

Wave function discord

Alexander Vilenkin*

Institute of Cosmology, Department of Physics and Astronomy, Tufts University, Medford, Massachusetts 02155

(Received 29 April 1998; published 27 August 1998)

Linde's proposal of a Euclidean path integral with the "wrong" sign of Euclidean action is often identified with the tunneling proposal for the wave function of the universe. However, the two proposals are in fact quite different. I illustrate the difference and point out that the recent criticism by Hawking and Turok does not apply to the tunneling proposal. [S0556-2821(98)07218-X]

PACS number(s): 98.80.Hw, 98.80.Bp

The debate about the form of the wave function of the universe has recently intensified [1–5]. Its most recent round was initiated by Bousso and Hawking [1] who claimed that the tunneling proposal for the wave function of the universe leads to a catastrophic instability of de Sitter space with respect to the pair production of black holes. This claim was analyzed in Ref. [2] and was shown to be unfounded. Linde [3] has argued that in models of the quantum creation of an open universe, the Hartle-Hawking wave function leads to unacceptably low values of the density parameter, and therefore the tunneling proposal should be preferred. In response, Hawking and Turok [4] asserted that with the choice of the tunneling wave function all perturbations about a homogeneous cosmological background become unstable and therefore this wave function can be meaningfully defined only in homogeneous minisuperspace models. This led them to conclude that the Hartle-Hawking wave function is the only proposal "with some pretensions to completeness" [5].

While the level of rhetoric is high, I think that progress on this issue may be helped by pointing out that what goes under the name of "tunneling wave function" is in fact two completely different wave functions. The debate will certainly gain in clarity if this fact is fully recognized.

Let me now briefly review the different proposals. The Hartle-Hawking wave function is given by the integral [6]

$$\psi_{HH} = \int e^{-S_E}, \quad (1)$$

where S_E is the Euclidean action and the integration is taken over compact Euclidean geometries and matter fields with a specified field configuration at the boundary. The Euclideanization is achieved by the standard Wick rotation, $t \rightarrow -i\tau$, from a Lorentzian path integral. Linde [7] suggested that the Euclidean rotation should be performed in the opposite sense, yielding

$$\psi_L = \int e^{+S_E}. \quad (2)$$

Finally, I introduced the tunneling wave function ψ_T which is specified either by the tunneling boundary condition [8] or by a Lorentzian path integral,

$$\psi_T = \int e^{iS}, \quad (3)$$

interpolating between "nothing" and a specified field configuration [9,10]. Arguments that the two definitions are equivalent were presented in [12].

The Hawking-Turok criticism [4] is directed against Linde's wave function (2). They point out that while the anti-Wick rotation $t \rightarrow +i\tau$ may work in simple minisuperspace models, it leads to disastrous consequences when inhomogeneous modes are included. All such modes become unstable, resulting in a breakdown of the semiclassical description of the universe. Linde has argued that this problem might be avoided if the Euclidean rotation is performed separately for the background fields and for the perturbations [3], but no prescription of this sort has yet been suggested that would apply in the general case. On the other hand, for the tunneling wave function it was shown in [13] that the gravitational and matter fields are stable and are in the same quantum state as for the Hartle-Hawking wave function.

If the reader needs further convincing that the tunneling and Linde's wave functions are indeed different, I would like to illustrate the difference for the simplest de Sitter minisuperspace model where both wave functions are well behaved. In this model, the universe is assumed to be homogeneous, isotropic, closed, and filled with a vacuum of constant energy density ρ_v [14]. The radius of the universe, a , is the only variable of the model, and the wave function $\psi(a)$ satisfies the Wheeler-DeWitt equation

$$\left[\frac{d^2}{da^2} - a^2(1 - H^2 a^2) \right] \psi(a) = 0. \quad (4)$$

Here, $H^2 = 8\pi G\rho_v/3$ and I have disregarded the ambiguity in the ordering of non-commuting operators a and d/da . (This ambiguity is unimportant in the semiclassical approximation which I am going to use below.)

Equation (4) has the form of a one-dimensional Schrödinger equation for a "particle" described by a coordinate $a(t)$, having zero energy, and moving in a potential

$$U(a) = a^2(1 - H^2 a^2). \quad (5)$$

The classically allowed region is $a \geq H^{-1}$, and the WKB solutions of Eq. (4) in this region are

*Electronic address: vilenkin@cosmos2.phy.tufts.edu

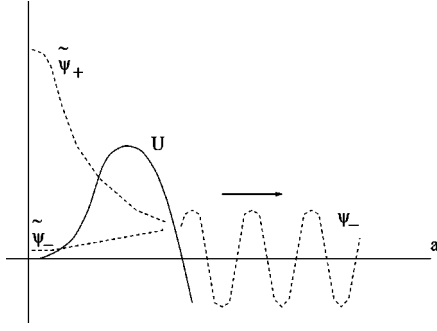


FIG. 1. Tunneling wave function for the de Sitter minisuperspace model. The potential $U(a)$ is shown by a solid line and the wave function by a dashed line.

$$\psi_{\pm}(a) = [p(a)]^{-1/2} \exp\left[\pm i \int_{H^{-1}}^a p(a') da' \mp i\pi/4\right], \quad (6)$$

where $p(a) = [-U(a)]^{1/2}$. The under-barrier, $a < H^{-1}$, solutions are

$$\tilde{\psi}_{\pm}(a) = |p(a)|^{-1/2} \exp\left[\pm \int_a^{H^{-1}} |p(a')| da'\right]. \quad (7)$$

The classical momentum conjugate to a is $p_a = -a\dot{a}$. For $a \gg H^{-1}$, we have

$$(-i d/da) \psi_{\pm}(a) \approx \pm p(a) \psi_{\pm}(a), \quad (8)$$

and thus $\psi_-(a)$ and $\psi_+(a)$ describe an expanding and a contracting universe, respectively. The tunneling boundary condition requires that only the expanding component should be present at large a ,

$$\psi_T(a > H^{-1}) = \psi_-(a). \quad (9)$$

The under-barrier wave function is found from the WKB connection formula,

$$\psi_T(a < H^{-1}) = \tilde{\psi}_+(a) - \frac{i}{2} \tilde{\psi}_-(a). \quad (10)$$

The growing exponential $\tilde{\psi}_-(a)$ and the decreasing exponential $\tilde{\psi}_+(a)$ have comparable amplitudes at the nucleation point $a = H^{-1}$, but away from that point the decreasing exponential dominates (see Fig. 1).

The Hartle-Hawking wave function is specified by requiring that it be given by $\exp(-S_E)$ in the Euclidean under-barrier regime. This gives [6]

$$\psi_H(a < H^{-1}) = \tilde{\psi}_-(a), \quad (11)$$

$$\psi_H(a > H^{-1}) = \psi_+(a) - \psi_-(a). \quad (12)$$

Linde's wave function is obtained by reversing the sign of the exponential in the Euclidean regime,

$$\psi_L(a < H^{-1}) = \tilde{\psi}_+(a), \quad (13)$$

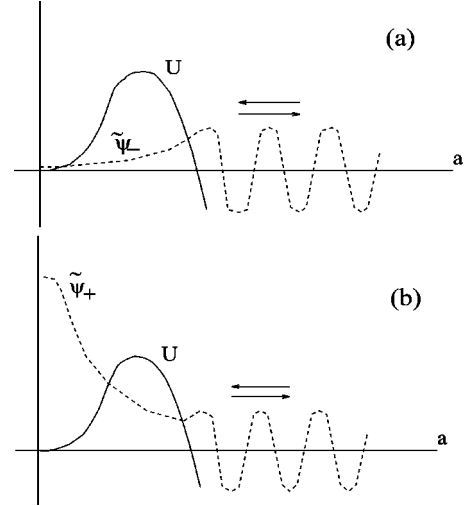


FIG. 2. Hartle-Hawking (a) and Linde (b) wave functions for de Sitter minisuperspace model.

and the continuation to the classically allowed range of a gives

$$\psi_L(a > H^{-1}) = \frac{1}{2} [\psi_+(a) + \psi_-(a)]. \quad (14)$$

The Hartle-Hawking and Linde wave functions are schematically represented in Fig. 2.

We can now compare Linde's and tunneling wave functions. Under the barrier, both of them are dominated by the decaying exponential $\tilde{\psi}_+(a)$. For this reason, ψ_L and ψ_T give similar expressions for the probability of the creation of the universe [3,9]. But that is where the similarity ends. The sub-dominant growing exponential, which is absent in Linde's wave function, was crucial in the derivation of the quantum state of inhomogeneous modes for the tunneling wave function [13]. Moreover, outside the barrier, Linde's wave function includes expanding and contracting universe components with equal amplitudes, and in this respect it is more similar to the Hartle-Hawking than to the tunneling wave function.

In conclusion, Linde's wave function ψ_L and tunneling wave function ψ_T are two different wave functions. Using the term "tunneling wave function" for ψ_L and ψ_T indiscriminately, as was done, e.g., in Refs. [1, 3], has introduced much confusion in the debate. The recent objection raised by Hawking and Turok [4] applies to ψ_L but not to ψ_T , and so the tunneling wave function remains a viable choice for the wave function of the universe.

I would like to add that, in following this wave function debate, the reader should be aware that all three wave functions are far from being rigorously defined mathematical objects. Except in the simplest models, the actual calculations of these wave functions involve additional assumptions which may appear reasonable, but are not really well justified. For a recent discussion of problems associated with defining and interpreting the cosmological wave function see, e.g., Ref. [12].

- [1] R. Bousso and S. W. Hawking, *Phys. Rev. D* **54**, 6312 (1996).
- [2] J. Garriga and A. Vilenkin, *Phys. Rev. D* **56**, 2464 (1997).
- [3] A. D. Linde, [gr-qc/9802038](#).
- [4] S. W. Hawking and N. G. Turok, [gr-qc/9802062](#).
- [5] N. G. Turok and S. W. Hawking, [hep-th/9803156](#).
- [6] J. B. Hartle and S. W. Hawking, *Phys. Rev. D* **28**, 2960 (1983).
- [7] A. D. Linde, *Lett. Nuovo Cimento* **39**, 401 (1984).
- [8] A. Vilenkin, *Phys. Rev. D* **33**, 3560 (1986).
- [9] A. Vilenkin, *Phys. Rev. D* **30**, 509 (1984).
- [10] In my first papers on quantum cosmology [11] I used the standard Wick rotation, but later I realized [9] that it did not correspond to the picture of “tunneling from nothing” which I had in mind from the beginning.
- [11] A. Vilenkin, *Phys. Lett.* **117B**, 25 (1982); *Phys. Rev. D* **27**, 2848 (1983).
- [12] A. Vilenkin, *Phys. Rev. D* **50**, 2581 (1994).
- [13] T. Vachaspati and A. Vilenkin, *Phys. Rev. D* **37**, 898 (1988).
- [14] This difference was first pointed out to me by Slava Mukhanov (private communication).