¹ This renewal of interest was stimulated by the investigations of 't Hooft and B. W. Lee² which indicate that models of this type, where the masses of the mediating weak bosons are generated by spontaneous breakdown of a gauge symmetry,³ may be renormalizable in the conventional sense. Furthermore, a number of higher-order perturbation calculations have since been performed and shown to be finite.⁴ These results all suggest that the model has considerable theoretical interest.

The experimental predictions embodied in Weinberg's model are equally interesting inasmuch as they differ significantly from the predictions of the conventional (V-A) theory of leptonic weak interactions. While the available weak-interac-