Production of π^+, π^-, K^+, K^-, p, and $\bar{p}$ in light (uds), c, and b jets from Z^0 decays

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We present improved measurements of the differential production rates of stable charged particles in hadronic Z^0 decays, and of charged pions, kaons, and protons identified over a wide momentum range using the SLD Cherenkov ring imaging detector. In addition to flavor-inclusive Z^0 decays, measurements are made for Z^0 decays into light (u, d, s), c, and b primary flavors, selected using the upgraded vertex detector. Large differences between the flavors are observed that are qualitatively consistent with expectations based upon previously measured production and decay properties of heavy hadrons. These results are used to test the predictions of QCD in the modified leading logarithm approximation, with the ansatz of local parton-hadron duality, and the predictions of three models of the hadronization process. The light-flavor results provide improved tests of these predictions; the heavy-flavor results provide complementary model tests. In addition we have compared hadron and antihadron production in light quark (as opposed to antiquark) jets. Differences are observed at high momentum for all three charged hadron species, providing direct probes of leading particle effects, and stringent constraints on models.

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I. INTRODUCTION

The production of jets of hadrons from hard partons produced in high energy collisions is understood qualitatively in terms of three stages. Considering the process $e^+e^- \to q\bar{q}$, the first stage, ''fragmentation,'' involves the radiation of gluons from the primary quark and antiquark, which in turn may radiate gluons or split into $q\bar{q}$ pairs until their virtuality approaches the hadron mass scale. This process is in principle calculable in perturbative QCD, and three approaches have been taken so far: (i) differential cross sections have been calculated [1] for the production of up to 4 partons to second order in the strong coupling α_s, and leading order calculations have been performed recently for as many as 6 partons (see e.g. [2]); (ii) certain parton distributions have been calculated to all orders in α_s in the modified leading logarithm approximation (MLLA) [3]; (iii) ''parton shower'' calculations [4] have been implemented in Monte Carlo simulations that consist of an arbitrary number of $q \to qg$, $g \to gg$ and $g \to q\bar{q}$ branchings, with each branching probability determined from QCD at up to next-to-leading logarithmic order.

In the second stage these partons transform into ''primary'' hadrons. This ''hadronization'' process is not understood quantitatively and there are few theoretical predictions for light hadrons, i.e., those that do not contain a c or b quark. One approach is to make the ansatz of local parton-hadron duality (LPHD) [3] that inclusive distributions of primary hadrons are the same, up to a normalization factor, as those for partons. Calculations using MLLA QCD, cut off at a virtual parton mass comparable with the mass of the hadron in question, have been used in combination with LPHD to predict properties of the distribution of $\xi = -\ln(2p/E_{CM})$, where p and E_{CM} are the hadron momentum and the e^+e^- energy, respectively, in the e^+e^- center of mass frame. For a given primary hadron species, the shape should be described well by a Gaussian function within about one unit of the peak position; a wider ξ range should be described with the addition of small distortion terms; and the peak position should depend inversely on the hadron mass and logarithmically on E_{CM}. At sufficiently high $x_p = 2p/E_{CM}$, perturbative QCD has also been used to calculate the E_{CM} dependence of the x_p distributions [5]. It is desirable to test these predictions experimentally and to encourage deeper theoretical understanding of the hadronization process.

In the third stage unstable primary hadrons decay into the stable particles that traverse particle detectors. This stage is understood inasmuch as proper lifetimes and decay branching ratios have been measured for many hadron species. However, these decays complicate fundamental hadronization measurements because many of the stable particles are decay products rather than primary hadrons, and it is typically not possible to determine whether a given detected hadron is primary. Previous measurements at e^+e^- colliders (see e.g. [6,7]) indicate that decays of vector mesons, strange baryons and decuplet baryons produce roughly two-thirds of the stable particles; scalar and tensor mesons, and radially excited baryons have also been observed [7], which contribute additional secondaries. Ideally one would measure every possible hadron species and distinguish primary hadrons from decay products on a statistical basis. A body of knowledge could be assembled by reconstructing heavier and heavier states, and subtracting their known decay products from the measured differential production rates of lighter hadrons. The measurements presented here constitute a first step.

Additional complications arise in jets initiated by heavy quarks, since the leading heavy hadrons carry a large fraction of the beam energy, restricting that available to other primary hadrons, and their decays produce a sizable fraction of the stable particles in the jet. Although the production and decay of some B and D hadrons, containing a b and c quark, respectively, have been studied inclusively [8] and the effect of heavy quark masses on the charged multiplicity in their jets has been observed [9,10], there remain large uncertainties in, e.g., B_s^0 and heavy baryon decays, and heavy quark masses. The removal of heavy flavor events will therefore allow the direct study of the hadronization of light quark jets, and measurements of c and b jets may shed additional light on some of the above issues.

A particularly interesting aspect of light quark hadronization is the question of what happens to the quark or antiquark that initiated a given jet. Recent experimental results [11–13] have confirmed the notion that it is ''contained'' as a valence constituent of a particular hadron, and that this ''leading'' hadron has on average a higher momentum than the other

TABLE I. Efficiencies for simulated events in the three flavor categories to be tagged as light (*uds*), *c* or *b* flavor, or none of these. The three rightmost columns indicate the composition of each simulated tagged sample assuming standard model relative flavor production.

	Efficiency for $Z^0\to$			Purity of $Z^0\to$		
	$u\bar{u}$, $d\bar{d}$, $s\bar{s}$	$c\bar{c}$	$b\bar{b}$	$u\bar{u}$, $d\bar{d}$, $s\bar{s}$	$c\bar{c}$	$b\bar{b}$
uds tag	0.703	0.190	0.010	0.925	0.070	0.005
c tag	0.061	0.551	0.105	0.204	0.638	0.158
b tag	0.001	0.024	0.815	0.005	0.023	0.973
no tag	0.235	0.235	0.070	0.721	0.201	0.078

hadrons of the same type in the jet. More precise measurements of such leading particle effects could lead to methods for identifying jets of a specific light primary flavor, with a number of applications in *ep* and hadron-hadron collisions as well as in e^+e^- annihilations.

There are several phenomenological models of jet production that combine modelling of all three stages of particle production and it is important to test their predictions. To model the parton production stage the HERWIG 5.8 [14], JETSET 7.4 [15] and UCLA 4.1 [16] event generators use a combination of first order matrix elements and a parton shower. To model the hadronization stage, the HERWIG model splits the gluons produced in the first stage into $q\bar{q}$ pairs, and these quarks and antiquarks are paired up locally to form colorless clusters that decay into the primary hadrons. The JETSET model takes a different approach, representing the color field between the partons by a semi-classical string; this string is broken, according to an iterative algorithm, into several pieces, each of which corresponds to a primary hadron. In the UCLA model whole events are generated according to weights derived from the phase space available to their final states and the relevant Clebsch-Gordan coefficients. Each of these models contains arbitrary parameters that control various aspects of hadronization and have been tuned to reproduce data from e^+e^- annihilations. The JETSET model includes a large number of parameters that control, on average, the species of primary hadron produced at each string break, giving it the potential to model the observed properties of identified hadron species in great detail. In the HERWIG model, clusters are decayed into pairs of primary hadrons according to phase space, and the relative production of different hadrons is effectively governed by two parameters that control the distribution of cluster masses. In the UCLA model, there is only one such free parameter, which controls the degree of locality of baryon-antibaryon pair formation.

In this paper we present an analysis of inclusive charged particle and identified $\pi^\pm$, $K^\pm$, and $p/\bar{p}$ production in hadronic Z^0 decays collected by the SLC Large Detector (SLD). The analysis is based upon the approximately 400 000 hadronic events obtained in runs of the SLC between 1996 and 1998, and supersedes our previous measurements [12]. We measure differential production rates for these hadron species in a flavor-inclusive sample of hadronic Z^0 decays and use the results to test the predictions of MLLA QCD +LPHD and of the three models just described. We also measure these differential production rates separately in Z^0 decays into light flavors ($u\bar{u}$, $d\bar{d}$ and $s\bar{s}$), $c\bar{c}$ and $b\bar{b}$, which provide improved tests of the QCD predictions and new tests of the models that separate the heavy hadron production and decay modelling from that of the rest of the hadronization process. In addition we measure hadron and antihadron differential production rates in light quark jets, thereby obtaining precise information on leading particles and additional stringent tests of the models.

In Sec. II we describe the SLD, in particular its precision vertexing and particle identification capabilities. Section III describes the selection of hadronic events, their separation into samples of different primary flavors using displaced vertices from heavy hadron decays, and the selection of light quark and antiquark hemispheres using the large production asymmetry in polar angle induced by the polarization of the SLC electron beam. In Sec. IV we present a measurement of the inclusive stable charged particle production rate. We describe the hadron species identification and present results for flavor-inclusive events in Sec. V, and Sec. VI presents results separately for light- ($Z^0\to u\bar{u},d\bar{d},s\bar{s}$), *c*- ($Z^0\to c\bar{c}$) and *b*-flavor ($Z^0\to b\bar{b}$) events. In Secs. VII and VIII we use these results to test the predictions of MLLA QCD +LPHD, and extract total yields of each hadron species per event of each flavor, respectively. We present measurements of leading particle production in light-flavor jets in Sec. IX and summarize the results in Sec. X.

II. THE SLD

This analysis of data from the SLD [17] used charged tracks measured in the central drift chamber (CDC) [18] and upgraded charge-coupled device vertex detector (VXD3) [19], and identified in the Cherenkov ring imaging detector (CRID) [12,20]. Energy deposits reconstructed in the liquid argon calorimeter [21] were used in the initial hadronic event selection and in the calculation of the event thrust [22] axis. Momentum measurement was provided by an axial magnetic field of 0.6 T. The CDC and VXD3 gave a momentum resolution of $\sigma_{p_\perp}/p_\perp = 0.01 \oplus 0.0026 p_\perp$, where $p_\perp$ is the track momentum transverse to the beam axis in GeV/c.

In the plane normal to the beamline the centroid of the micrometer-sized SLC interaction region was reconstructed from tracks in sets of approximately thirty sequential hadronic Z^0 decays to a precision of $\sigma_{IP}^{r\phi} \simeq 3\ \mu$m and used as an estimate of the primary interaction point (IP). The IP position

TABLE II. Production rate for all stable charged particles in terms of momentum p, scaled momentum x_p and $\xi = -\ln x_p$. For momentum we give statistical errors and the systematics arising from track-finding efficiency (including an overall 1% normalization term) and the sum of other sources, which is dominated by backgrounds (momentum resolution) at low (high) momentum; in the other columns these have been summed in quadrature.

Inclusive stable charged particle production rate					
p (GeV/c)	$\langle p \rangle$	ξ	$\frac{1}{N_{evts}}\frac{dn_{chg}}{dp}$ $\pm$stat. $\pm$ effic. $\pm$ other	$\frac{1}{N_{evts}}\frac{dn_{chg}}{dx_p}$	$\frac{1}{N_{evts}}\frac{dn_{chg}}{d\xi}$
0.25–0.35	0.300	4.87–5.21	$11.17 \pm 0.02 \pm 0.12 \pm 0.17$	509.2 ± 9.5	3.234 ± 0.062
0.35–0.45	0.400	4.62–4.87	$11.27 \pm 0.02 \pm 0.12 \pm 0.13$	513.9 ± 7.8	4.399 ± 0.068
0.45–0.55	0.499	4.42–4.62	$10.65 \pm 0.02 \pm 0.11 \pm 0.09$	485.5 ± 6.6	5.221 ± 0.072
0.55–0.65	0.599	4.25–4.42	$9.724 \pm 0.021 \pm 0.099 \pm 0.073$	443.4 ± 5.6	5.665 ± 0.073
0.65–0.75	0.699	4.11–4.25	$8.744 \pm 0.020 \pm 0.089 \pm 0.057$	398.7 ± 4.9	6.110 ± 0.075
0.75–1.00	0.870	3.82–4.11	$7.364 \pm 0.011 \pm 0.075 \pm 0.040$	335.8 ± 3.9	6.324 ± 0.074
1.00–1.25	1.121	3.60–3.82	$5.860 \pm 0.010 \pm 0.059 \pm 0.024$	267.2 ± 2.9	6.486 ± 0.072
1.25–1.50	1.371	3.41–3.60	$4.775 \pm 0.009 \pm 0.048 \pm 0.015$	217.8 ± 2.3	6.465 ± 0.070
1.50–1.75	1.622	3.26–3.41	$3.962 \pm 0.008 \pm 0.040 \pm 0.011$	180.7 ± 1.9	6.338 ± 0.068
1.75–2.00	1.872	3.13–3.26	$3.360 \pm 0.008 \pm 0.034 \pm 0.009$	153.2 ± 1.6	6.202 ± 0.066
2.00–2.25	2.122	3.01–3.13	$2.865 \pm 0.007 \pm 0.029 \pm 0.007$	130.6 ± 1.4	5.990 ± 0.063
2.25–2.50	2.375	2.90–3.01	$2.488 \pm 0.007 \pm 0.025 \pm 0.005$	113.4 ± 1.2	5.807 ± 0.062
2.50–2.75	2.625	2.81–2.90	$2.173 \pm 0.006 \pm 0.022 \pm 0.004$	99.1 ± 1.0	5.605 ± 0.059
2.75–3.00	2.875	2.72–2.81	$1.920 \pm 0.006 \pm 0.019 \pm 0.004$	87.57 ± 0.91	5.419 ± 0.058
3.00–3.25	3.125	2.64–2.72	$1.701 \pm 0.005 \pm 0.017 \pm 0.003$	77.56 ± 0.81	5.214 ± 0.056
3.25–3.50	3.375	2.57–2.64	$1.530 \pm 0.005 \pm 0.015 \pm 0.003$	69.78 ± 0.73	5.061 ± 0.054
3.50–3.75	3.625	2.50–2.57	$1.378 \pm 0.005 \pm 0.014 \pm 0.002$	62.84 ± 0.66	4.892 ± 0.053
3.75–4.00	3.875	2.43–2.50	$1.244 \pm 0.005 \pm 0.013 \pm 0.002$	56.72 ± 0.60	4.713 ± 0.051
4.00–4.50	4.244	2.32–2.43	$1.072 \pm 0.003 \pm 0.011 \pm 0.001$	48.90 ± 0.51	4.450 ± 0.047
4.50–5.00	4.744	2.21–2.32	$0.894 \pm 0.003 \pm 0.009 \pm 0.001$	40.78 ± 0.43	4.141 ± 0.045
5.00–5.50	5.244	2.12–2.21	$0.754 \pm 0.003 \pm 0.008 \pm 0.001$	34.39 ± 0.37	3.856 ± 0.042
5.50–6.50	5.975	1.95–2.12	$0.600 \pm 0.002 \pm 0.006 \pm 0.001$	27.35 ± 0.29	3.492 ± 0.038
6.50–7.50	6.977	1.81–1.95	$0.4502 \pm 0.0014 \pm 0.0049 \pm 0.0005$	20.53 ± 0.23	3.051 ± 0.035
7.50–8.50	7.980	1.68–1.81	$0.3429 \pm 0.0012 \pm 0.0039 \pm 0.0003$	15.64 ± 0.18	2.647 ± 0.032
8.50–9.50	8.982	1.57–1.68	$0.2687 \pm 0.0011 \pm 0.0032 \pm 0.0003$	12.25 ± 0.15	2.328 ± 0.029
9.50–10.50	9.980	1.47–1.57	$0.2120 \pm 0.0010 \pm 0.0026 \pm 0.0002$	9.67 ± 0.12	2.039 ± 0.027
10.5–11.5	10.99	1.38–1.47	$0.1700 \pm 0.0009 \pm 0.0022 \pm 0.0002$	7.75 ± 0.11	1.792 ± 0.025
11.5–12.5	11.99	1.29–1.38	$0.1351 \pm 0.0008 \pm 0.0018 \pm 0.0001$	6.161 ± 0.088	1.551 ± 0.023
12.5–13.5	12.99	1.22–1.29	$0.1103 \pm 0.0007 \pm 0.0016 \pm 0.0001$	5.029 ± 0.076	1.368 ± 0.022
13.5–14.5	13.99	1.15–1.22	$0.0889 \pm 0.0006 \pm 0.0013 \pm 0.0001$	4.053 ± 0.065	1.184 ± 0.020
14.5–16.0	15.22	1.05–1.15	$0.0688 \pm 0.0005 \pm 0.0011 \pm 0.0001$	3.139 ± 0.052	0.997 ± 0.017
16.0–17.5	16.72	0.96–1.05	$0.0513 \pm 0.0004 \pm 0.0009 \pm 0.0001$	2.338 ± 0.042	0.812 ± 0.016
17.5–19.0	18.23	0.88–0.96	$0.0383 \pm 0.0003 \pm 0.0007 \pm 0.0001$	1.748 ± 0.034	0.661 ± 0.014
19.0–20.5	19.72	0.80–0.88	$0.0291 \pm 0.0003 \pm 0.0006 \pm 0.0001$	1.326 ± 0.028	0.544 ± 0.012
20.5–22.0	21.21	0.73–0.80	$0.0221 \pm 0.0003 \pm 0.0005 \pm 0.0001$	1.008 ± 0.023	0.444 ± 0.011
22.0–24.0	22.96	0.64–0.73	$0.0159 \pm 0.0002 \pm 0.0004 \pm 0.0001$	0.724 ± 0.018	0.3468 ± 0.0090
24.0–26.0	24.94	0.56–0.64	$0.0105 \pm 0.0001 \pm 0.0003 \pm 0.0001$	0.480 ± 0.013	0.2513 ± 0.0073
26.0–30.0	27.74	0.42–0.56	$0.0063 \pm 0.0001 \pm 0.0002 \pm 0.0001$	0.285 ± 0.009	0.1700 ± 0.0056
30.0–35.0	32.14	0.26–0.42	$0.0025 \pm 0.0000 \pm 0.0001 \pm 0.0001$	0.114 ± 0.005	0.0800 ± 0.0035
35.0–45.0	38.23	0.01–0.26	$0.0005 \pm 0.0000 \pm 0.0000 \pm 0.0001$	0.024 ± 0.001	0.0207 ± 0.0013

along the beam axis was determined event by event using charged tracks, with an average resolution of $\sigma_{IP}^{z} \simeq 20$ μm. Including the uncertainty on the IP position, the resolution on the charged track impact parameter δ was parametrized in the plane perpendicular to the beamline as $\sigma_{\delta}^{r\phi} = 8 \oplus 29/(p \sin^{3/2}\theta)$ μm, and in any plane containing the beam axis as $\sigma_{\delta}^{z} = 9 \oplus 29/(p \sin^{3/2}\theta)$ μm, where θ is the track polar angle with respect to the beamline.

The barrel CRID covered the polar angle range $|\cos\theta| < 0.68$, and comprised liquid and gaseous radiator systems

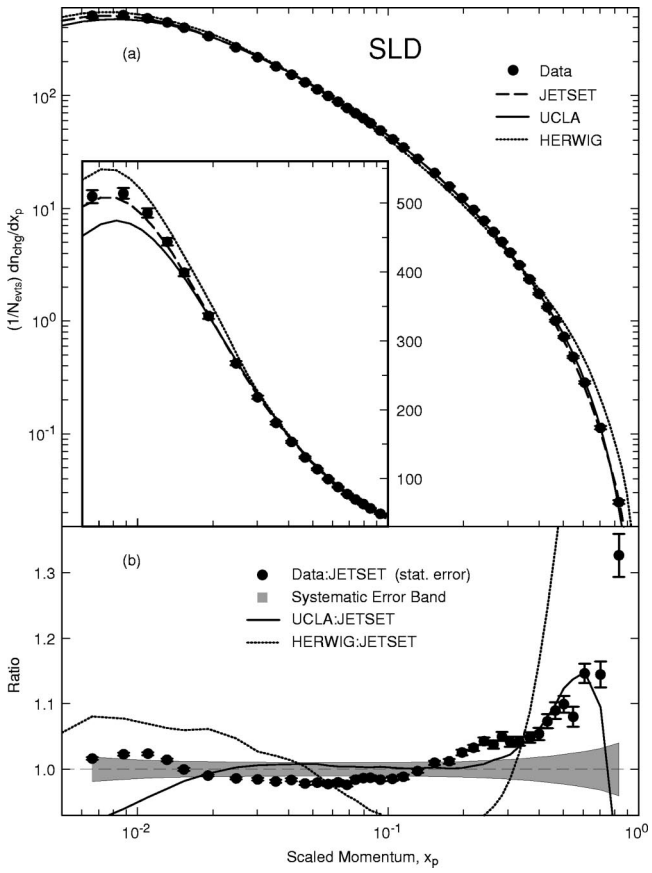

FIG. 1. (a) Charged track x_p distribution in hadronic events (dots) with logarithmic and (inset) linear vertical scales, compared with the predictions of three models (lines) using their respective default parameter values. (b) The same data divided by the JETSET prediction. The error bars in (b) are statistical and the shaded band represents the systematic uncertainty; all errors except an overall 1% normalization uncertainty are included in (a).

with refractive indices of 1.278 and 1.00176, respectively. In the high momentum limit, an average of 13 (10) Cherenkov photons was reconstructed per track traversing the liquid (gas) system, with an average Cherenkov angle resolution of 15 (4.3) mrad. The combination of these systems gave efficient and pure separation of pions, kaons, and protons over much of the kinematic range, as detailed in Sec. V.

III. EVENT SELECTION

The trigger and initial selection of hadronic events are described in [23]. The analysis presented here is based on charged tracks measured in the CDC and VXD3. A set of cuts was applied in order to select charged tracks and events well contained within the detector acceptance. Tracks were required to have (i) a closest approach to the beam axis within 5 cm, and within 10 cm along the beam axis of the measured IP; (ii) a polar angle θ with respect to the beam axis with $|\cos\theta|<0.80$; (iii) a momentum transverse to this axis $p_\perp>150$ MeV/c; and (iv) a momentum $p<50$ GeV/c. Events were required to have a minimum of seven such tracks; a visible energy $E_{vis}>18$ GeV, calculated from the accepted tracks, assigned the charged pion mass; a thrust axis polar angle θ_t with respect to the beam axis with $|\cos\theta_t|<0.71$; and the CDC and VXD3 active and a well-measured IP position. A sample of 284 494 events passed these cuts. For the identified hadrons, the CRID was also required to be active, giving a sample of 232 802 events. The non-hadronic background was estimated to be 0.1%, dominated by $Z^0\rightarrow\tau^+\tau^-$ events.

Samples of events enriched in light, c, and b primary flavors were selected using tracks that were measured well near the IP [24]. For each event we defined n_{sig} as the number of tracks with an impact parameter to the IP greater than three times its estimated error, $\delta>3\sigma_\delta$. We also ran a topological vertex finding algorithm [25] on the set of tracks in each hemisphere, and for each secondary vertex (significantly separated from the IP) found, we considered the total momentum P_{vtx} of the tracks assigned to the vertex and their invariant mass, calculated assuming each was from a charged pion and corrected [25] for missing momentum transverse to the line joining the IP and vertex position, M_{pt}. Cuts on these quantities were chosen [26] to maximize efficiency while maintaining reasonable sample purity: any event containing a vertex with $M_{pt}>2$ GeV/c^2 was assigned to the b-tagged sample; if not b tagged, an event was c tagged if either it contained a vertex with $0.5<M_{pt}<2$ GeV/c^2, $P_{vtx}>2$ GeV/c and $cP_{vtx}-14c^2M_{pt}>-10$ GeV, or it contained no secondary vertex but had $2\leqslant n_{sig}\leqslant 3$; events with no secondary vertex and $n_{sig}=0$ were assigned to the light-tagged sample; the remaining events were kept as an untagged sample. The 15% of the data taken in 1996 were excluded due to uncertainties in the simulation, and the light-, c- and b-tagged samples comprised roughly 83 000, 28 000 and 33 000 events, respectively. Selection efficiencies were calculated from a detailed detector simulation based on a version of JETSET tuned to the world's data on inclusive particle production and D and B hadron production and decay [27]; this tuning is essential for the reliable simulation of the flavor tagging, but has no effect on the track finding or particle identification. Efficiencies and sample purities are listed in Table I.

Samples of hemispheres enriched in light-quark and light-antiquark jets were selected by exploiting the large electroweak forward-backward production asymmetry with respect to the beam direction induced by the high polarization of the SLC electron beam. Here a looser light-flavor event tag of $n_{sig}=0$ was applied, giving a simulated *uds* efficiency of 84% and purity of 89%, with the background dominated by c-flavor events. The event thrust axis was used to approximate the initial $q\bar{q}$ axis and was signed such that its z component was along the electron beam direction, $\hat{t}_z>0$. Events in the central region of the detector, where the production asymmetry is small, were removed by the requirement $|\hat{t}_z|>0.15$, leaving 109 000 events. The quark-tagged hemisphere in events with left- (right-)handed electron beam polarization was defined to comprise the set of tracks with positive (negative) momentum projection along the signed thrust axis. The remaining tracks in each event were defined to be in the antiquark-tagged hemisphere. For the selected event

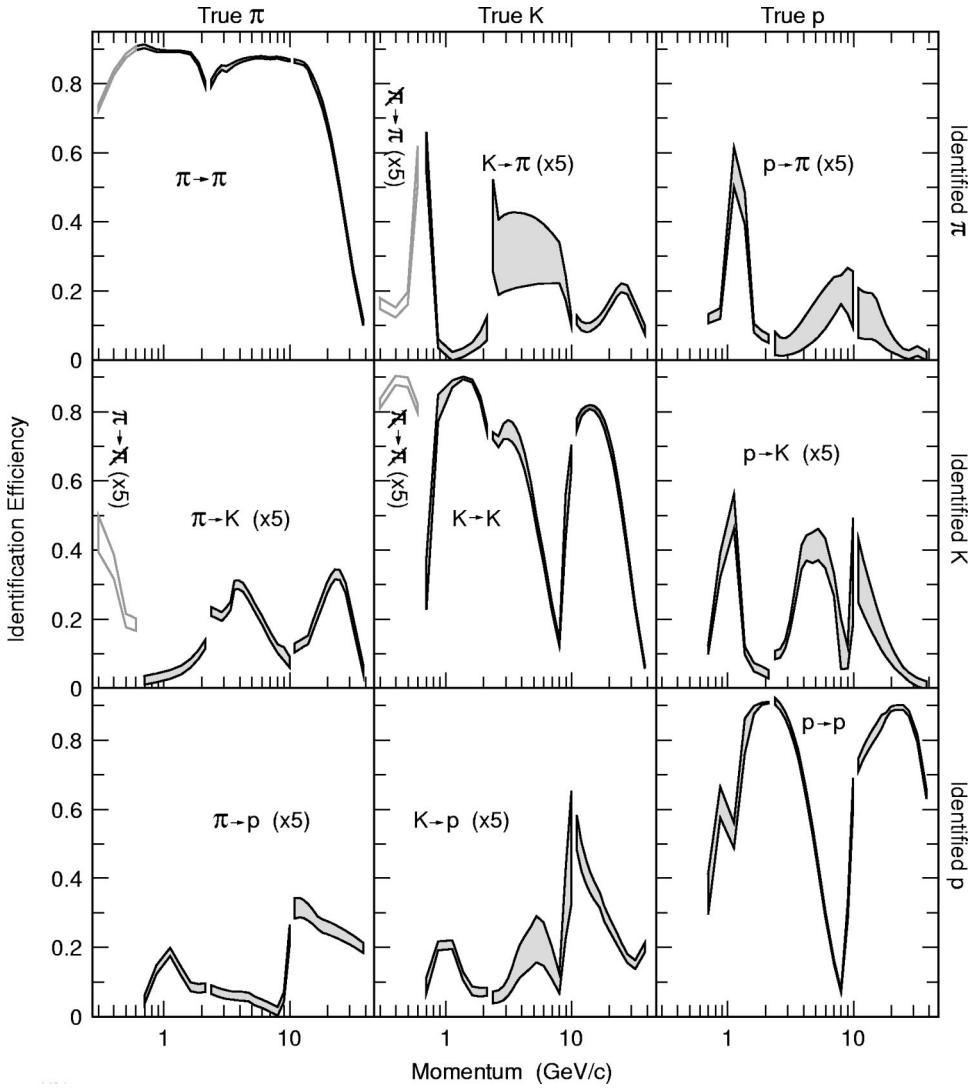

FIG. 2. Calibrated identification efficiencies for selected tracks. The half-widths of the gray bands represent the systematic uncertainties, which are strongly correlated between momenta. The off-diagonal efficiencies have been scaled by factors of five for clarity. The white bands in the four upper left plots represent the 2×2 matrix used below kaon threshold in the liquid system.

sample, the average magnitude of the e^- beam polarization was 0.73. Using this value and assuming standard model couplings, a tree-level calculation gives a quark (antiquark) purity of 0.73 in the quark- (antiquark-)tagged sample.

IV. STABLE CHARGED PARTICLE RATES

We first measured the differential production rate of all stable charged particles defined, following convention, as the sum of electrons, muons, pions, kaons and protons that are either primary hadrons or products of a chain of decays of particles with proper lifetime less than 3×10^{-10} s. Tracks satisfying the requirements in Sec. III were counted in momentum bins and corrected, using our detector simulation, for the track finding and selection efficiency and resolution, non-hadronic event background, and spurious tracks from beam-related backgrounds and interactions in the detector material. The resulting integrated rate was constrained to be 20.95 ± 0.21 charged particles per event, an average of previous measurements [10,28,29] in Z^0 decays; this corresponded to a 3.2% correction to our detector simulation.

The momentum dependence of the selection procedure was constrained by comparing the properties of reconstructed charged tracks in measured and simulated τ lepton decays, for which the momentum and particle type distributions are well known. We selected $e^+e^-\to\tau^+\tau^-$ events according to the criteria in [30], with the additional requirement of at least one track in the event with $p>7$ GeV/c, to reduce beam-related and two-photon event background. Comparisons were made for 1-, 3-, and 5-prong decays separately, as well as for hemispheres in which 2 or 4 tracks were found, giving constraints on both isolated tracks and those in close proximity to others. We also checked the momentum distributions for tagged electrons, muons, pions, and kaons, and the overall multiplicity distribution in selected events. In all cases the simulation was consistent with the data. Linear fits were made to the ratios of data:simulated distributions, and an uncertainty on the momentum dependence of $0.10\times|p-2.7\ \mathrm{GeV}/c|\%$ was assigned, reflecting the error on a typical fitted slope. Here p is the particle momentum in GeV/c and 2.7 GeV/c is the average momentum of all charged particles, so that this uncertainty is independent of the normalization uncertainty noted above. Variations of the background levels, detector material and momentum resolution were also considered and found to have much smaller

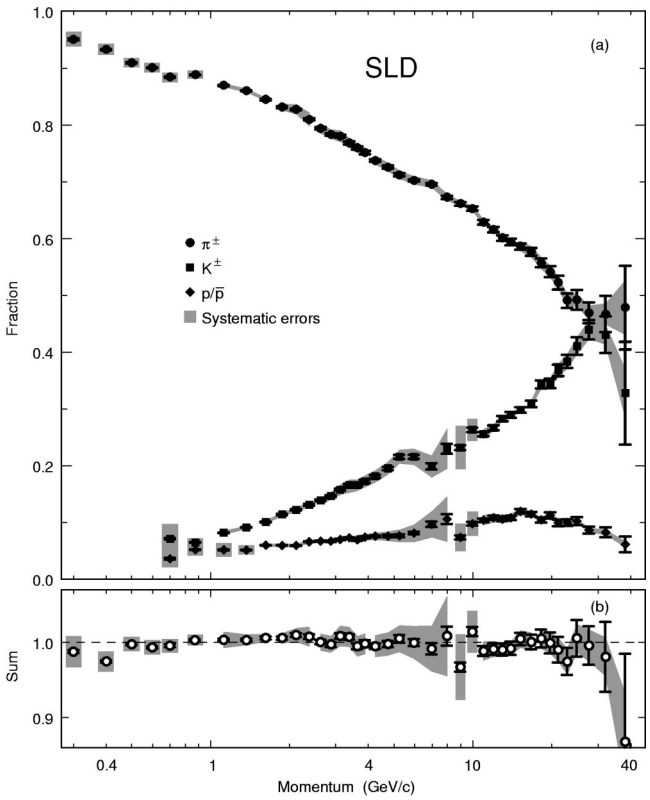

FIG. 3. (a) Measured charged hadron production fractions in hadronic Z^0 decays. The circles represent the $\pi^\pm$ fraction, the squares the $K^\pm$ fraction, the diamonds the $p/\bar{p}$ fraction. The error bars are statistical only; the shaded areas indicate the systematic errors, and are connected across momentum regions where there is a strong positive correlation. (b) The sum of the three fractions.

effects, except at the lowest and highest momenta.

The inclusive charged particle differential production rate is listed in Table II, in terms of the variables momentum p, scaled momentum $x_p = 2p/E_{CM}$ and $\xi = -\ln x_p$. The x_p distribution is shown in Fig. 1, and compared with the predictions of the JETSET, UCLA and HERWIG event generators described in Sec. I, using the default parameter values for each model. The JETSET model is the most consistent with the data, but predicts a slightly softer spectrum; the UCLA model describes the data over most of the x_p region, but falls low for $x_p > 0.7$ and $x_p < 0.015$ units; the HERWIG model predicts a spectrum quite different from that seen in the data. Our results for the shape of the spectrum are consistent with those published previously [10,28,29]

V. CHARGED PION, KAON, AND PROTON PRODUCTION

Additional track selection cuts [26] were applied to remove tracks that interacted or scattered through large angles before exiting the CRID and to ensure that the CRID performance was well modelled by the simulation. Good information from the liquid (gas) system was required for tracks in the "liquid" ("gas") region, with momenta below 2.25 GeV/c (above 10.5 GeV/c); good information from both systems was required for momenta in the "combined" region, 2.25–10.5 GeV/c. Tracks were required to have at least 40 CDC hits, at least one of which was at a radius of at least 92 cm, to extrapolate through an active region of the appropriate radiator(s), and to have at least 40 (70) % of their expected liquid (gas) ring contained within a sensitive region of the CRID time projection chamber (TPCs). The latter requirement included dead and inefficient regions within the TPCs and also rejected tracks with $p > 2.25$ GeV/c for which there was a saturated CRID hit within a 2 cm radius ($\sim$ the maximum ring radius) of a point 1 cm behind the expected gas ring center in the photoelectron drift direction. Saturated hits arise from the passage of tracks through the TPCs and from various backgrounds, and mask single photoelectron hits in their vicinity. Liquid region tracks that extrapolated through an active TPC were required to have a saturated hit within 1 cm of the extrapolated track, indicating that the track traversed the TPC and hence the liquid radiator; gas region tracks were required to have either such a saturated hit or the presence of at least four hits consistent with a liquid ring. These cuts accepted 71%, 68% and 74% of the tracks within the CRID geometrical acceptance in the liquid, combined and gas regions, respectively.

Tracks were identified using a global likelihood technique [12,31]. For each track and each of the five hypotheses $i = e, \mu, \pi, K, p$, a likelihood L_i was calculated based upon the number of detected photoelectrons and their measured Cherenkov angles; the expected number of photons; the expected Cherenkov angle; and terms accounting for random backgrounds and hits consistent with Cherenkov radiation from other tracks in the event. The best hypothesis for each track was used to determine its contribution to the background for other tracks, and the calculation was iterated until there was no change in any best hypothesis. Particle separation was based upon differences between logarithms of the three likelihoods, $\mathcal{L}_i = \ln L_i$, $i = \pi, K, p$. A track in the liquid (combined, gas) region was identified as species j if $\mathcal{L}_j$ exceeded both of the other log-likelihoods by at least 3 (2) units. The electron and muon likelihoods are generally quite similar to the pion likelihood, and the leptons were included in the pion category at this stage.

We quantified the identification performance in terms of a momentum-dependent identification efficiency matrix **E**, each element E_{ij} of which represents the probability that the selected track from a true i hadron is identified as a j hadron, with $i, j = \pi, K, p$. The elements of this matrix were determined where possible from the data [26,32]. For example, tracks from selected K^0_s and τ decays were used as "pion" test samples, having estimated kaon plus proton contents of 0.3% and 1.7% respectively. Probabilities for such tracks to be identified as pions, kaons and protons were compared with the predictions of our detailed Monte Carlo simulation, which was found to describe the momentum dependence of the efficiencies well and to reproduce their amplitudes to within a few percent. Functional forms, chosen to describe the momentum dependence of both data and simulated test samples as well as that of simulated true pions in hadronic events, were fitted to the data, except for momenta below 0.8 GeV/c where there was significant structure on the scale of the bin size and bin-by-bin corrections were used. The simulation was used to correct the fitted parameters for non-

TABLE III. Charged pion fraction f_π and differential production rate $(1/N_{evts})dn_\pi/dx_p$ per hadronic Z^0 decay. The first error is statistical, the second systematic. $\langle x_p \rangle$ is the average x_p value of charged tracks in each bin. The last row gives the integral over the x_p range of the measurement. A 1.0% normalization uncertainty is included in the systematic error on the integral, but not in those on the differential rate.

x_p range	$\langle x_p \rangle$	f_π	$(1/N_{evts})dn_\pi/dx_p$
0.005–0.008	0.0066	$0.951 \pm 0.002 \pm 0.010$	$471.8 \pm 1.3 \pm 9.2$
0.008–0.010	0.0088	$0.933 \pm 0.001 \pm 0.007$	$470.4 \pm 1.1 \pm 6.6$
0.010–0.012	0.0109	$0.910 \pm 0.001 \pm 0.006$	$434.6 \pm 1.1 \pm 5.0$
0.012–0.014	0.0131	$0.901 \pm 0.001 \pm 0.006$	$388.8 \pm 1.0 \pm 4.0$
0.014–0.016	0.0153	$0.885 \pm 0.001 \pm 0.006$	$352.7 \pm 0.9 \pm 3.3$
0.016–0.022	0.0191	$0.888 \pm 0.001 \pm 0.004$	$294.8 \pm 0.5 \pm 2.2$
0.022–0.027	0.0246	$0.870 \pm 0.001 \pm 0.003$	$229.6 \pm 0.5 \pm 1.3$
0.027–0.033	0.0301	$0.860 \pm 0.001 \pm 0.003$	$185.0 \pm 0.4 \pm 0.9$
0.033–0.038	0.0356	$0.845 \pm 0.001 \pm 0.003$	$150.6 \pm 0.4 \pm 0.7$
0.038–0.044	0.0411	$0.831 \pm 0.002 \pm 0.005$	$125.6 \pm 0.4 \pm 0.9$
0.044–0.049	0.0465	$0.828 \pm 0.002 \pm 0.008$	$106.5 \pm 0.4 \pm 1.1$
0.049–0.055	0.0521	$0.810 \pm 0.002 \pm 0.012$	$90.40 \pm 0.35 \pm 1.33$
0.055–0.060	0.0576	$0.794 \pm 0.002 \pm 0.009$	$77.38 \pm 0.31 \pm 0.91$
0.060–0.066	0.0630	$0.784 \pm 0.002 \pm 0.008$	$67.39 \pm 0.29 \pm 0.70$
0.066–0.071	0.0685	$0.780 \pm 0.002 \pm 0.010$	$59.40 \pm 0.27 \pm 0.75$
0.071–0.077	0.0740	$0.768 \pm 0.003 \pm 0.009$	$52.57 \pm 0.25 \pm 0.60$
0.077–0.082	0.0795	$0.760 \pm 0.003 \pm 0.008$	$46.76 \pm 0.24 \pm 0.50$
0.082–0.088	0.0850	$0.752 \pm 0.003 \pm 0.008$	$41.70 \pm 0.23 \pm 0.43$
0.088–0.099	0.0931	$0.738 \pm 0.002 \pm 0.007$	$35.26 \pm 0.15 \pm 0.36$
0.099–0.110	0.1040	$0.726 \pm 0.002 \pm 0.007$	$28.89 \pm 0.13 \pm 0.29$
0.110–0.121	0.1150	$0.712 \pm 0.003 \pm 0.007$	$23.88 \pm 0.12 \pm 0.25$
0.121–0.143	0.1310	$0.702 \pm 0.002 \pm 0.007$	$18.69 \pm 0.08 \pm 0.19$
0.143–0.164	0.1530	$0.696 \pm 0.002 \pm 0.006$	$13.85 \pm 0.07 \pm 0.14$
0.164–0.186	0.1750	$0.673 \pm 0.003 \pm 0.006$	$10.16 \pm 0.06 \pm 0.11$
0.186–0.208	0.1970	$0.662 \pm 0.003 \pm 0.004$	$7.812 \pm 0.050 \pm 0.069$
0.208–0.230	0.2189	$0.653 \pm 0.004 \pm 0.004$	$6.076 \pm 0.044 \pm 0.061$
0.230–0.252	0.2410	$0.629 \pm 0.004 \pm 0.005$	$4.674 \pm 0.039 \pm 0.053$
0.252–0.274	0.2629	$0.616 \pm 0.005 \pm 0.005$	$3.632 \pm 0.035 \pm 0.044$
0.274–0.296	0.2849	$0.601 \pm 0.005 \pm 0.004$	$2.886 \pm 0.031 \pm 0.037$
0.296–0.318	0.3068	$0.594 \pm 0.006 \pm 0.004$	$2.292 \pm 0.028 \pm 0.031$
0.318–0.351	0.3338	$0.586 \pm 0.006 \pm 0.009$	$1.749 \pm 0.021 \pm 0.034$
0.351–0.384	0.3666	$0.577 \pm 0.007 \pm 0.010$	$1.275 \pm 0.018 \pm 0.028$
0.384–0.417	0.3997	$0.557 \pm 0.008 \pm 0.010$	$0.921 \pm 0.016 \pm 0.022$
0.417–0.450	0.4325	$0.542 \pm 0.010 \pm 0.010$	$0.680 \pm 0.014 \pm 0.018$
0.450–0.482	0.4651	$0.523 \pm 0.012 \pm 0.011$	$0.499 \pm 0.013 \pm 0.014$
0.482–0.526	0.5035	$0.491 \pm 0.013 \pm 0.010$	$0.338 \pm 0.010 \pm 0.010$
0.526–0.570	0.5470	$0.492 \pm 0.018 \pm 0.011$	$0.226 \pm 0.009 \pm 0.007$
0.570–0.658	0.6083	$0.469 \pm 0.018 \pm 0.012$	$0.130 \pm 0.005 \pm 0.005$
0.658–0.768	0.7047	$0.467 \pm 0.032 \pm 0.018$	$0.0526 \pm .0037 \pm .0029$
0.768–1.000	0.8383	$0.479 \pm 0.074 \pm 0.048$	$0.0113 \pm .0018 \pm .0013$
0.005–1.000			$15.74 \pm 0.01 \pm 0.17$

pion content in the test samples and differences in tracking performance between tracks in these samples and those passing our selection cuts in hadronic events. The resulting identification efficiency functions, $E_{\pi\pi}$, $E_{\pi K}$ and $E_{\pi p}$, are shown in the leftmost column of Fig. 2. A similar procedure using only π and p likelihoods was used to measure the π-p separation in the liquid (gas) system for $p > 2(17)$ GeV/c. This information was combined with that from test samples of protons and kaons from decays of Λ^0 hyperons and ϕ mesons, respectively, to derive the remaining efficiencies in Fig. 2.

The bands in Fig. 2 encompass the upper and lower systematic error bounds on the efficiencies. There are discontinuities between the liquid, combined and gas regions, and there are strong point-to-point correlations within each region. For the diagonal elements, these errors correspond to

TABLE IV. Charged kaon fraction and differential production rate per hadronic Z^0 decay.

x_p range	$\langle x_p \rangle$	f_K	$(1/N_{evts})\ dn_K/dx_p$
0.014–0.016	0.0153	$0.072 \pm 0.002 \pm 0.023$	$28.59 \pm 0.64 \pm 9.26$
0.016–0.022	0.0191	$0.065 \pm 0.001 \pm 0.005$	$21.57 \pm 0.20 \pm 1.57$
0.022–0.027	0.0246	$0.082 \pm 0.001 \pm 0.003$	$21.62 \pm 0.19 \pm 0.80$
0.027–0.033	0.0301	$0.091 \pm 0.001 \pm 0.002$	$19.65 \pm 0.18 \pm 0.53$
0.033–0.038	0.0356	$0.101 \pm 0.001 \pm 0.002$	$18.02 \pm 0.16 \pm 0.44$
0.038–0.044	0.0411	$0.114 \pm 0.001 \pm 0.003$	$17.27 \pm 0.17 \pm 0.43$
0.044–0.049	0.0465	$0.123 \pm 0.001 \pm 0.004$	$15.78 \pm 0.17 \pm 0.47$
0.049–0.055	0.0521	$0.131 \pm 0.002 \pm 0.004$	$14.664 \pm 0.194 \pm 0.442$
0.055–0.060	0.0576	$0.139 \pm 0.002 \pm 0.005$	$13.535 \pm 0.189 \pm 0.503$
0.060–0.066	0.0630	$0.147 \pm 0.002 \pm 0.006$	$12.599 \pm 0.176 \pm 0.558$
0.066–0.071	0.0685	$0.158 \pm 0.002 \pm 0.008$	$12.036 \pm 0.165 \pm 0.635$
0.071–0.077	0.0740	$0.166 \pm 0.002 \pm 0.009$	$11.349 \pm 0.162 \pm 0.622$
0.077–0.082	0.0795	$0.166 \pm 0.003 \pm 0.010$	$10.207 \pm 0.164 \pm 0.603$
0.082–0.088	0.0850	$0.172 \pm 0.033 \pm 0.010$	$9.571 \pm 0.160 \pm 0.566$
0.088–0.099	0.0931	$0.181 \pm 0.002 \pm 0.011$	$8.671 \pm 0.113 \pm 0.505$
0.099–0.110	0.1040	$0.196 \pm 0.003 \pm 0.011$	$7.784 \pm 0.114 \pm 0.440$
0.110–0.121	0.1150	$0.216 \pm 0.004 \pm 0.012$	$7.237 \pm 0.120 \pm 0.395$
0.121–0.143	0.1310	$0.216 \pm 0.003 \pm 0.014$	$5.746 \pm 0.089 \pm 0.369$
0.143–0.164	0.1530	$0.199 \pm 0.005 \pm 0.019$	$3.959 \pm 0.102 \pm 0.381$
0.164–0.186	0.1750	$0.230 \pm 0.009 \pm 0.035$	$3.473 \pm 0.134 \pm 0.532$
0.186–0.208	0.1970	$0.232 \pm 0.004 \pm 0.035$	$2.739 \pm 0.047 \pm 0.419$
0.208–0.230	0.2189	$0.264 \pm 0.004 \pm 0.017$	$2.452 \pm 0.037 \pm 0.163$
0.230–0.252	0.2410	$0.256 \pm 0.004 \pm 0.008$	$1.903 \pm 0.030 \pm 0.063$
0.252–0.274	0.2629	$0.267 \pm 0.004 \pm 0.006$	$1.574 \pm 0.027 \pm 0.036$
0.274–0.296	0.2849	$0.283 \pm 0.005 \pm 0.004$	$1.360 \pm 0.024 \pm 0.026$
0.296–0.318	0.3068	$0.290 \pm 0.005 \pm 0.004$	$1.118 \pm 0.022 \pm 0.020$
0.318–0.351	0.3338	$0.298 \pm 0.005 \pm 0.004$	$0.890 \pm 0.016 \pm 0.017$
0.351–0.384	0.3666	$0.309 \pm 0.006 \pm 0.006$	$0.683 \pm 0.014 \pm 0.016$
0.384–0.417	0.3997	$0.343 \pm 0.007 \pm 0.008$	$0.567 \pm 0.013 \pm 0.015$
0.417–0.450	0.4325	$0.345 \pm 0.008 \pm 0.009$	$0.433 \pm 0.012 \pm 0.014$
0.450–0.482	0.4651	$0.368 \pm 0.010 \pm 0.011$	$0.351 \pm 0.011 \pm 0.012$
0.482–0.526	0.5035	$0.384 \pm 0.011 \pm 0.013$	$0.264 \pm 0.008 \pm 0.010$
0.526–0.570	0.5470	$0.411 \pm 0.016 \pm 0.015$	$0.188 \pm 0.008 \pm 0.008$
0.570–0.658	0.6083	$0.439 \pm 0.017 \pm 0.017$	$0.122 \pm 0.005 \pm 0.006$
0.658–0.768	0.7047	$0.431 \pm 0.032 \pm 0.018$	$0.0485 \pm .0037 \pm .0027$
0.768–1.000	0.8383	$0.328 \pm 0.090 \pm 0.042$	$0.0078 \pm .0022 \pm .0011$
0.014–1.000			$2.074 \pm 0.006 \pm 0.066$

those on the fitted parameters, and fall into four categories. The overall amplitude of the efficiency is driven by the average and rms of the number of photons detected per track; the corresponding uncertainty is common to all momenta in a given region and is also correlated between the liquid and combined regions for protons and the combined and gas regions for pions. The Cherenkov angle resolution affects the positions of the falling edges; its uncertainty is therefore correlated across the range of a given edge, as well as between pions and kaons in the regions 1.5–2.25 GeV/c and 12–45 GeV/c, and between kaons and protons in the region 3–8 GeV/c. Performance near a Cherenkov threshold depends on the relevant index of refraction and its stability; since the efficiencies change rapidly on the scale of our bin sizes, bin-by-bin calibrations were done for E_{KK} (E_{pp}) in the range $p<1$ GeV/c ($p<1.5$ GeV/c), and E_{KK} and E_{pp} in the range $7.5<p<10.5$ GeV/c, which are completely independent. In the region 10.5–18 GeV/c, protons are below threshold in the gas, whereas pions and kaons are well above threshold; here E_{pp} depends largely on the background level, and its uncertainty is correlated across this region.

For the off-diagonal elements, representing misidentification rates, the errors on fitted parameters were also used, but subject to a minimum value of 0.0025 to account for the limited statistics of the test sample constraints on the momentum dependences. These uncertainties correspond to a combination of the effects listed above for the diagonal elements, and each is typically dominated by one effect in a given momentum region. These errors should also therefore be considered strongly positively correlated across each of

TABLE V. Proton plus antiproton fraction and differential production rate per hadronic Z^0 decay.

x_p range	$\langle x_p \rangle$	f_p	$(1/N_{evts})\ dn_p/dx_p$
0.014–0.016	0.0153	$0.036 \pm 0.001 \pm 0.013$	$14.51 \pm 0.52 \pm 5.08$
0.016–0.022	0.0191	$0.052 \pm 0.001 \pm 0.008$	$17.32 \pm 0.27 \pm 2.58$
0.022–0.027	0.0246	$0.052 \pm 0.001 \pm 0.009$	$13.75 \pm 0.29 \pm 2.50$
0.027–0.033	0.0301	$0.052 \pm 0.001 \pm 0.006$	$11.12 \pm 0.17 \pm 1.24$
0.033–0.038	0.0356	$0.060 \pm 0.001 \pm 0.003$	$10.75 \pm 0.14 \pm 0.47$
0.038–0.044	0.0411	$0.060 \pm 0.001 \pm 0.002$	$9.048 \pm 0.123 \pm 0.350$
0.044–0.049	0.0465	$0.060 \pm 0.001 \pm 0.002$	$7.669 \pm 0.117 \pm 0.298$
0.049–0.055	0.0521	$0.066 \pm 0.001 \pm 0.003$	$7.410 \pm 0.113 \pm 0.294$
0.055–0.060	0.0576	$0.068 \pm 0.001 \pm 0.003$	$6.587 \pm 0.109 \pm 0.259$
0.060–0.066	0.0630	$0.067 \pm 0.001 \pm 0.003$	$5.788 \pm 0.105 \pm 0.238$
0.066–0.071	0.0685	$0.070 \pm 0.001 \pm 0.003$	$5.344 \pm 0.100 \pm 0.228$
0.071–0.077	0.0740	$0.073 \pm 0.002 \pm 0.003$	$4.987 \pm 0.104 \pm 0.229$
0.077–0.082	0.0795	$0.069 \pm 0.002 \pm 0.004$	$4.278 \pm 0.100 \pm 0.242$
0.082–0.088	0.0850	$0.074 \pm 0.002 \pm 0.005$	$4.117 \pm 0.101 \pm 0.253$
0.088–0.099	0.0931	$0.076 \pm 0.002 \pm 0.006$	$3.633 \pm 0.072 \pm 0.269$
0.099–0.110	0.1040	$0.076 \pm 0.002 \pm 0.008$	$3.036 \pm 0.076 \pm 0.300$
0.110–0.121	0.1150	$0.077 \pm 0.002 \pm 0.011$	$2.568 \pm 0.081 \pm 0.357$
0.121–0.143	0.1310	$0.081 \pm 0.003 \pm 0.015$	$2.165 \pm 0.069 \pm 0.398$
0.143–0.164	0.1530	$0.097 \pm 0.005 \pm 0.023$	$1.931 \pm 0.096 \pm 0.452$
0.164–0.186	0.1750	$0.106 \pm 0.009 \pm 0.039$	$1.603 \pm 0.133 \pm 0.594$
0.186–0.208	0.1970	$0.074 \pm 0.004 \pm 0.022$	$0.871 \pm 0.045 \pm 0.255$
0.208–0.230	0.2189	$0.098 \pm 0.003 \pm 0.019$	$0.912 \pm 0.030 \pm 0.179$
0.230–0.252	0.2410	$0.104 \pm 0.003 \pm 0.008$	$0.775 \pm 0.025 \pm 0.062$
0.252–0.274	0.2629	$0.108 \pm 0.004 \pm 0.007$	$0.639 \pm 0.022 \pm 0.044$
0.274–0.296	0.2849	$0.106 \pm 0.004 \pm 0.007$	$0.511 \pm 0.019 \pm 0.033$
0.296–0.318	0.3068	$0.109 \pm 0.004 \pm 0.006$	$0.419 \pm 0.016 \pm 0.024$
0.318–0.351	0.3338	$0.120 \pm 0.004 \pm 0.006$	$0.358 \pm 0.011 \pm 0.018$
0.351–0.384	0.3666	$0.115 \pm 0.004 \pm 0.005$	$0.254 \pm 0.009 \pm 0.012$
0.384–0.417	0.3997	$0.105 \pm 0.004 \pm 0.005$	$0.173 \pm 0.008 \pm 0.008$
0.417–0.450	0.4325	$0.113 \pm 0.005 \pm 0.004$	$0.141 \pm 0.007 \pm 0.005$
0.450–0.482	0.4651	$0.100 \pm 0.006 \pm 0.003$	$0.0950 \pm .0055 \pm .0036$
0.482–0.526	0.5035	$0.100 \pm 0.005 \pm 0.003$	$0.0688 \pm .0039 \pm .0027$
0.526–0.570	0.5470	$0.103 \pm 0.007 \pm 0.003$	$0.0470 \pm .0032 \pm .0018$
0.570–0.658	0.6083	$0.087 \pm 0.006 \pm 0.003$	$0.0241 \pm .0017 \pm .0012$
0.658–0.768	0.7047	$0.083 \pm 0.009 \pm 0.004$	$0.0093 \pm .0010 \pm .0006$
0.768–1.000	0.8383	$0.062 \pm 0.014 \pm 0.005$	$0.0015 \pm .0003 \pm .0001$
0.014–1.000			$0.984 \pm 0.006 \pm 0.035$

the liquid, combined and gas regions.

The identification efficiencies in Fig. 2 peak near or above 0.9 and are greater than 0.8 over wide ranges. The pion (kaon and proton) coverage is continuous from 0.25 GeV/c (0.65 GeV/c) up to the beam energy, although the efficiencies fall below 0.2 for pions and kaons above about 30 GeV/c, and kaons and protons in the 6–9 GeV/c range. Misidentification rates are typically at the few percent level, with peak values of up to 0.1.

In each momentum bin we measured the fractions of the selected tracks that were identified as pions, kaons and protons. The observed fractions were related to the true production fractions by an efficiency matrix, composed of the values shown in Fig. 2. This matrix was inverted and used to unfold our observed identified hadron fractions. This analysis procedure does not require that the sum of the charged hadron fractions be unity; instead the sum was used as a consistency check, which was found to be satisfied at all momenta (see Fig. 3). For momenta below 0.65 GeV/c, we could not distinguish kaons from protons, but pions could be identified down to 0.25 GeV/c. An analogous 2×2 analysis of pions and non-pions was used in this region, and we present only the pion fraction.

The background from electrons and muons was estimated from the simulation to be about 5% of the tracks in the inclusive flavor sample, predominantly from c- and b-flavor events. The fractions were corrected for the lepton backgrounds using the simulation, as well as for the effects of beam-related backgrounds, particles interacting in the detector material, and particles decaying outside the tracking vol-

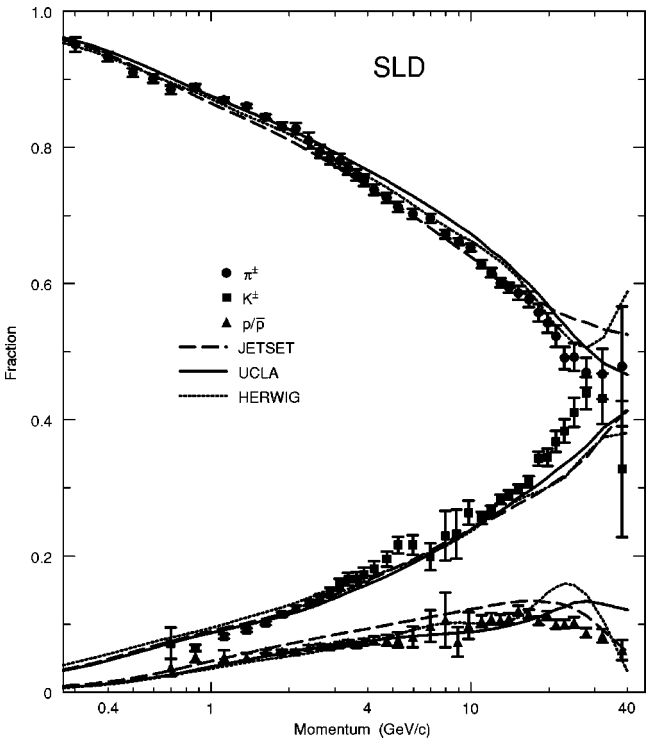

FIG. 4. Comparison of our measured charged hadron fractions (symbols) with the predictions of the JETSET (dashed lines), UCLA (solid lines) and HERWIG (dotted lines) models.

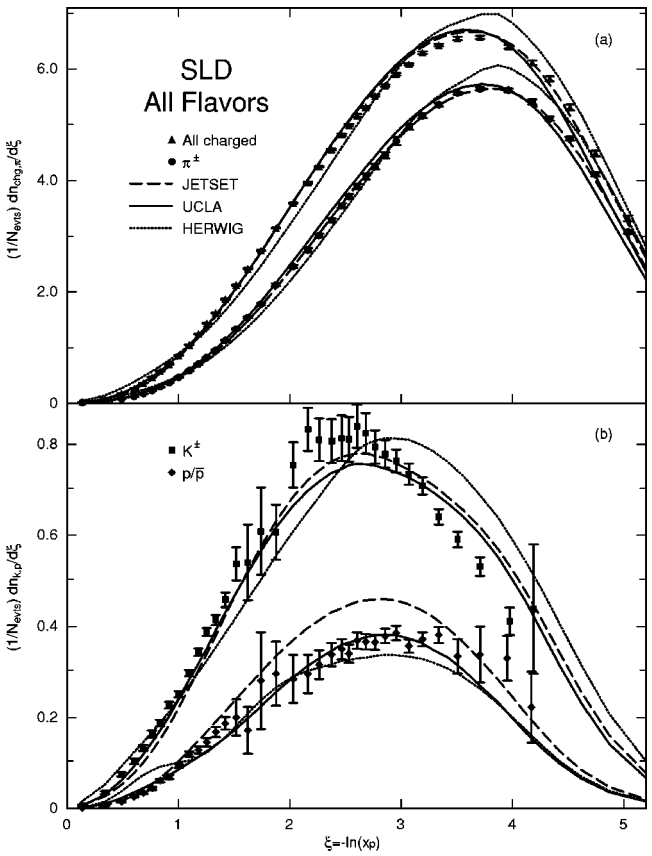

FIG. 5. Comparison of our measured (a) charged particle and pion and (b) kaon and proton production rates (symbols) with the predictions of the JETSET (dashed lines), UCLA (solid lines) and HERWIG (dotted lines) models.

ume. The conventional definition of a final-state charged hadron was used, namely a charged pion, kaon or proton that is either from the primary interaction or a product of a chain of decays of hadrons each with a proper lifetime less than 3×10^{-10} s.

The measured charged hadron fractions in inclusive hadronic Z^0 decays are shown in Fig. 3 and listed in Tables III–V. The systematic errors were determined by propagating the errors on the calibrated efficiency matrix and are strongly correlated as described above. They are indicated by the shaded regions in Fig. 3. The errors on the points below $\sim$15 GeV/c are dominated by the systematic uncertainties; those above $\sim$30 GeV/c are dominated by statistical uncertainties. The sum of fractions is consistent with unity everywhere.

Pions are seen to dominate the charged hadron production at low momentum, and to decline steadily in fraction as momentum increases. The kaon fraction rises steadily, approaching the pion fraction at high momentum. The proton fraction rises to a plateau value of about one-tenth at about 10 GeV/c, then declines at the highest momenta.

In Fig. 4 we compare our results with the predictions of the JETSET, UCLA and HERWIG models. The momentum dependence for each of the three hadron species is reproduced qualitatively by all models. All three models overestimate the kaon fraction for momenta below about 1.5 GeV/c, and underestimate it for momenta above about 15 GeV/c. The UCLA model overestimates the pion fraction by about 2σ (taking into account the correlation in the experimental errors) in the 5–15 GeV/c range, but is the only model consistent with the behavior above about 20 GeV/c. The JETSET model overestimates the proton fraction at all momenta, but describes the momentum dependence; HERWIG and UCLA predict behavior at high momentum that is inconsistent with the proton data.

Since the JETSET model has a number of parameters that control specific aspects of relative particle production, we have investigated some simple changes. We find that reducing the diquark:quark ratio [PARJ(1)] from 0.10 to 0.08, while leaving all other parameters unchanged, results in a good description of the proton fraction. The kaon fraction is sensitive to both the s-quark probability [PARJ(2)] and the strange vector:pseudoscalar ratio [PARJ(12)]; however, no combination of these two parameters was found to give a good description of the kaon fraction over the entire measured momentum range.

In order to obtain charged hadron production rates, the simulation was used to subtract the contribution of all particles (mostly leptons) that were not charged pions, kaons or protons from our measured total charged production rate (table II). The pion, kaon and proton fractions were multiplied by this adjusted rate to obtain the individual rates tabulated in Tables III–V as a function of scaled momentum. In Fig. 5 we compare our ξ distributions with the predictions of the three models. The features for pions are similar to those for all charged tracks: all models describe the data qualita-

TABLE VI. Differential production rates $(1/N_{evts})dn_{\pi^\pm}/dx_p$ for charged pions per Z^0 decay into light (u, d, s), c and b primary flavors. The errors are the sums in quadrature of statistical errors and those systematic uncertainties arising from the unfolding procedure; systematic errors common to the three flavors are not included in the rates and cancel in the ratios. The $\langle x_p \rangle$ values for the three flavor samples are consistent, and have been averaged.

	Pion production rates			Ratios	
$\langle x_p \rangle$	$u\bar{u},d\bar{d},s\bar{s}$	$c\bar{c}$	$b\bar{b}$	c:uds	b:uds
0.0066	474.0±13.9	425.5±26.6	478.1±15.8	0.898±0.058	1.009±0.019
0.0088	467.3±10.5	440.5±23.2	488.4±11.9	0.943±0.054	1.045±0.017
0.0109	418.2±8.4	453.8±20.0	463.7±9.5	1.085±0.055	1.109±0.018
0.0131	375.5±6.9	409.2±17.2	432.2±7.7	1.090±0.053	1.151±0.019
0.0153	327.7±5.7	372.8±14.6	382.4±6.5	1.138±0.052	1.167±0.019
0.0191	275.8±4.2	306.4±10.9	333.3±4.6	1.111±0.047	1.208±0.017
0.0246	216.0±3.0	234.6±8.0	261.7±3.3	1.086±0.044	1.212±0.018
0.0301	171.2±2.2	197.4±6.2	214.2±2.7	1.153±0.043	1.251±0.017
0.0356	140.4±1.9	155.8±6.0	175.2±2.3	1.109±0.050	1.247±0.016
0.0411	116.4±1.5	132.5±4.3	145.4±1.9	1.38±0.043	1.249±0.018
0.0465	99.9±1.2	109.3±3.5	121.4±1.7	1.095±0.040	1.215±0.018
0.0521	85.4±1.0	92.9±2.9	103.3±1.5	1.088±0.038	1.209±0.019
0.0576	72.85±0.89	77.56±2.48	89.24±1.36	1.065±0.037	1.225±0.020
0.0630	64.51±0.79	68.23±2.17	75.47±1.21	1.058±0.036	1.170±0.020
0.0685	56.82±0.72	60.06±1.97	65.97±1.12	1.057±0.038	1.161±0.022
0.0740	50.84±0.66	51.76±1.81	59.39±1.04	1.018±0.038	1.168±0.023
0.0795	45.34±0.61	45.28±1.67	52.11±0.97	0.999±0.039	1.149±0.024
0.0850	40.71±0.56	40.04±1.55	45.86±0.90	0.984±0.041	1.127±0.025
0.0931	34.60±0.40	33.50±1.12	38.29±0.65	0.968±0.034	1.106±0.020
0.1040	28.99±0.35	27.45±0.99	30.57±0.58	0.947±0.036	1.054±0.021
0.1150	24.19±0.31	22.92±0.87	24.34±0.51	0.947±0.038	1.006±0.022
0.1310	18.97±0.22	18.73±0.63	18.21±0.36	0.987±0.036	0.960±0.019
0.1530	14.52±0.17	13.72±0.50	12.27±0.28	0.945±0.038	0.845±0.020
0.1750	11.06±0.14	10.18±0.41	8.25±0.22	0.920±0.040	0.746±0.021
0.1970	8.67±0.12	7.53±0.34	5.83±0.18	0.868±0.042	0.673±0.022
0.2189	6.79±0.10	5.76±0.29	4.14±0.15	0.848±0.046	0.609±0.023
0.2410	5.341±0.085	4.381±0.235	2.984±0.128	0.820±0.048	0.559±0.025
0.2629	4.214±0.073	3.358±0.202	2.303±0.110	0.797±0.051	0.547±0.027
0.2849	3.452±0.064	2.487±0.171	1.642±0.094	0.720±0.052	0.475±0.028
0.3068	2.727±0.056	1.947±0.148	1.365±0.085	0.714±0.057	0.501±0.032
0.3338	2.138±0.042	1.436±0.108	0.886±0.063	0.672±0.053	0.414±0.030
0.3666	1.652±0.036	0.817±0.087	0.631±0.052	0.495±0.054	0.382±0.032
0.3997	1.164±0.031	0.614±0.074	0.490±0.047	0.527±0.066	0.421±0.042
0.4325	0.874±0.027	0.386±0.063	0.276±0.038	0.441±0.074	0.316±0.045
0.4651	0.622±0.024	0.429±0.061	0.187±0.033	0.689±0.102	0.301±0.054
0.5035	0.441±0.019	0.206±0.043	0.111±0.025	0.467±0.101	0.252±0.058
0.5470	0.300±0.017	0.142±0.037	0.045±0.019	0.472±0.128	0.151±0.062
0.6083	0.178±0.010	0.066±0.021	0.039±0.010	0.371±0.122	0.222±0.060
0.7047	0.081±0.007	0.003±0.010	0.011±0.005	0.043±0.125	0.141±0.062
0.8383	0.016±0.003	0.003±0.006	0.003±0.002	0.160±0.356	0.190±0.150
Total	15.294±0.250	15.783±0.465	16.841±0.278		

tively; JETSET is within about 2% except for $\xi<1.5$ and $\xi>4.5$; UCLA shows a spectrum very similar to the data but shifted slightly toward lower ξ values; HERWIG is several percent high (low) for $\xi<1$ and $3.3<\xi<4.3$ ($1.5<\xi<2.5$). For the protons we see features similar to those seen in the proton fraction (Fig. 4): JETSET describes the shape of the spectrum but is high by about 20%; UCLA describes the data well except for the structure for $\xi<0.8$; HERWIG is consistent with the data for $\xi>1.5$, but shows a pronounced structure at lower values of ξ that is inconsistent with the data. For kaons, all three predictions are too high for $\xi>3.3$, especially HERWIG, which is also quite high for

TABLE VII. Differential production rates $(1/N_{evts})dn_{K^\pm}/dx_p$ of $K^\pm$ mesons per Z^0 decay into light, c and b primary flavors.

	Kaon production rates			Ratios	
$\langle x_p \rangle$	$u\bar{u}$, $d\bar{d}$, $s\bar{s}$	$c\bar{c}$	$b\bar{b}$	c:uds	b:uds
0.0153	27.05±1.27	30.92±3.86	30.21±1.99	1.143±0.155	1.117±0.090
0.0191	20.00±0.42	22.43±1.37	23.06±0.62	1.121±0.078	1.153±0.040
0.0246	19.74±0.40	22.04±1.27	22.89±0.60	1.116±0.074	1.159±0.039
0.0301	17.52±0.37	20.82±1.17	21.64±0.55	1.189±0.076	1.236±0.041
0.0356	16.08±0.37	16.79±1.15	21.36±0.53	1.044±0.083	1.328±0.046
0.0411	15.04±0.34	16.68±1.15	21.36±0.56	1.110±0.086	1.420±0.048
0.0465	13.54±0.34	16.46±1.06	19.90±0.57	1.215±0.088	1.469±0.056
0.0521	11.87±0.34	15.81±1.08	18.91±0.60	1.332±0.103	1.593±0.068
0.0576	11.44±0.33	12.62±0.99	18.46±0.58	1.103±0.094	1.613±0.069
0.0630	10.64±0.30	12.24±0.92	17.43±0.54	1.151±0.095	1.639±0.069
0.0685	10.24±0.29	11.42±0.87	16.92±0.53	1.115±0.093	1.652±0.070
0.0740	9.67±0.29	10.95±0.85	15.62±0.52	1.133±0.096	1.616±0.072
0.0795	8.13±0.27	10.88±0.84	15.11±0.52	1.339±0.114	1.859±0.090
0.0850	7.98±0.28	9.62±0.81	13.18±0.50	1.206±0.111	1.651±0.084
0.0931	7.00±0.19	9.84±0.59	12.43±0.36	1.406±0.094	1.778±0.070
0.1040	6.36±0.19	8.08±0.58	11.56±0.37	1.271±0.100	1.819±0.080
0.1150	5.85±0.20	8.98±0.63	9.96±0.38	1.535±0.122	1.704±0.088
0.1310	4.89±0.15	6.59±0.45	7.17±0.27	1.349±0.104	1.467±0.072
0.1530	3.41±0.17	5.50±0.51	4.58±0.29	1.614±0.175	1.343±0.108
0.1750	2.84±0.22	5.12±0.68	4.20±0.36	1.805±0.289	1.480±0.173
0.1970	2.564±0.082	3.850±0.245	2.541±0.126	1.502±0.110	0.991±0.059
0.2189	2.401±0.067	3.087±0.190	2.009±0.096	1.286±0.090	0.837±0.047
0.2410	1.973±0.054	2.074±0.145	1.627±0.078	1.051±0.081	0.825±0.046
0.2629	1.643±0.048	1.960±0.132	1.116±0.062	1.193±0.090	0.679±0.043
0.2849	1.481±0.044	1.681±0.119	0.830±0.053	1.135±0.090	0.560±0.039
0.3068	1.211±0.039	1.368±0.104	0.640±0.045	1.129±0.096	0.529±0.041
0.3338	1.001±0.029	1.043±0.076	0.452±0.032	1.042±0.084	0.451±0.034
0.3666	0.746±0.025	0.874±0.068	0.337±0.028	1.171±0.102	0.451±0.040
0.3997	0.666±0.023	0.600±0.058	0.245±0.024	0.900±0.094	0.367±0.038
0.4325	0.559±0.022	0.408±0.050	0.149±0.020	0.730±0.094	0.266±0.036
0.4651	0.426±0.020	0.408±0.050	0.108±0.018	0.957±0.126	0.253±0.043
0.5035	0.363±0.016	0.243±0.037	0.057±0.012	0.669±0.106	0.156±0.035
0.5470	0.261±0.015	0.173±0.034	0.061±0.013	0.663±0.136	0.233±0.050
0.6083	0.183±0.010	0.064±0.020	0.012±0.005	0.351±0.109	0.066±0.030
0.7047	0.079±0.007	0.009±0.011	0.002±0.003	0.113±0.141	0.026±0.035
0.8383	0.008±0.004	0.008±0.008	−.001±0.001	0.989±1.126	−0.098±0.102
Total	1.869±0.062	2.273±0.093	2.377±0.080		

$\xi<0.6$ and low for $1.2<\xi<2.4$; the other two predictions are low for $\xi<0.8$, and consistent with the data elsewhere, although they peak at higher ξ values than the data.

Our fractions and production rates are generally consistent with those from previous experiments at the Z^0 [28,29,33]. The ALEPH pion spectrum lies above ours at large ξ; the OPAL proton spectrum is lower than all others at small ξ; all other differences are within two standard deviations, considering correlations in the systematic errors. Our measurement is the most precise in several regions, most notably for pions in the range $3<\xi<4$, kaons in the range $2.5<\xi<4$ and protons in the range $2.5<\xi<3.5$. Measurements based on ring imaging and those based on ionization energy loss rates cover complementary momentum ranges and could be combined to provide continuous, high-precision coverage over the range from the beam momentum ($\xi=0$) down to $p=0.22$ GeV/c ($\xi=5.3$).

VI. FLAVOR-DEPENDENT ANALYSIS

The analyses described above were repeated on the light-, c-, b-, and untagged event samples described in Sec. III, to yield differential production rates R_h^{ktag}, $k=l,c,b,un$, respectively; the correction for leptons was not performed at this point, so that $h=e\,\mu\,\pi$, K,p. True differential rates R_h^m in events of the three flavor types, $m=l,c,b$, representing

TABLE VIII. Differential production rates $(1/N_{evts})dn_{p/\bar{p}}/dx_p$ of p and $\bar{p}$ per Z^0 decay into light, c and b primary flavors.

	Proton Production Rates			Ratios	
$\langle x_p \rangle$	$u\bar{u}$, $d\bar{d}$, $s\bar{s}$	$c\bar{c}$	$b\bar{b}$	c:uds	b:uds
0.0153	13.98±0.99	13.28±2.94	13.79±1.49	0.950±0.226	0.987±0.127
0.0191	17.63±0.58	15.22±1.89	17.93±0.78	0.863±0.120	1.017±0.056
0.0246	13.42±0.60	13.32±1.86	16.41±0.93	0.992±0.152	1.223±0.089
0.0301	10.57±0.36	9.60±1.16	12.11±0.52	0.909±0.122	1.146±0.063
0.0356	9.98±0.31	11.64±1.01	10.32±0.42	1.166±0.118	1.034±0.052
0.0411	8.37±0.26	10.07±0.87	9.52±0.40	1.203±0.118	1.138±0.059
0.0465	7.33±0.24	8.10±0.76	7.72±0.36	1.105±0.117	1.054±0.059
0.0521	7.79±0.23	6.09±0.72	6.86±0.32	0.781±0.099	0.879±0.049
0.0576	6.62±0.22	6.54±0.67	6.19±0.32	0.988±0.111	0.935±0.057
0.0630	5.88±0.20	6.36±0.63	4.96±0.29	1.082±0.117	0.845±0.056
0.0685	5.39±0.19	4.62±0.59	4.82±0.29	0.857±0.115	0.895±0.062
0.0740	5.22±0.19	4.43±0.58	4.57±0.29	0.848±0.117	0.875±0.064
0.0795	4.42±0.18	4.08±0.55	4.07±0.29	0.924±0.133	0.920±0.076
0.0850	4.44±0.19	3.67±0.55	3.82±0.29	0.827±0.131	0.860±0.074
0.0931	3.65±0.13	4.07±0.42	3.29±0.21	1.116±0.123	0.901±0.066
0.1040	3.11±0.13	2.98±0.41	2.68±0.22	0.960±0.139	0.861±0.079
0.1150	2.73±0.15	2.30±0.43	2.24±0.23	0.844±0.166	0.819±0.096
0.1310	2.15±0.12	2.39±0.36	1.84±0.19	1.110±0.180	0.856±0.100
0.1530	1.83±0.16	1.72±0.50	1.91±0.27	0.940±0.287	1.048±0.174
0.1750	1.84±0.24	0.31±0.71	1.25±0.36	0.171±0.392	0.683±0.219
0.1970	0.905±0.078	0.561±0.235	0.867±0.121	0.619±0.265	0.958±0.157
0.2189	1.065±0.054	0.978±0.163	0.739±0.072	0.918±0.161	0.695±0.077
0.2410	0.822±0.044	0.907±0.136	0.645±0.060	1.104±0.176	0.784±0.085
0.2629	0.762±0.038	0.652±0.116	0.392±0.047	0.855±0.159	0.515±0.067
0.2849	0.628±0.033	0.572±0.101	0.252±0.038	0.911±0.169	0.401±0.064
0.3068	0.486±0.029	0.494±0.089	0.266±0.035	1.016±0.195	0.547±0.079
0.3338	0.446±0.022	0.454±0.069	0.146±0.022	1.016±0.165	0.327±0.052
0.3666	0.306±0.018	0.314±0.054	0.102±0.018	1.026±0.190	0.333±0.061
0.3997	0.230±0.015	0.170±0.041	0.020±0.012	0.741±0.186	0.085±0.053
0.4325	0.197±0.013	0.103±0.033	0.034±0.011	0.522±0.171	0.171±0.055
0.4651	0.145±0.011	0.064±0.028	−0.004±0.007	0.445±0.196	−0.026±0.048
0.5035	0.108±0.008	0.015±0.018	0.016±0.006	0.142±0.166	0.150±0.058
0.5470	0.070±0.006	0.044±0.017	−0.003±0.003	0.626±0.243	−0.043±0.049
0.6083	0.036±0.003	0.007±0.007	0.004±0.002	0.186±0.193	0.107±0.061
0.7047	0.013±0.002	0.015±0.005	−0.001±0.001	1.150±0.434	−0.089±0.045
0.8383	0.003±0.001	0.001±0.002	0.000±0.000	0.221±0.620	0.159±0.160
Total	1.008±0.038	0.930±0.056	0.909±0.037		

events of the types $Z^0 \to u\bar{u}, d\bar{d}, s\bar{s}$, $Z^0 \to c\bar{c}$, and $Z^0 \to b\bar{b}$, respectively, were extracted by solving for each species h the relations

$$R_h^{ktag} = \frac{\Sigma_m B_{mk}^h \epsilon_{mk} F_m R_h^m}{\Sigma_m \epsilon_{mk} F_m}. \quad (1)$$

Here, F_m is the fraction of hadronic Z^0 decays of flavor type m, taken from the standard model, ϵ_{mk} is an element of the event tagging efficiency matrix (see Table I), and B_{mk}^h represents the momentum-dependent bias of tag k toward selecting events of flavor m that contain hadrons of species h. Ideally all biases would be unity in this formulation. The biases were calculated from the simulation as $B_{mk}^h = (n_{m,ktag}^h / N_{m,ktag})/(n_m^h / N_m)$, where N_m (n_m^h) is the number of simulated events (hadrons of species h in events) of true flavor m and $N_{m,ktag}$ ($n_{m,ktag}^h$) is the number of (h hadrons in) those events that are tagged as flavor k. The diagonal bias values [26] are within a few percent of unity, reflecting the small dependences of the flavor tags on the charged multiplicity of the event and tracks from strange mesons and baryons that decay close to the IP. The sum of the products of biases and efficiencies must be unity for events of a given true flavor m, $\Sigma_k B_{mk}^h \epsilon_{mk} = 1$; some off-diagonal bias values therefore deviate substantially from

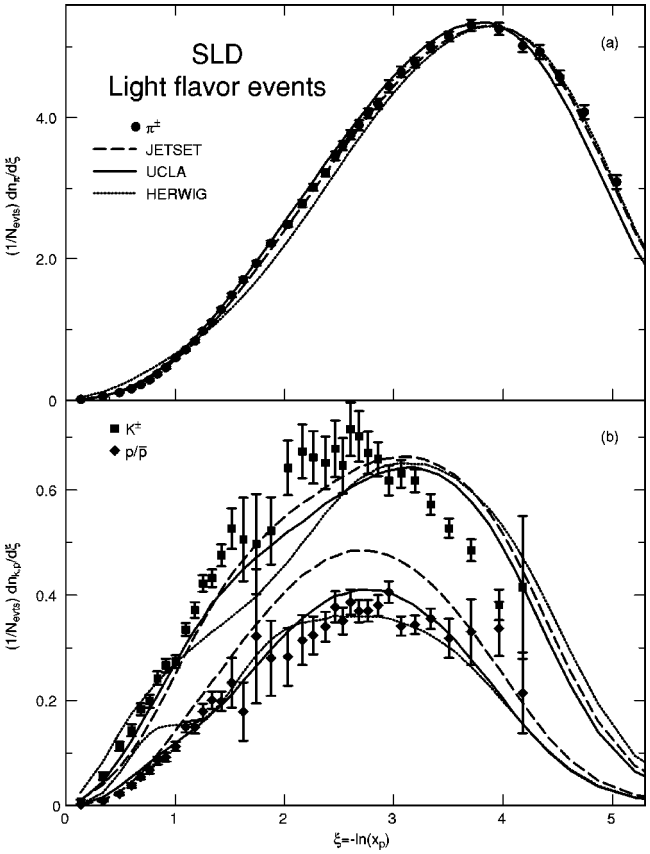

FIG. 6. Distributions of ξ in light-flavor events for (a) pions (circles) and (b) kaons (squares) and protons (diamonds), compared with the predictions of the three models.

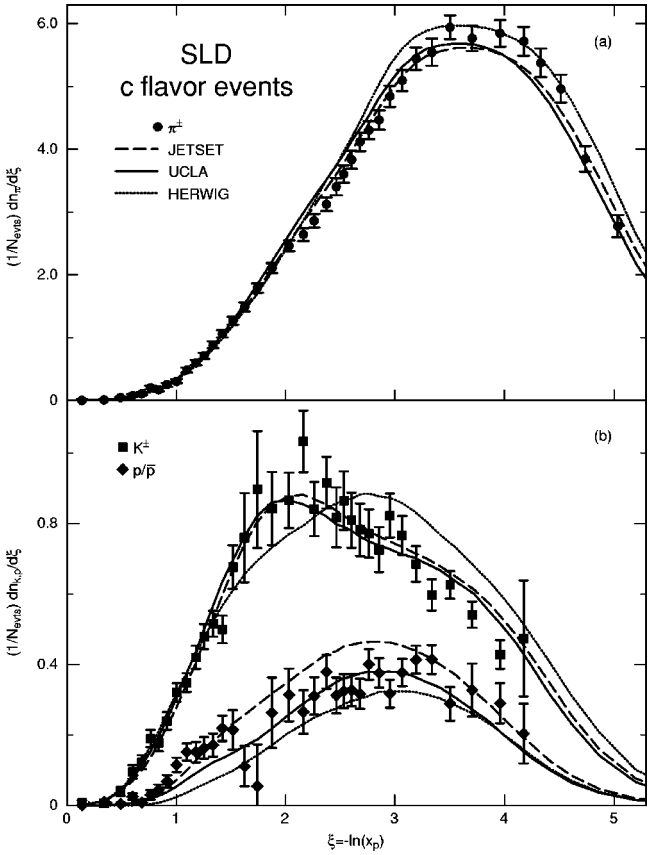

FIG. 7. Distributions of ξ in c-flavor events for (a) pions (circles) and (b) kaons (squares) and protons (diamonds), compared with the predictions of the three models.

unity since the corresponding mistagging rate can be small, but these do not affect the results as the figure of merit is the diagonal element.

The unfolded pion rates were corrected for the contributions from leptons by subtracting the absolute lepton rate predicted by the simulation. At low momentum this background is dominated by electrons from photon conversions, and is a 6% contribution at 1 GeV/c that falls rapidly with increasing momentum. For the heavy flavors, electrons and muons from semi-leptonic decays of heavy hadrons cause the correction to increase with momentum above about 5 GeV/c, reaching 13% (60%) for c- (b-) flavor events in the highest momentum bins.

The resulting differential production rates are listed in Tables VI–VIII. The systematic errors listed are only those relevant for the comparison of different flavors, namely those due to uncertainties in the unfolding procedure; the relative systematic errors given in the preceding section are also applicable, but are common to all three flavor categories; we also list ratios, for which these common errors cancel. The flavor unfolding systematic errors were evaluated by varying each diagonal element of the event tagging efficiency matrix ϵ_{ii} by ± 0.01, the heavy quark production fractions R_b and R_c by the errors on their respective world averages, each diagonal bias value B^h_{ii} by the larger of ± 0.005 and $\pm 20\%$ of its difference from unity, the photon conversion rate in the simulation by $\pm 15\%$, and the simulated physics lepton rates in light-, c-, and b-flavor events by $\pm 20\%$, $\pm 10\%$ and $\pm 5\%$, respectively. These variations correspond to uncertainties on measurements from our data or other experiments except for the variation of the bias, which was chosen conservatively to be larger than any change seen in a set of comparisons of relevant quantities in data and simulation when selection cuts were varied. The unfolding systematic errors are similar in magnitude to (smaller than) the statistical errors at low (high) momenta, and are generally dominated by the bias for the relevant flavor. There are also substantial contributions from the lepton correction at high momentum for R^π_c and R^π_b, and from all three bias terms for R^K_c and R^p_c.

In Figs. 6–8 we show the ξ distributions for the three flavor categories, and in Fig. 9 we show the ratios of production in b-flavor to light-flavor events and c-flavor to light-flavor events vs x_p. At low momentum (high ξ), there is substantially higher production of charged pions in b- and c-flavor events than in light-flavor events, and much higher production of charged kaons in b-flavor events than in light- or c-flavor events; proton production is roughly equal at low momentum. As momentum increases, the production of all three charged hadron species falls much more rapidly in b-flavor events than in light-flavor events, and that in c-flavor events also drops off sharply at very low ξ.

These features are consistent with expectations based on

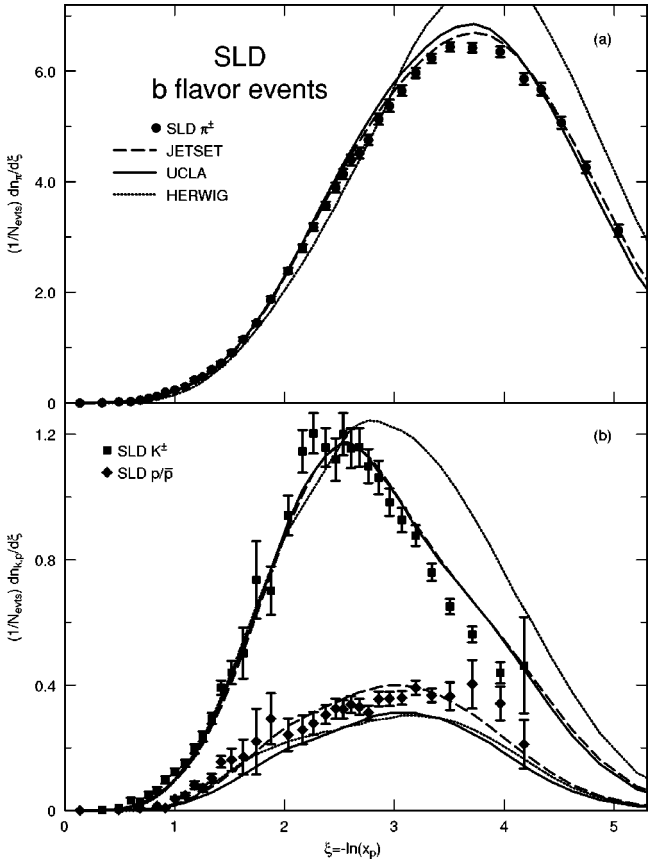

FIG. 8. Distributions of ξ in *b*-flavor events for (a) pions (circles) and (b) kaons (squares) and protons (diamonds), compared with the predictions of the three models.

the known properties of $e^+e^- \to b\bar{b}$ events, namely that a large fraction of the event energy is carried by the leading B and $\bar{B}$ hadrons [34], leaving little energy available to produce nonleading hadrons. The B hadrons decay into a large number of lighter particles, including on average 5.5 stable charged hadrons [8], which are expected to populate primarily the region $2<\xi<4$. The mixture of particle types might be similar to that in light-flavor jets, except that the weak decay chain of the B hadron should produce one ''additional'' $K^{\pm}$ per event, and baryon production might be suppressed since a typical baryon-antibaryon pair mass is not small compared with the B hadron mass. Similar effects in c jets result in an intermediate situation: each jet contains a charmed hadron with on average about half the beam energy, a lower fraction than for B hadrons, which leaves more energy available for nonleading hadrons than in b jets. A D meson decay produces an additional kaon that often carries a large fraction of its momentum, and fewer additional particles than a B decay. Our results are consistent with those published previously [29] and considerably more precise.

Also shown in Figs. 6–9 are the predictions of the three models; they all reproduce the observed flavor dependence qualitatively. In the case of the light flavors, problems with the models very similar to those seen in the preceding section for the flavor inclusive sample are observed, indicating problems in the modelling of hadronization, and not just in that of heavy hadron decay. Some of the discrepancies are larger in the light flavors than in all flavors; in particular the structure at small ξ in the HERWIG model is now very pronounced for all three particle types.

In c-flavor events, all models predict the pion yield within a few percent, but the spectra are slightly too hard. JETSET again predicts the shape of the proton spectrum reasonably well, but is slightly high on the amplitude; UCLA and HERWIG describe the data at large ξ, but fall well below the data at small ξ. JETSET and UCLA predict a kaon spectrum in which those from leading D hadron decays peak at $\xi \approx 2$, and those from the hadronization of the remaining jet form a broad shoulder at larger ξ values. They are qualitatively consistent with the data, although the data prefer a lower shoulder. HERWIG predicts a quite different shape that is inconsistent with the data, probably reflecting known problems in the modelling of heavy hadron production and decay.

In b-flavor events, JETSET and UCLA describe the pion and kaon spectra reasonably well, predicting a few too many pions in the range $2.5<\xi<4$, and too many kaons for $\xi > 3.5$; both also predict the shape of the proton spectrum well, but UCLA is too low in amplitude and JETSET is consistent with the data. Again, HERWIG predicts very different spectra, all of which are inconsistent with the data at large ξ values.

The rate for all charged tracks in each of the flavor samples was derived by summing the rates for the three charged hadron species and adding the simulated contribution from leptons. These rates and their ratios are listed in Table IX; no plots are shown since the features correspond to those of the pion data and models in Figs. 6–9. Our data are consistent with those published by DELPHI [29] and OPAL [10], and substantially more precise.

VII. COMPARISON WITH QCD PREDICTIONS

We tested the predictions of QCD in the modified leading logarithm approximation, combined with the ansatz of local parton-hadron duality, by fitting Gaussian and distorted Gaussian functions to our measured ξ distributions. Examples of such fits are shown in Figs. 10 and 11. In each case, we first fitted a simple Gaussian to each set of points within a region of ξ at least 0.5 units in size centered on the approximate peak position. We chose the largest such symmetric range for which the confidence level of the χ^2 of the fit exceeded 0.01, and then added points on one side (in all cases at lower ξ values) as long as the confidence level remained above 0.01. We found that the Gaussian function could describe the data over a symmetric range of at least 0.8–1.6 units about the peak position, consistent with one prediction of MLLA QCD; the fit could be extended toward lower ξ by as much as 1.7 additional units, although it must be noted that our errors on the kaon and proton spectra are rather large in this range. The largest fittable ranges are given in Table X and the corresponding fitted functions are shown on Figs. 10 and 11.

We next introduced a skewness term into the function (G^+) and repeated the above procedure. We found that the symmetric range could be extended in only some cases, and

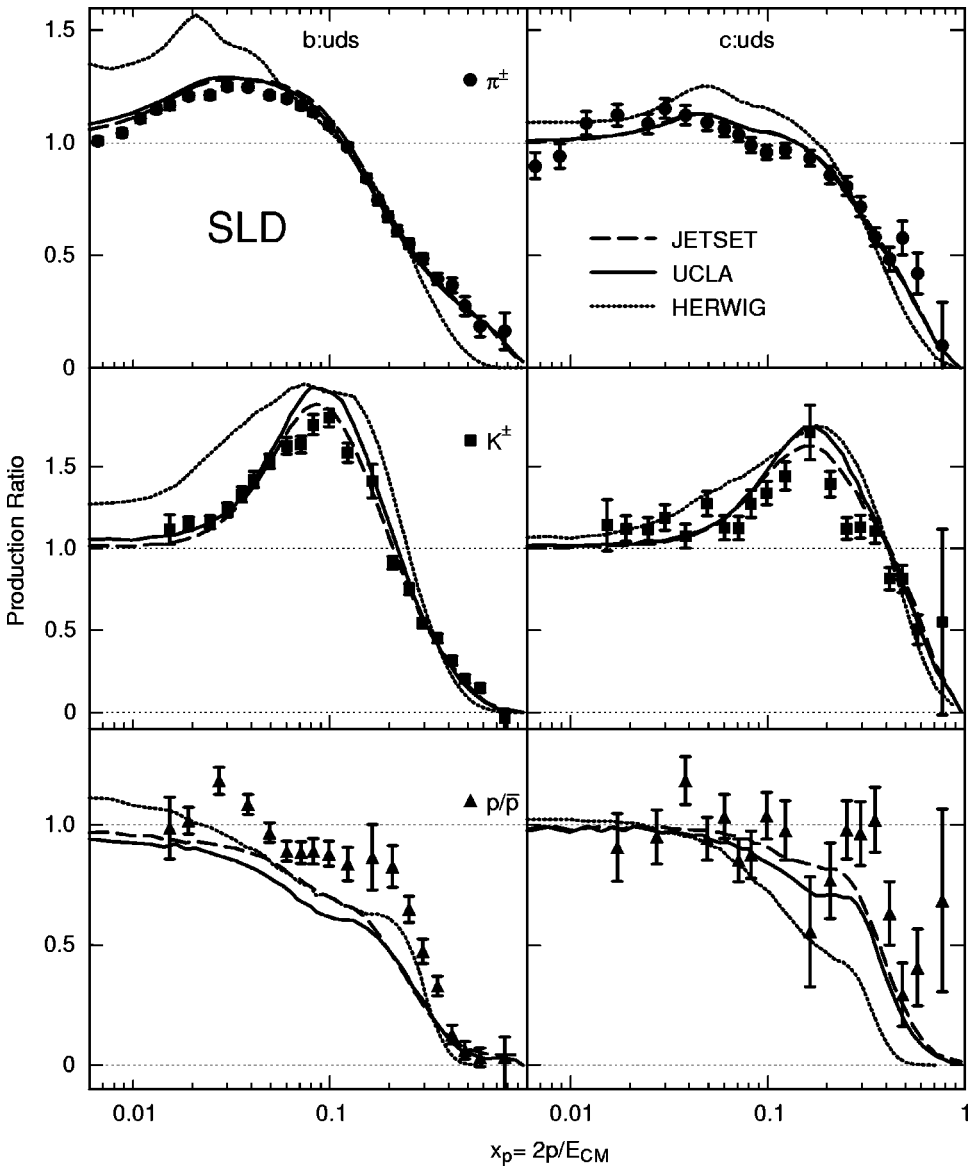

FIG. 9. Ratios of pion (top), kaon (middle) and proton (bottom) production in *b*-flavor events to that in light-flavor events (left) and in *c*-flavor:light-flavor events (right). Some bins have been combined for clarity. Also shown are the predictions of the three models.

by at most 0.6 units (see Table X). In all-flavor and light-flavor events there was always an increase in the fittable range on either the low or high ξ end, and the fitted skewness values were small. In *c*- and *b*-flavor events, however, there were some cases in which no increase in the fittable range could be obtained, and others in which the range could be increased but the skewness value increased rapidly.

The addition of a kurtosis term (G^{++}) had similar results (see Table X); the fittable range was increased greatly, in many cases to the entire measured range; however both the skewness and kurtosis values became large for *c*- and *b*-flavor events. The resulting G^{++} functions are shown on Figs. 10 and 11 over the entire ξ range; the large distortions are evident in the *c*- and *b*-flavor events for kaons and protons. The MLLA prediction that a Gaussian with small distortion terms should describe the data over a range substantially larger than one unit about the peak position holds for the light flavors, but does not hold for heavy flavors; this might be expected since the calculation assumes massless partons.

The peak ξ^* of the ξ distribution is predicted to decrease exponentially with increasing particle mass. Following convention, we took the mean of the fitted Gaussian over a range of one unit about the peak as an estimate of ξ^*, and, in addition to statistical and experimental systematic errors, we considered a variation of the fit range. A fit was performed to each set of contiguous points with a smallest (largest) ξ value between 0.75 and 1.25 units below (above) the peak position. Half the difference between the highest and lowest of the fitted means was taken as an estimate of the systematic uncertainty due to the fit range. The resulting values of ξ^* are listed in Table X, where a considerable flavor dependence is seen.

The values in all-flavor and light-flavor events are shown in Fig. 12 along with previous results [7,12,28,29]. As observed previously, the ξ^* values for the measured hadron species in all-flavor events do not fall on a single trajectory as a function of mass, although roughly parallel trajectories can be postulated for mesons and baryons. For light flavors, the ξ^* value for pions is quite similar to that for all flavors,

TABLE IX. Differential production rates $(1/N_{evts})dn_{chg}/dx_p$ of stable charged particles per Z^0 decay into light, c and b primary flavors.

	Charged particle production rates			Ratios	
$\langle x_p \rangle$	$u\bar{u}$, $d\bar{d}$, $s\bar{s}$	$c\bar{c}$	$b\bar{b}$	c:uds	b:uds
0.0066	507.8±11.9	468.7±25.3	546.1±14.2	0.923±0.053	1.075±0.014
0.0088	505.2±9.2	485.4±22.6	558.5±10.3	0.961±0.050	1.105±0.013
0.0109	465.3±7.5	507.2±20.2	531.9±8.0	1.090±0.052	1.143±0.014
0.0131	421.9±6.3	464.0±17.6	490.8±6.3	1.100±0.050	1.163±0.014
0.0153	371.7±5.8	422.9±16.9	436.5±6.7	1.137±0.054	1.174±0.020
0.0191	315.5±4.2	349.1±12.7	382.8±4.4	1.107±0.049	1.213±0.017
0.0246	250.5±3.2	274.1±9.5	308.6±3.4	1.094±0.046	1.232±0.018
0.0301	200.3±2.3	231.5±7.3	254.8±2.6	1.156±0.044	1.272±0.017
0.0356	167.3±2.0	187.5±6.9	213.2±2.1	1.121±0.050	1.275±0.016
0.0411	140.4±1.6	162.3±5.2	182.1±1.9	1.156±0.044	1.297±0.017
0.0465	121.2±1.3	136.6±4.2	154.5±1.7	1.127±0.041	1.275±0.018
0.0521	105.5±1.1	117.3±3.6	134.3±1.6	1.112±0.040	1.273±0.019
0.0576	91.2±1.0	99.1±3.0	118.6±1.4	1.087±0.038	1.301±0.020
0.0630	81.29±0.89	89.21±2.65	102.4±1.3	1.097±0.037	1.260±0.020
0.0685	72.69±0.81	78.25±2.42	91.92±1.22	1.077±0.037	1.265±0.021
0.0740	65.92±0.76	69.26±2.23	83.63±1.14	1.051±0.038	1.269±0.022
0.0795	58.06±0.70	62.25±2.06	75.06±1.10	1.072±0.040	1.293±0.024
0.0850	53.26±0.66	55.28±1.93	66.58±1.04	1.038±0.040	1.250±0.024
0.0931	45.37±0.48	49.14±1.42	57.31±0.74	1.083±0.035	1.263±0.020
0.1040	38.55±0.43	40.11±1.29	47.80±0.71	1.040±0.037	1.240±0.022
0.1150	32.84±0.41	35.64±1.22	39.19±0.67	1.085±0.041	1.194±0.025
0.1310	26.05±0.28	28.94±0.83	29.54±0.48	1.111±0.035	1.134±0.022
0.1530	19.79±0.29	21.99±0.86	20.69±0.48	1.111±0.048	1.045±0.029
0.1750	15.75±0.35	16.51±1.03	15.36±0.55	1.048±0.071	0.975±0.041
0.1970	12.16±0.17	12.69±0.50	10.65±0.26	1.044±0.046	0.876±0.024
0.2189	10.27±0.14	10.41±0.40	8.06±0.20	1.013±0.044	0.785±0.022
0.2410	8.14±0.11	7.86±0.32	6.28±0.16	0.965±0.044	0.772±0.022
0.2629	6.62±0.10	6.37±0.28	4.69±0.13	0.961±0.047	0.709±0.023
0.2849	5.565±0.087	5.060±0.241	3.490±0.110	0.909±0.048	0.627±0.022
0.3068	4.428±0.076	4.080±0.210	2.935±0.098	0.922±0.052	0.663±0.025
0.3338	3.588±0.057	3.123±0.156	2.041±0.068	0.870±0.048	0.569±0.021
0.3666	2.706±0.049	2.141±0.126	1.534±0.058	0.791±0.051	0.567±0.024
0.3997	2.062±0.042	1.472±0.106	1.111±0.051	0.714±0.055	0.539±0.027
0.4325	1.631±0.037	0.952±0.090	0.736±0.042	0.584±0.058	0.451±0.028
0.4651	1.193±0.034	0.935±0.085	0.510±0.037	0.783±0.076	0.427±0.033
0.5035	0.912±0.026	0.485±0.061	0.330±0.027	0.532±0.070	0.362±0.032
0.5470	0.632±0.023	0.372±0.055	0.188±0.022	0.588±0.091	0.298±0.037
0.6083	0.398±0.015	0.140±0.031	0.089±0.012	0.353±0.081	0.224±0.031
0.7047	0.172±0.011	0.027±0.016	0.017±0.006	0.158±0.096	0.099±0.034
0.8383	0.027±0.005	0.011±0.010	0.003±0.002	0.412±0.386	0.103±0.095

and those for the other mesons (baryons) tend to be higher (lower), as would be expected if the MLLA prediction is correct for primary hadrons, but the observed distributions are distorted by secondary particles. However it is still not possible to draw a single trajectory through all the light-flavor points and this apparent deficiency in the theory remains; it is possible that the exclusion of additional non-primary hadrons might remove it.

Our previous results [12] were used in conjunction with data from other experiments to test the perturbative QCD calculations of the E_{CM} dependence of the rates at high x_p [5]. The results presented here can be used to make a slightly more precise test; however more precise data at lower and/or higher E_{CM} are needed to confront these calculations in a more meaningful way.

VIII. TOTAL PRODUCTION RATES

We have integrated our differential production rates over their respective measurement ranges, taking into account the bin-to-bin correlations in the systematic errors. These integrated rates per event are listed in Tables II–V and VI–VIII;

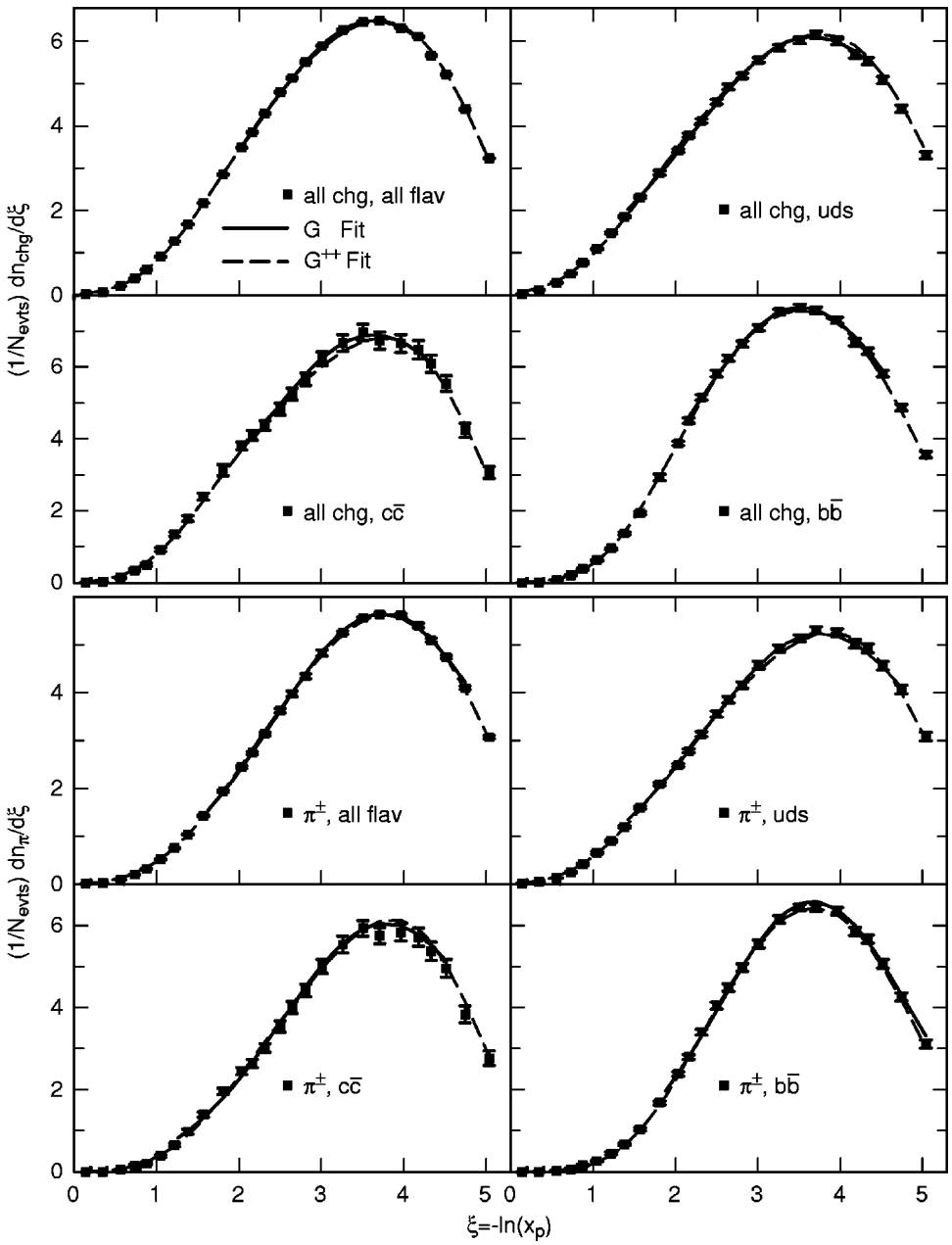

FIG. 10. Measured ξ distributions for all charged particles (top) and pions (bottom) in all-, light-, c- and b-flavor events. Some bins have been combined for clarity. The solid (dashed) lines represent the results of the maximal Gaussian (G) and G^{++} fits described in the text.

the errors are dominated by overall normalization uncertainties corresponding to the uncertainty in our track reconstruction efficiency. In order to quote total rates, we must extrapolate into the unmeasured regions, $\xi > 5.21$ for inclusive charged particles and pions, and $\xi > 4.25$ for kaons and protons. As can be seen in Figs. 5–8, 10, and 11, this is a non-trivial effect, and none of the models or functions provides an ideal estimate of the accepted fraction. However, a set of four estimates derived from the three models and the G^{++} function might be considered to cover a reasonable range of possibilities. For the c- and b-flavor events, we used in addition the generator used for our detector simulation [27] a version of JETSET tuned to the world's data on D and B hadron production and decay.

From the pions and protons generated using each of the models, we calculated the fraction that were generated in the range of our measurement, as the predictions, when normalized to the data, provided reasonable descriptions of the shape. For the charged kaons (Figs. 5–8), all models predict spectra that are too soft; we therefore scaled the predictions along the ξ axis until the best agreement with the data was found. This procedure changed the acceptance correction by a few percent for b- and c-flavor events, and by as much as 12% for the HERWIG model in light-flavor events. The fitted G^{++} function was integrated to obtain another estimate of the accepted fraction in those cases where it gave an acceptable χ^2 over the entire measured range. Otherwise, it was used to calculate a fraction above the lower bound given in Table X, and the data below that bound were added to obtain an overall fraction.

For each hadron species and flavor category these four or five estimates of the fraction were found to be similar, with a typical rms value of about 1% relative. We took their average as our central value of the fraction in each case, and took

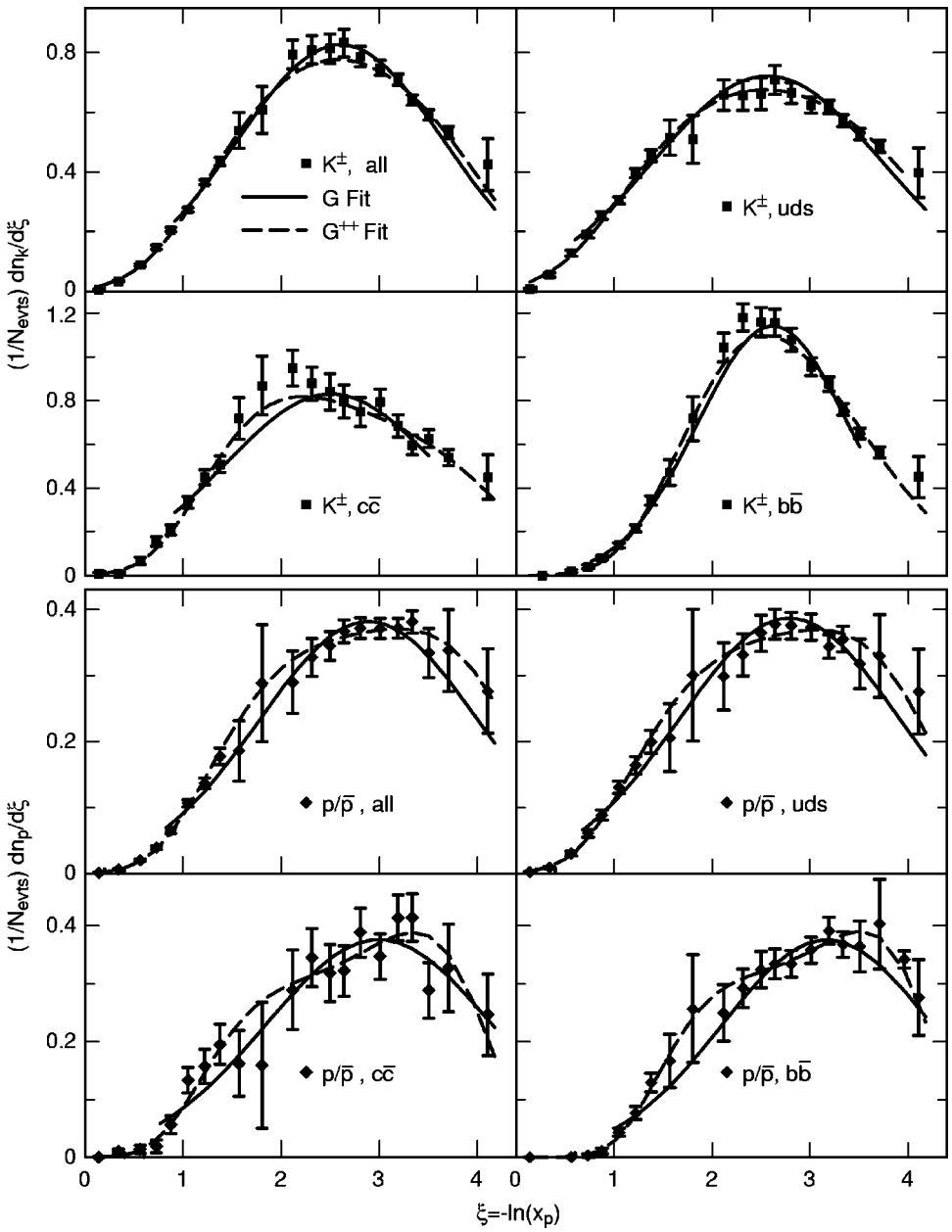

FIG. 11. Measured ξ distributions for charged kaons (top) and protons (bottom) in all-, light-, *c*- and *b*-flavor events. Some bins have been combined for clarity. The solid (dashed) lines represent the results of the maximal Gaussian (G) and G^{++} fits described in the text.

their rms as the uncertainty due to the extrapolation procedure. These values along with the corrected total rates are listed in Table XI; also shown are differences between the three flavor categories, for which some of the uncertainties cancel. The results were found to be consistent with previous measurements; the precision is similar for all-flavor events, and substantially improved for light and *b* flavors; ours are the only measurements for *c*-flavor events.

Roughly 10% more pions are produced in *b*-flavor events than in *c*- or light-flavor events; roughly 20% more kaons are produced in both *b*- and *c*-flavor events than in light-flavor events; roughly 10% fewer protons are produced in *b*- than in light-flavor events. The total charged multiplicities and differences between flavors are consistent with previous, dedicated measurements [9,10]; they have comparable precision with different systematic error sources.

IX. LEADING PARTICLE EFFECTS

We extended these studies to look for differences between particle and antiparticle production in light quark (as opposed to antiquark) jets, in order to address the question of whether e.g. a primary *u*-initiated jet contains more hadrons that contain a valence *u* quark (e.g. π^+, K^+, p) than hadrons that do not (e.g. $\pi^-, K^-, \bar{p}$). To this end we used the light quark- and antiquark-tagged hemispheres described in Sec. III.

We measured the differential production rates per light quark jet

$$R_h^q = \frac{1}{2N_{evts}} \frac{d}{dx_p} [N(q \to h) + N(\bar{q} \to \bar{h})], \qquad (2)$$

TABLE X. The widest ranges in ξ (see text) over which a Gaussian function (G) was able to describe the data alone, and with the addition of skewness (G^+) and kurtosis terms (G^{++}). Peak positions ξ^* from the Gaussian fits described in the text; the errors are statistical, experimental systematic, and due to variation of the fit range.

	Maximum fit range			$\xi^* \pm \text{stat} \pm \text{syst} \pm \text{fit}$
	G	G^+	G^{++}	
			$\pi^\pm$	
All	1.62–4.74	1.88–5.21	0.92–5.21	$3.752 \pm 0.004 \pm 0.008 \pm 0.010$
uds	1.40–4.85	1.55–5.21	0.64–5.21	$3.787 \pm 0.008 \pm 0.027 \pm 0.014$
c	1.30–4.60	1.47–5.21	0.87–5.21	$3.726 \pm 0.014 \pm 0.090 \pm 0.039$
b	1.81–5.21	1.81–5.21	0.00–5.21	$3.684 \pm 0.008 \pm 0.005 \pm 0.001$
			$K^\pm$	
All	0.92–4.18	0.49–4.18	0.34–4.18	$2.549 \pm 0.018 \pm 0.042 \pm 0.062$
uds	0.55–4.25	0.24–4.25	0.24–4.25	$2.592 \pm 0.032 \pm 0.091 \pm 0.021$
c	0.87–3.61	0.00–4.25	0.00–4.25	$2.412 \pm 0.030 \pm 0.018 \pm 0.012$
b	0.72–3.61	0.00–4.25	0.00–4.25	$2.587 \pm 0.014 \pm 0.017 \pm 0.004$
			$p\bar{p}$	
All	0.84–4.18	0.76–4.18	0.14–4.18	$3.084 \pm 0.101 \pm 0.056 \pm 0.031$
uds	0.64–4.25	0.42–4.25	0.00–4.25	$2.859 \pm 0.076 \pm 0.028 \pm 0.026$
c	0.72–4.25	0.72–4.25	0.00–4.25	$3.079 \pm 0.113 \pm 0.023 \pm 0.068$
b	0.97–4.25	0.97–4.25	0.00–4.25	$3.513 \pm 0.100 \pm 0.254 \pm 0.011$

$$R_h^q = \frac{1}{2N_{evts}} \frac{d}{dx_p} [N(q \to \bar{h}) + N(\bar{q} \to h)], \qquad (3)$$

where q and $\bar{q}$ represent light-flavor quark and antiquark jets respectively; N_{evts} is the total number of events in the sample; h represents π^+, K^+ or p, and $\bar{h}$ the corresponding antihadron. Then, for example, $N(q \to h)$ is the number of hadrons of species h in light quark jets. This formulation assumes *CP* symmetry, i.e. $N(q \to h) = N(\bar{q} \to \bar{h})$, which was found to be satisfied in the data in all cases.

The charged hadron analysis was repeated on the sample of positively charged tracks in the quark-tagged jets and negatively charged tracks in the antiquark-tagged jets, yielding measured values of $R^q_{\pi^+}$, $R^q_{K^+}$, and R^q_p in the tagged samples. The same procedure applied to the remaining tracks yielded $R^q_{\pi^-}$, $R^q_{K^-}$, and $R^{\bar{q}}_p$.

The decays of the leading heavy hadrons in simulated heavy flavor background events give rise to differences between hadron and antihadron production in the quark-tagged sample over the entire x_p range, including an effect in high-momentum pions of opposite sign to that seen in the light-

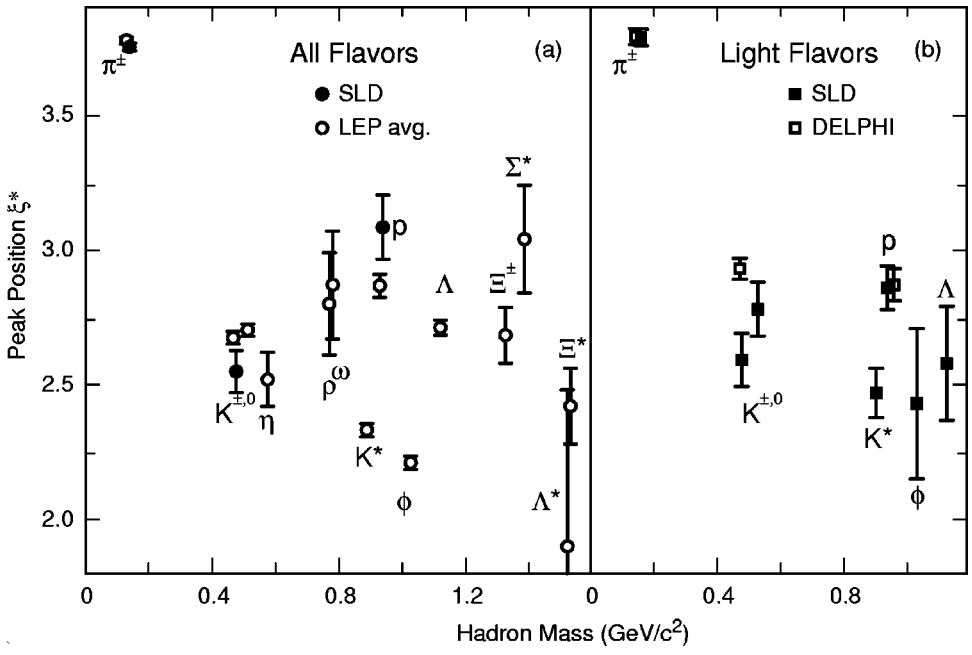

FIG. 12. Peak positions ξ^* as a function of hadron mass in (a) all- and (b) light-flavor events, along with previous results from experiments at the Z^0.

TABLE XI. Estimated fractions of the total production rates contained within the range of the measurements; corrected yields of charged hadrons per Z^0 decay into events of each flavor category. The rightmost columns show differences between flavors, for which some uncertainties cancel.

		Coverage	Yield/event		Difference
	All	0.925 ± 0.007	17.007 ± 0.209		
	uds	0.923 ± 0.008	16.597 ± 0.304	$c-uds$	0.375 ± 0.587
$\pi^{\pm}$	*c*	0.931 ± 0.013	16.954 ± 0.556	$b-uds$	1.557 ± 0.232
	b	0.929 ± 0.007	18.136 ± 0.330	$b-c$	1.182 ± 0.498
	All	0.941 ± 0.010	2.203 ± 0.071		
	uds	0.934 ± 0.007	2.000 ± 0.068	$c-uds$	0.427 ± 0.074
$K^{\pm}$	*c*	0.937 ± 0.005	2.427 ± 0.100	$b-uds$	0.510 ± 0.037
	b	0.947 ± 0.006	2.510 ± 0.086	$b-c$	0.083 ± 0.075
	All	0.933 ± 0.003	1.054 ± 0.035		
	uds	0.921 ± 0.010	1.094 ± 0.043	$c-uds$	-0.060 ± 0.074
$p\bar{p}$	*c*	0.899 ± 0.038	1.034 ± 0.077	$b-uds$	-0.091 ± 0.034
	b	0.906 ± 0.020	1.004 ± 0.046	$b-c$	-0.031 ± 0.074
	uds	0.927 ± 0.007	20.048 ± 0.316	$c-uds$	1.048 ± 0.718
All charged	*c*	0.940 ± 0.015	21.096 ± 0.653	$b-uds$	3.050 ± 0.311
	b	0.933 ± 0.010	23.098 ± 0.378	$b-c$	2.002 ± 0.643

flavor data (below). It is essential to understand and/or suppress this contribution; our simulation has been tuned to the available data on charmed and bottom hadron decays, and our light-flavor tag reduces this background to 9% $c\bar{c}$ and 2% $b\bar{b}$ events. At this level the simulated contribution to any difference is well below the other uncertainties. The simulated contribution to each rate was applied as a correction, yielding differential production rates per light-quark-tagged jet. For each hadron species, differential production rates in light quark jets were then extracted by correcting for the light-tag bias (see Sec. VI) and unfolding for the effective quark (vs antiquark) purity. This purity depends on polar angle, and hence on the acceptance of the CRID as well as the event selection discussed in Sec. III; its average value was estimated from the simulation to be 0.72 for the selected track sample.

The measured differential production rates per light quark jet are listed in Table XII and shown in Fig. 13; as for the flavor dependent results (Sec. VI), the errors given are the sums in quadrature of the statistical error and those systematic errors arising from the tagging and correction procedures. The latter include variation of the event tagging efficiencies and biases as described in Sec. VI; the electroweak parameters R_b, R_c, A_b and A_c by the errors on their respective world average values; the effective quark purity by ± 0.01; the sum of h and $\bar{h}$ rates in *c*- and *b*-flavor events by a smooth parametrization of the errors in Tables VI–VIII; and their difference by $\pm 20\%$ of itself to cover the uncertainty in the electron beam polarization and statistical error on the simulated purities. The systematic errors are small compared with the statistical errors, and are typically dominated by the uncertainty on the effective quark purity.

In all cases the hadron and antihadron rates are consistent at low x_p. For kaons (protons), significant differences appear for x_p above about 0.15 (0.2) units and grow with increasing x_p. For pions, differences are smaller, but significant for x_p above about 0.3, also appearing to grow with x_p. It is convenient to show these results in the form of the difference between hadron h and antihadron $\bar{h}$ production normalized by the sum

$$D_h = \frac{R_h^q - R_h^{\bar{q}}}{R_h^q + R_h^{\bar{q}}}. \tag{4}$$

The common systematic errors cancel explicitly in this variable, which is shown for the hadrons $h = \pi^-, K^-, p$ in Fig. 14 and listed in Table XII. A value of zero corresponds to equal production of hadron and antihadron, whereas a value of $+$ $(-)$ 1 corresponds to complete dominance of (anti)particle production in light quark jets.

The results for the protons afford the most straightforward interpretation. Since baryons contain valence quarks and not antiquarks, the positive values of D_p for $x_p > 0.2$ are clear evidence for the production of leading protons. The data are consistent with a steady increase with x_p to a plateau of $D_p \approx 0.5$ for $x_p > 0.5$, although the errors on the highest x_p points are quite large. For $x_p < 0.1$ the data are consistent, within common systematic errors, with equal production of baryons and antibaryons; however it must be noted that the contribution from non-leading hadrons is very high in this region and we cannot exclude the possibility that some leading baryons are produced at low x_p.

The interpretation of our results for mesons is more complicated, since a meson contains one valence quark along with one valence antiquark. In standard model (SM) Z^0 decays all primary down-type quarks are produced equally and with the same forward-backward asymmetry, so that if a

TABLE XII. Differential production rates $R_h^q=(1/2N_{evts})dn_h/dx_p$ for positively and negatively charged hadrons $h=\pi^\pm,K^\pm,p,\bar{p}$ in light (u, d, and s) quark jets from hadronic Z^0 decays, along with the normalized differences $D_h=(R_h^q-R_{\bar{h}}^q)/(R_h^q+R_{\bar{h}}^q)$. The errors are the sum in quadrature of statistical errors and those systematic errors arising from the light quark tagging and unfolding procedure.

$\langle x_p \rangle$	R_{π^+}	R_{π^-}	D_{π^-} (%)	R_{K^+}	R_{K^-}	D_{K^-} (%)	R_p	$R_{\bar{p}}$	D_p (%)
0.0066	243.9±4.0	237.5±4.0	−1.3±1.5	—	—	—	—	—	—
0.0088	236.7±3.6	235.5±3.5	−0.3±1.3	—	—	—	—	—	—
0.0110	210.5±3.1	215.3±3.2	1.1±1.3	—	—	—	—	—	—
0.0131	194.5±2.9	189.4±2.8	−1.3±1.3	—	—	—	—	—	—
0.0153	168.1±2.6	166.4±2.6	−0.5±1.4	13.69±1.39	12.89±1.39	−3.0±9.4	8.09±1.19	6.67±1.18	9.6±14.3
0.0191	137.8±1.5	142.8±1.6	1.8±1.0	9.23±0.45	10.10±0.45	4.5±4.1	8.98±0.62	8.82±0.62	0.9±5.9
0.0246	108.1±1.4	109.5±1.4	0.7±1.1	10.10±0.44	9.73±0.44	−1.9±4.0	7.42±0.72	6.41±0.72	7.4±9.3
0.0301	84.99±1.18	88.51±1.19	2.0±1.2	8.67±0.40	8.87±0.40	1.1±4.1	5.81±0.42	5.05±0.42	7.0±6.6
0.0356	68.97±1.06	71.82±1.07	2.0±1.3	8.55±0.38	7.80±0.38	−4.6±4.2	5.52±0.34	5.19±0.34	3.1±5.6
0.0411	57.26±0.96	59.36±0.97	1.8±1.5	7.91±0.38	8.08±0.38	1.1±4.3	4.40±0.29	4.18±0.28	2.6±6.0
0.0466	49.27±0.91	49.63±0.91	0.4±1.6	7.81±0.40	7.37±0.40	2.9±4.7	3.74±0.26	3.40±0.26	4.8±6.6
0.0521	43.81±0.86	42.11±0.86	−2.0±1.8	5.27±0.40	7.04±0.41	14.3±5.9	4.11±0.27	3.74±0.27	4.7±6.3
0.0576	37.47±0.79	36.45±0.78	−1.4±1.9	5.46±0.37	5.59±0.38	1.2±6.2	3.48±0.25	3.32±0.25	2.4±6.6
0.0631	32.00±0.72	32.79±0.72	1.2±2.0	4.91±0.35	5.69±0.35	7.4±6.0	3.01±0.23	3.01±0.24	−0.1±7.7
0.0685	28.67±0.68	28.02±0.67	−1.1±2.1	5.24±0.34	4.96±0.34	−2.9±6.0	2.87±0.23	2.47±0.23	7.6±7.7
0.0740	24.88±0.63	26.16±0.63	2.5±2.2	4.65±0.33	5.14±0.34	5.1±6.2	2.58±0.23	2.58±0.22	0.1±7.9
0.0795	21.33±0.59	24.09±0.60	6.1±2.4	4.09±0.32	4.31±0.33	2.6±7.0	2.34±0.22	1.94±0.21	9.2±9.1
0.0850	20.51±0.56	20.15±0.56	−0.9±2.5	4.31±0.32	3.84±0.33	−5.9±7.3	1.87±0.22	2.50±0.23	−14.5±9.4
0.0931	17.82±0.37	17.17±0.36	−1.8±1.9	3.13±0.22	3.89±0.23	10.9±5.8	1.94±0.16	1.87±0.16	2.1±7.5
0.1041	14.52±0.33	14.61±0.33	0.3±2.0	3.07±0.22	3.28±0.23	3.3±6.4	1.43±0.16	1.62±0.16	−6.1±9.5
0.1150	12.25±0.30	11.99±0.30	−1.1±2.2	3.00±0.24	3.05±0.25	0.8±7.3	1.45±0.18	1.33±0.17	4±11
0.1311	9.454±0.188	9.663±0.189	1.1±1.8	2.02±0.18	2.91±0.19	17.9±6.8	1.30±0.15	0.96±0.14	15±12
0.1530	7.231±0.163	7.386±0.164	1.1±2.0	1.50±0.21	1.77±0.22	8±12	0.88±0.19	1.07±0.20	−10±18
0.1750	5.594±0.141	5.433±0.139	−1.5±2.3	1.09±0.28	1.77±0.30	24±19	1.00±0.27	0.82±0.28	10±28
0.1970	4.330±0.123	4.360±0.122	0.4±2.6	0.997±0.104	1.528±0.111	21.0±6.9	0.501±0.087	0.417±0.087	9±18
0.2189	3.302±0.108	3.503±0.108	3.0±2.9	0.856±0.078	1.593±0.086	30.1±5.9	0.612±0.064	0.374±0.063	24±14
0.2410	2.619±0.092	2.721±0.093	1.9±3.1	0.784±0.061	1.180±0.067	20.1±5.9	0.469±0.050	0.384±0.050	10±11
0.2629	2.059±0.082	2.142±0.083	2.0±3.6	0.571±0.054	1.105±0.060	31.9±6.2	0.475±0.045	0.275±0.043	27±11
0.2850	1.623±0.074	1.844±0.075	6.4±3.9	0.508±0.049	0.959±0.055	30.7±6.4	0.460±0.040	0.191±0.038	41±11
0.3070	1.294±0.065	1.378±0.065	3.2±4.4	0.366±0.043	0.824±0.049	38.4±7.0	0.321±0.033	0.174±0.032	30±12
0.3337	0.974±0.048	1.161±0.049	8.8±4.1	0.320±0.033	0.711±0.038	27.9±6.2	0.298±0.025	0.147±0.024	34±10
0.3665	0.721±0.042	0.908±0.044	11.5±4.8	0.249±0.029	0.529±0.033	35.9±7.1	0.198±0.020	0.123±0.020	23±11
0.3996	0.523±0.036	0.637±0.037	9.9±5.7	0.194±0.027	0.503±0.031	44.4±7.4	0.163±0.017	0.095±0.017	26±12
0.4326	0.424±0.032	0.451±0.032	3.1±6.7	0.167±0.024	0.372±0.027	37.9±8.5	0.114±0.015	0.083±0.014	16±13
0.4653	0.278±0.028	0.375±0.030	15.0±8.1	0.132±0.022	0.294±0.025	38.1±9.9	0.105±0.013	0.036±0.012	49±16
0.5034	0.221±0.021	0.222±0.022	0.2±8.8	0.089±0.018	0.267±0.021	50.0±9.5	0.085±0.010	0.022±0.009	59±15
0.5464	0.107±0.018	0.192±0.020	28±12	0.050±0.016	0.208±0.019	61±12	0.057±0.008	0.011±0.007	67±20
0.6085	0.066±0.011	0.112±0.012	26±12	0.063±0.011	0.118±0.012	30±11	0.024±0.004	0.009±0.004	48±21
0.7046	0.020±0.007	0.054±0.009	46±20	0.022±0.008	0.056±0.009	43±20	0.006±0.002	0.009±0.002	−18±29
0.8342	0.011±0.004	0.006±0.004	−27±40	0.005±0.005	0.006±0.005	9±83	0.000±0.001	0.003±0.001	−164±144

leading neutral hadron such as $\bar{K}^{*0}$ ($s\bar{d}$) were produced equally in s and $\bar{d}$ jets then we would measure $D_{\bar{K}^{*0}}$ to be zero. We previously measured [12] a significantly positive value (see Fig. 14) at high x_p, indicating both that there is leading $\bar{K}^{*0}$ production *and* that more leading $\bar{K}^{*0}$ are produced in s jets than in $\bar{d}$ jets. This is an expected consequence of strangeness suppression in the hadronization process. That is, it is expected to be less likely for an $s\bar{s}$ pair to pop out of the vacuum and the s to pair up with an initial $\bar{d}$ than it is for a $d\bar{d}$ pair to pop out and the $\bar{d}$ to pair up with the initial s.

In the case of charged mesons such as π^- ($d\bar{u}$), there is a nonzero dilution of leading particle effects because of the different Z^0 branching ratios and forward-backward asymmetries of up- and down-type quarks. If we assume SM couplings to the Z^0, equal production of leading π^+ in u jets and

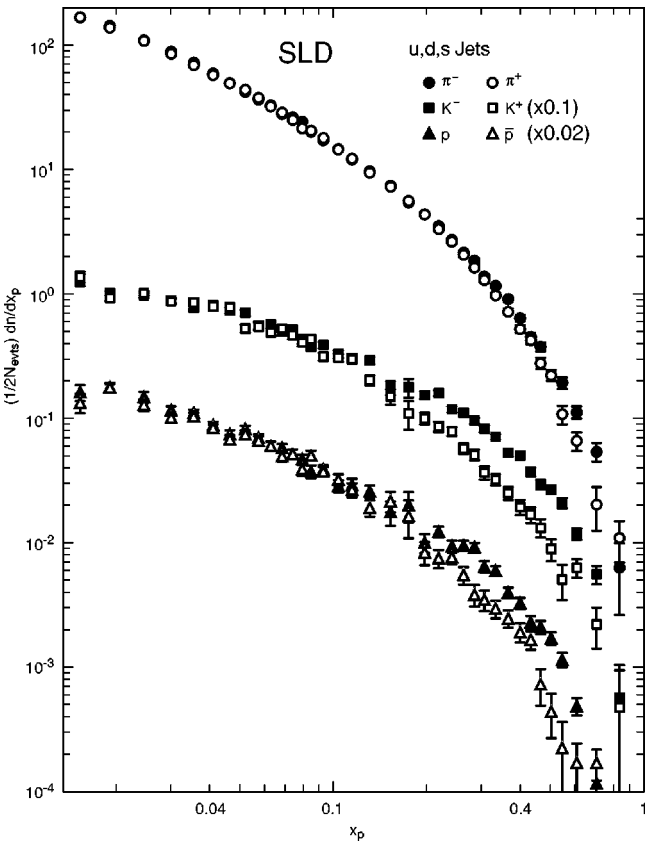

FIG. 13. Fully corrected production rates for positively and negatively charged hadrons in light quark jets. The error bars include statistics and those systematic uncertainties that affect the difference between the charges.

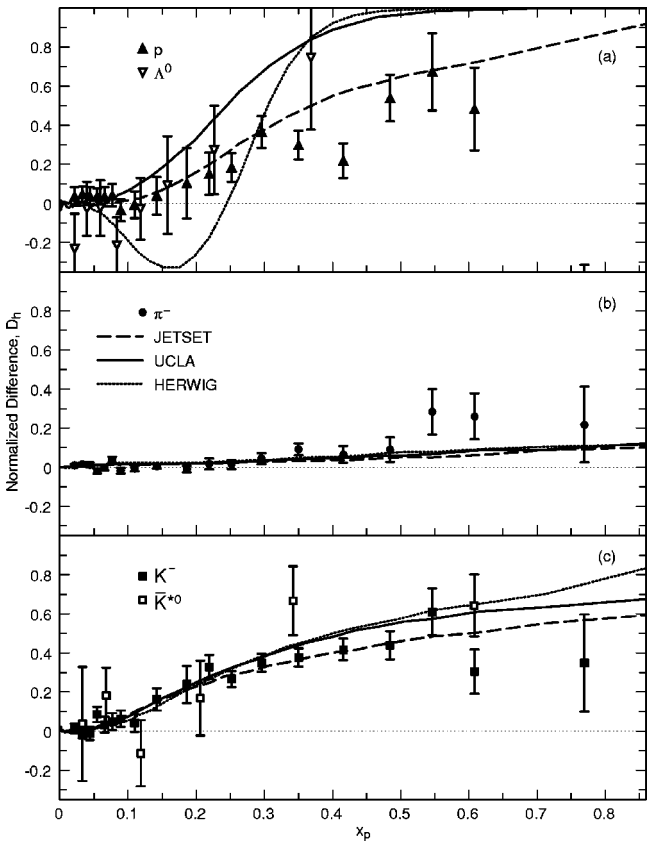

FIG. 14. Normalized differences between hadron and antihadron production in light quark jets, compared with the predictions of the three models. Some bins have been combined for clarity. Our previous results [12] for Λ^0 (K^{*0}) are overlaid on the proton (kaon) plot; the corresponding model predictions are similar.

π^- in d jets, and no contribution from s jets, then we calculate a dilution factor for this analysis of 0.27. That is, we would expect to observe a value of $D_{\pi^-}^{obs}=0.27D_{d\to\pi^-}^{true}=0.27D_{u\to\pi^+}^{true}$.

Our measured D_{π^-} are significantly positive at high x_p, indicating leading pion production. The data give no information on the relative contributions of the three light flavors, but are consistent with $0.27D_p$ at all x_p, and hence with the notion that leading pion effects in u and d jets are of similar strength to those of leading baryons in all light quark jets. Our measured D_{K^-} are consistently positive for $x_p>0.15$, and well above both $0.27D_p$ and D_{π^-} at high x_p. This indicates both production of leading charged kaons and greater production of leading K^- in s jets than in $\bar{u}$ jets.

The predictions of the three models are also shown in Fig. 14. For protons, the HERWIG prediction drops below zero for $0.05<x_p<0.25$, then rises rapidly to unity at higher x_p; this structure is inconsistent with our data. The UCLA prediction also rises to unity at high x_p, becoming inconsistent with the data for $x_p>0.25$. The JETSET prediction is consistent with the data. For the mesons, all models predict positive values of D_{π^-} and D_{K^-} even at very low x_p; they are all consistent with our data in that region, but we are not able to resolve the difference from zero. At high x_p, all three predictions for D_{π^-} lie roughly 2σ below the data, and the measured D_{K^-} favor the JETSET prediction slightly over those of HERWIG and UCLA.

X. SUMMARY AND CONCLUSIONS

We have improved our measurements of the production of the charged hadron species $\pi^\pm$, $K^\pm$ and $p/\bar{p}$, as well as that of inclusive stable charged particles, in hadronic Z^0 decays, taking advantage of the sample of 400 000 decays recorded with the upgraded vertex detector to reduce both the statistical and systematic errors substantially. The SLD Cherenkov ring imaging detector enabled the clean and efficient identification of charged tracks over a wide momentum range, yielding precise measurements of their total and differential production rates. Our measurements in flavor-inclusive hadronic Z^0 decays are consistent with, complementary to, and in some regions more precise than previous measurements. Deficiencies in popular hadronization models, using their respective default parameter values, have been confirmed.

The precision of the vertex detector allowed us to isolate very high-purity light-, c- and b-tagged event samples, and to compare production characteristics of the hadron species in light-, c- and b-flavor events. Significant differences between flavors were found, consistent with expectations based on the known properties of B and D hadron production and decay.

The known problems with hadronization models were all present in the light-flavor events, confirming that they are indeed in the hadronization itself and not just artifacts of heavy hadron modelling. Additional problems with the models were observed in *c*- and *b*-flavor events.

The shape of the $\xi = -\ln x_p$ distribution near its peak for each hadron species in events of each flavor is consistent with the Gaussian form predicted by MLLA QCD+LPHD. For *c*- and *b*-flavor events, however, the Gaussian form cannot accommodate the data over a wider range without the addition of very large distortion terms. In light-flavor events, the Gaussian with small additional distortion terms is able to describe the data over a wider range. The peak positions ξ^* for each hadron species in light-flavor events are more consistent with a monotonic dependence on hadron mass than those in flavor-inclusive events. Our data are thus consistent with the predictions of MLLA QCD for the light flavors, and indicate that the presence of heavy hadrons distorts the observed spectra.

Using the large forward-backward asymmetry induced by the polarized SLC electron beam to separate light quark from light antiquark hemispheres, we have compared hadron and antihadron production in light quark jets. A large excess of proton over antiproton production at high x_p is direct evidence for the production of leading baryons, i.e. baryons that carry the quantum numbers of the quark that initiated the jet. A similarly large excess of K^- over K^+ production indicates not only leading kaon production but also that leading kaons are produced more often from initial *s* quarks than from initial *u* quarks. A smaller excess of π^- over π^+ production was also observed. This is expected if leading pions account for a large fraction of high momentum pion production but with equal π^- rates from initial *d* and $\bar{u}$ jets: $u\bar{u}$ and $d\bar{d}$ events then contribute to the observed difference with opposite signs, but since down-type quarks are produced more often and with a higher electroweak asymmetry than up-type quarks, a net positive difference is observed. These data provide unique and stringent tests of hadronization models. All models tested were able to reproduce the pion and kaon data, though the latter favor the JETSET model over the other two; JETSET is also consistent with the proton data, but the other two models predict values much higher than the data for $x_p > 0.3$.

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