

Parity and time reversal in the spin-rotation interaction

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A recently reported discrepancy between experimental and theoretical values of the muon's $g-2$ factor is interpreted as due to small violations of the conservation of P and T in the spin-rotation coupling. The experiments place an upper limit on these violations and on the weight change of spinning gyroscopes.

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The spin-rotation effect described by Mashhoon [1] attributes an energy $-(\hbar/2)\vec{\omega}\cdot\vec{\sigma}$ to a spin- $\frac{1}{2}$ particle in a frame rotating with angular velocity ω relative to an inertial frame. Identical results have been derived directly from the Dirac equation by means of the tetrad formalism [2–4]. The effect extends our knowledge of rotational inertia to the quantum level and violates the principle of equivalence [5] while preserving invariance under P and T . It has physical and astrophysical implications [3,6–8] and also plays a fundamental role in precise measurements of the $g-2$ factor of the muon [8].

Recently, a discrepancy $a_\mu(\text{exp})-a_\mu(\text{SM})=43\times 10^{-10}$ has been observed [9] between the experimental and standard model values of the muon's anomalous g value, $a_\mu=(g-2)/2$. This discrepancy can be used to set an upper limit on P and T invariance violations in spin-rotation coupling.

The possibility that discrete symmetries in gravitation are not conserved has been considered by some authors [10–13]. Attention has in general focused on the potential (in units $\hbar=c=1$)

$$U(\vec{r}) = \frac{GM}{r} [\alpha_1 \vec{\sigma} \cdot \hat{r} + \alpha_2 \vec{\sigma} \cdot \vec{v} + \alpha_3 \hat{r} \cdot (\vec{v} \times \vec{\sigma})], \quad (1)$$

which applies to a particle of generic spin $\vec{\sigma}$. The first term, introduced by Leitner and Okubo [11], violates the conservation of P and T . The same authors determined the upper limit $\alpha_1 \leq 10^{-11}$ from the hyperfine splitting of the ground state of hydrogen. The upper limit $\alpha_2 \leq 10^{-3}$ was determined in Ref. [13] from SN 1987A data. The corresponding potential violates the conservation of P and C . Conservation of C and T is violated by the last term, while Eq. (1), as a whole, conserves CPT . There is, as yet, no upper limit on α_3 . These studies are extended here to the Mashhoon term.

The $g-2$ experiment involves muons in a storage ring [14]. As the muons decay, the angular distribution of those electrons projected forward in the direction of motion reflect the precession of the muon spin along the cyclotron orbits.

Assume that all quantities in the effective Hamiltonian are time-independent and referred to a left-handed turn of axes comoving with the muons and rotating about the x_2 axis in the clockwise direction of motion of the muons. The x_3 axis

is tangent to the orbits and in the direction of the muon momentum. The magnetic field is $B_2 = -B$. Of all the terms that appear in the Dirac Hamiltonian, only the Mashhoon term couples the helicity states of the muon. The remaining terms contribute to the overall energy E of the states and the corresponding part of the Hamiltonian is indicated by H_0 [8].

Before decay, the muon states can be represented by

$$|\psi(t)\rangle = a(t)|\psi_+\rangle + b(t)|\psi_-\rangle, \quad (2)$$

where $|\psi_+\rangle$ and $|\psi_-\rangle$ are the right and left helicity states of the Hamiltonian H_0 and satisfy the equation

$$H_0|\psi_{+,-}\rangle = E|\psi_{+,-}\rangle. \quad (3)$$

Assume now that the coupling of rotation to $|\psi_+\rangle$ differs in strength from that to $|\psi_-\rangle$. Then the Mashhoon term can be altered by means of a matrix $A = \begin{pmatrix} \kappa_1 & 0 \\ 0 & \kappa_2 \end{pmatrix}$ that reflects the different coupling of rotation to the two helicity states. The total effective Hamiltonian is $H_{\text{eff}} = H_0 + H'$, where

$$H' = -\frac{1}{2}A\omega_2\sigma_2 + \mu B\sigma_2, \quad (4)$$

$\mu = (1 + a_\mu)\mu_0$ represents the total magnetic moment of the muon and μ_0 is the Bohr magneton. A violation of P and T in Eq. (4) would arise through $\kappa_2 - \kappa_1 \neq 0$. The constants κ_1 and κ_2 are assumed to differ from unity by small amounts ϵ_1 and ϵ_2 .

The coefficients $a(t)$ and $b(t)$ in Eq. (1) evolve in time according to

$$i\frac{\partial}{\partial t} \begin{pmatrix} a(t) \\ b(t) \end{pmatrix} = M \begin{pmatrix} a(t) \\ b(t) \end{pmatrix}, \quad (5)$$

where

$$M = \begin{pmatrix} E - i\frac{\Gamma}{2} & i\left(\kappa_1\frac{\omega_2}{2} - \mu B\right) \\ -i\left(\kappa_2\frac{\omega_2}{2} - \mu B\right) & E - i\frac{\Gamma}{2} \end{pmatrix}, \quad (6)$$

and Γ represents the width of the muon. The spin-rotation term is off-diagonal in Eq. (6) and does not therefore couple to matter universally. It violates Hermiticity [15]. It also violates T , P and PT , as stated, while nothing can be said about

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CPT conservation which requires H_{eff} to be Hermitian [16,17]. Because of the non-Hermitian nature of Eq. (4), one expects Γ itself to be non-Hermitian. The resulting corrections to the width of the muon are, however, of second order in the ϵ 's and are neglected.

M has eigenvalues

$$\begin{aligned} h_1 &= E - i\frac{\Gamma}{2} + R \\ h_2 &= E - i\frac{\Gamma}{2} - R, \end{aligned} \quad (7)$$

where

$$R = \sqrt{\left(\kappa_1 \frac{\omega_2}{2} - \mu B\right) \left(\kappa_2 \frac{\omega_2}{2} - \mu B\right)}, \quad (8)$$

and eigenstates

$$\begin{aligned} |\psi_1\rangle &= b_1[\eta_1|\psi_+\rangle + |\psi_-\rangle], \\ |\psi_2\rangle &= b_2[\eta_2|\psi_+\rangle + |\psi_-\rangle]. \end{aligned} \quad (9)$$

One also finds

$$\begin{aligned} |b_1|^2 &= \frac{1}{1 + |\eta_1|^2} \\ |b_2|^2 &= \frac{1}{1 + |\eta_2|^2} \end{aligned} \quad (10)$$

and

$$\eta_1 = -\eta_2 = \frac{i}{R} \left(\kappa_1 \frac{\omega_2}{2} - \mu B \right). \quad (11)$$

Then the muon states (2) are

$$|\psi(t)\rangle = \frac{1}{2} e^{-iEt - \Gamma t/2} [-2i\eta_1 \sin Rt |\psi_+\rangle + 2 \cos Rt |\psi_-\rangle], \quad (12)$$

where the condition $|\psi(0)\rangle = |\psi_-\rangle$ has been applied. The spin-flip probability is therefore

$$\begin{aligned} P_{\psi_-\rightarrow\psi_+} &= |\langle\psi_+|\psi(t)\rangle|^2 \\ &= \frac{e^{-\Gamma t}}{2} \frac{\kappa_1 \omega_2 - 2\mu B}{\kappa_2 \omega_2 - 2\mu B} [1 - \cos 2Rt]. \end{aligned} \quad (13)$$

When $\kappa_1 = \kappa_2 = 1$, Eq. (13) yields [8]

$$P_{\psi_-\rightarrow\psi_+} = \frac{e^{-\Gamma t}}{2} \left[1 - \cos \left(a_\mu \frac{eB}{m} t \right) \right], \quad (14)$$

that provides the appropriate description of the spin-rotation contribution to the spin-flip transition probability. Notice that the case $\kappa_1 = \kappa_2 = 0$ (no spin-rotation coupling) yields

$$P_{\psi_-\rightarrow\psi_+} = \frac{e^{-\Gamma t}}{2} \left[1 - \cos \left((1 + a_\mu) \frac{eB}{m} t \right) \right] \quad (15)$$

and does not therefore agree with the results of the $g-2$ experiments. Hence the necessity of accounting for spin-rotation coupling whose contribution cancels the factor eB/m in Eq. (15) [8].

Substituting $\kappa_1 = 1 + \epsilon_1, \kappa_2 = 1 + \epsilon_2$ into Eq. (13), one finds

$$P_{\psi_-\rightarrow\psi_+} \simeq \frac{e^{-\Gamma t}}{2} \left[1 - \cos \frac{eB}{m} (a_\mu - \epsilon) t \right], \quad (16)$$

where $\epsilon = \frac{1}{2}(\epsilon_1 + \epsilon_2)$. One may attribute the discrepancy between $a_\mu(\text{exp})$ and $a_\mu(SM)$ to a violation of the conservation of the discrete symmetries by the spin-rotation coupling term in Eq. (4). The upper limit on the violation of P, T and PT is derived from Eq. (16) assuming that the deviation from the current value of $a_\mu(SM)$ is wholly due to ϵ . The upper limit is therefore 43×10^{-10} .

Some more information can be extracted from current a_μ data. One may in fact assume that the coupling of rotation to the two helicity states of the fermion is opposite. In this case the parameters have values $\kappa_1 = 1, \kappa_2 = -1$. This is the anti-Hermitian limit of the interaction. The oscillation frequency is then

$$\begin{aligned} R &= \frac{1}{2} \sqrt{(2\mu B)^2 - \omega_2^2} = \frac{eB}{2m} \sqrt{2a_\mu + a_\mu^2} \\ &\simeq \frac{eB}{m} \sqrt{a_\mu/2} \end{aligned} \quad (17)$$

and Eq. (13) gives

$$P_{\psi_-\rightarrow\psi_+} \simeq \frac{e^{-\Gamma t}}{2} \frac{a_\mu}{2 + a_\mu} \left[1 - \cos \left(\frac{eB}{m} \sqrt{2a_\mu} t \right) \right]. \quad (18)$$

Equations (18) and (14) differ in amplitude and frequency. In fact the amplitude of Eq. (18) is much smaller than that of Eq. (14) while its frequency is higher than that actually observed. The choice $\kappa_1 = 1, \kappa_2 = -1$ is not therefore supported experimentally.

It also follows from Eqs. (4), (5) and (6) that the weight of a rotating object depends on its direction of rotation. The problem has been studied experimentally in [18]. No theoretical motivation for the study has ever been presented. An upper limit on this effect can be obtained in the present model from Eq. (4). The eigenstate energy difference due to spin-rotation coupling is in fact

$$-i \frac{\omega_2}{2} (\langle\psi_-|\sigma^2|\psi_+\rangle + \langle\psi_+|\sigma^2|\psi_-\rangle) = \frac{\omega_2}{2} (1 + \epsilon). \quad (19)$$

The additional energy difference is therefore $\epsilon\omega_2/2$, where $\epsilon = -43 \times 10^{-10}$. Returning to normal units, the correspond-

ing decrease in mass for a muon of positive helicity is $\Delta m = -\epsilon e B \hbar / mc^2 \approx -3.1 \times 10^{-48}$ g and the decrease in weight is $g \Delta m \approx 4 \times 10^{-45}$ dyn.

The fraction of total rotational energy associated with the effect is $\epsilon/2$. If one applies this result to all the particles of the gyroscope used in the experiment of [18], then one finds that the energy difference of the two rotation states of the

body is at most $(\epsilon/2) \frac{1}{2} I \omega^2 \approx 2.4$ erg, corresponding to a change in mass $\leq 2.6 \times 10^{-21}$ g and a change in weight $\leq 2.6 \times 10^{-18}$ dyn, in agreement with the null experimental results of [19,20].

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