

Leptophobic Z' in stringy flipped SU(5)

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We show that leptophobic Z' gauge bosons occur naturally in flipped SU(5) and may shift R_b in an interesting way without upsetting the good values of Γ_{had} and R_c . Within a string-derived version of the model, we study three possible scenarios and the constraints imposed on model building that would allow the new symmetry to remain unbroken down to low energies. Such a Z' gauge boson has generation nonuniversal couplings to quarks that violate parity maximally in the up-quark sector, and may contribute significantly to spin asymmetries in polarized pp scattering experiments now being prepared for BNL RHIC. [S0556-2821(97)02101-2]

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I. INTRODUCTION

Recent experimental results concerning the rate of b quark production at the CERN e^+e^- collider LEP 1 ($R_b^{\text{expt}}=0.2179\pm 0.0012$ [1], $R_b^{\text{SM}}=0.2157$) and the distribution of high- E_T jets at the Fermilab Tevatron [as reported by the Collider Detector at Fermilab (CDF) Collaboration [2]], that may indicate the first departure from standard model expectations, have revived interest in Z' models as possible beyond-the-standard-model explanations for these phenomena [3–10]. In order not to disturb the agreement with standard model expectations in the lepton sector at LEP, the new Z' scenarios call for a leptophobic Z' that mixes with the regular Z .¹ Such a Z' may shift R_b in an interesting way if its couplings to quarks are suitably chosen. However one must ensure that the total hadronic width (Γ_{had}) remains essentially unchanged relative to standard model expectations, as it agrees rather well with observations. A new constraint in this class of models has recently arisen, in that the R_c ratio is now found to agree rather well with standard model expectations ($R_c^{\text{expt}}=0.1715\pm 0.0012$ [1], $R_c^{\text{SM}}=0.172$). Traditional Z' searches at hadron colliders (i.e., via its decay into charged leptons [13]) are not sensitive to a leptophobic Z' , whereas $Z'\rightarrow jj$ searches [14,15] are hampered by much larger backgrounds, although excluded regions of parameter space can still be obtained. The influence of a leptophobic Z' may also be felt via additional contributions to the top-

quark cross section [7], enhanced $b\bar{b}$ event yields at the Tevatron [9], and spin asymmetries in polarized pp collisions at the BNL Relativistic Heavy Ion Collider (RHIC) [16].

If such a Z' explanation to the seemingly anomalous data is to be taken seriously, one must provide a consistent theoretical framework, where the new gauge boson and its required properties arise naturally. In our minds that should be taken to be string model building, where the thorny question of cancellation of anomalies is dealt with automatically. New light neutral gauge bosons were early on considered to be the “smoking guns” of string, back when E_6 was the proverbial string-inspired gauge group [17]. The popularity of string Z' s, however, waned as non-grand-unified gauge groups became the natural outcome of actual string-derived models, such as $SU(5)\times U(1)$ [“flipped SU(5)”], $SU(3)^3$, $SU(4)\times SU(2)\times SU(2)$ (“Pati-Salam” models), and $SU(3)\times SU(2)\times U(1)$ (“standardlike” models). In these models one in fact has an excess of U(1) gauge groups at the string scale, but they typically get broken in the vacuum shifting process required to cancel the anomalous $U_A(1)$ that is characteristic of this type of string constructions. It is then interesting to explore whether string models can accommodate such gauge bosons at all, a search that is likely to lead to new and restrictive constraints on string model building. In this paper we explore this question in the context of stringy flipped SU(5).

We first show that leptophobia is very natural in flipped SU(5), and that it may provide the shift in R_b that seems to be preferred experimentally (Sec. II), while at the same time keeping Γ_{had} and R_c essentially unchanged. We then explore three scenarios for possible leptophobic Z' gauge bosons in a string-derived version of flipped SU(5) (Sec. III). Our second scenario is particularly compelling, entailing generation non-universal Z' couplings to quarks; it is further studied in Sec. IV. We also discuss the current experimental limits on and the prospects for detecting such Z' gauge bosons, which pos-

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¹If the Z' is also to explain the CDF data, it should be unusually broad and heavy [3]. As D0 has studied a much larger data sample than that reported by CDF, and no anomalous distributions have been observed [11], we do not consider Z' gauge bosons with such characteristics here. For a comparison of the CDF and D0 results, see Ref. [12].

ness parity-violating Z' couplings to up-type quarks, and may yield observable contributions to spin asymmetries in polarized pp scattering experiments being prepared for RHIC (Sec. V). Section VI summarizes our conclusions.

II. FLIPPED LEPTOPHOBIC Z' AND R_b

Let us first point out that leptophobia is very *natural* in $SU(5)\times U(1)$, where the particle content of the standard model is contained in the representations

$$F = (\mathbf{10}, \frac{1}{2}) = \{Q, d^c, \nu^c\}, \quad \bar{f} = (\bar{\mathbf{5}}, -\frac{3}{2}) = \{L, u^c\},$$

$$\ell^c = (\mathbf{1}, \frac{5}{2}) = \{e^c\}. \quad (1)$$

The new Z' would be leptophobic if \bar{f} and ℓ^c (which contain the standard model leptons) are uncharged under the new $U(1)$, while most of the quarks (in F) could still couple to it. In contrast, in regular $SU(5)$ ($F = \{Q, u^c, e^c\}$, $\bar{f} = \{L, d^c\}$), $SO(10)$, or E_6 such a separation (symmetry-based leptophobia) is not possible. However leptophobia may still be achieved dynamically under certain circumstances, as demonstrated in Ref. [6] in the case of the subgroup of E_6 called the “ η model” [18].

The effect of a Z' mixing with the regular Z and its impact on LEP physics has been addressed previously, particularly in the context of R_b, R_c [3,5,6,9]. Here we present a succinct discussion, emphasizing the flipped $SU(5)$ novelties. Small Z - Z' mixing amounts to a shift in the vector and axial-vector couplings of the standard Z :

$$C_V = C_V^0 + \theta(g_{Z'}/g_Z)C_V', \quad C_A = C_A^0 + \theta(g_{Z'}/g_Z)C_A', \quad (2)$$

where θ is the small Z - Z' mixing angle, $g_Z = g/\cos\theta_w$ is the usual weak coupling, $g_{Z'}$ is the new U' gauge coupling, $C_{V,A}^0$ are the usual vector and axial-vector Z couplings to fermions, and $C'_{V,A}$ are related to the charges of the fermions under U' . As Eq. (1) shows, a leptophobic Z' in flipped $SU(5)$ violates parity in the up-quark sector, but not in the down-quark sector. Using $C_V' = Q(\psi_L) + Q(\psi_R)$ and $C_A' = -Q(\psi_L) + Q(\psi_R)$, we obtain

	C_V^0	C_A^0	$Q(\psi_L)$	$Q(\psi_R)$	C_V'	C_A'
up	$\frac{1}{2} - \frac{4}{3}x_w$	$\frac{1}{2}$	c	0	c	$-c$
down	$-\frac{1}{2} + \frac{2}{3}x_w$	$-\frac{1}{2}$	c	c	$2c$	0

(3)

where $x_w = \sin^2\theta_w$, and c is the U' charge of the $F = (\mathbf{10}, \frac{1}{2})$ representation. These results apply to each individual generation, which generally could have different charges under U' , i.e., $c_{1,2,3}$.

With the above information, which depends solely on the flipped $SU(5)$ origin of the leptophobic Z' , one can obtain the first-order shifts in $\Gamma_{bb} = \Gamma(Z \rightarrow b\bar{b})$, $\Gamma_{c\bar{c}} = \Gamma(Z \rightarrow c\bar{c})$, and $\Gamma_{\text{had}} = \Gamma(Z \rightarrow \text{hadrons}) = \Gamma_{d\bar{d}} + \Gamma_{s\bar{s}} + \Gamma_{b\bar{b}} + \Gamma_{u\bar{u}} + \Gamma_{c\bar{c}}$ (and therefore, R_b, R_c) in terms of the composite parameter

$$\delta \equiv \theta(g_{Z'}/g_Z). \quad (4)$$

Recalling that $\Gamma(Z \rightarrow q\bar{q}) = \Gamma_0(C_{V_q}^2 + C_{A_q}^2)$, where $\Gamma_0 = G_F M_Z^3 / (2\sqrt{2}\pi)$, one can easily determine the relevant first-order shifts (i.e., linear in δ) for arbitrary choices of the $c_{1,2,3}$ charges.

What should the U' charges of the $SU(5)$ multiplets be? This question can be answered exactly in explicit string models, as we discuss in Sec. IIIB. For now let us just state our choices:

$$F_0 \quad -\frac{1}{2}, \quad \bar{F}_4 \quad \frac{1}{2}, \quad \bar{f}_{2,3,5} \quad 0,$$

$$F_1 \quad -\frac{1}{2}, \quad \bar{F}_5 \quad 0, \quad \ell_{2,3,5}^c \quad 0,$$

$$F_2 \quad 0,$$

$$F_3 \quad 1,$$

$$F_4 \quad -\frac{1}{2}. \quad (5)$$

These charge choices reflect some principles: the leptons (in $\bar{f}_{2,3,5}, \ell_{2,3,5}^c$) are uncharged (leptophobia); there is a pair of $(\mathbf{10}, \mathbf{10})(F_2, \bar{F}_5)$ whose neutral components acquire GUT scale vevs and break $SU(5)\times U(1)$ in the standard way, and by virtue of being uncharged leave U' unbroken; the remaining representations enforce $\text{Tr}U' = 0$ and contain three generations of quarks and an extra $(\mathbf{10}, \mathbf{10})$ to allow string unification. The $c_{1,2,3}$ are then to be taken from the set $\{0(1), 1(1), -\frac{1}{2}(3)\}$, where the number in parentheses indicates the number of generations that may carry such a charge.

In Table I we display the 13 possible charge combinations, along with the fractional changes in Γ_{had} , R_b , and R_c in units of δ . The value of δ is constrained by its effect on the many electroweak observables, where it enters through a tree-level correction to the ρ parameter due to the Z - Z' mixing. Detailed fits to the electroweak data allow values of $|\delta|$ as large as $\sim 10^{-2}$ [3,6,9]. For the representative choice $\delta = 0.01$, in Table I we also display the actual shifts in Γ_{had} (in MeV), R_b , and R_c , all of which scale linearly with δ . To help decide which of the charge assignments satisfy the present experimental constraints, in Fig. 1 we plot the correlated values of ΔR_b versus $\Delta\Gamma_{\text{had}}$, for each of the 13 charge assignments (the lines are parametrized by δ). Note that different charge assignments may yield the same ΔR_b - $\Delta\Gamma_{\text{had}}$ correlation. Since the LEP measured value of Γ_{had} agrees with the standard model prediction within 1 MeV, and the experimental uncertainty in Γ_{had} is 3 MeV, we constrain possible charge assignments by requiring $|\Delta\Gamma_{\text{had}}| < 3$ MeV, as denoted by the vertical dashed lines in Fig. 1. Furthermore, we demand that ΔR_b be in the interval 0.0010–0.0034 (denoted by the horizontal dashed lines in Fig. 1), such that when added to the standard model prediction ($R_b^{\text{SM}} = 0.2157$, for $m_t = 175$ GeV) the resulting R_b falls within the experimental allowed window.² R_b^{expt}

²Such level of precision should suffice, as one still needs to include one-loop supersymmetric corrections to R_b , which are not expected to be enhanced [19], especially in supergravity models [20]. In any event, shifting ΔR_b somewhat does not change the set of allowed charge assignments.

TABLE I. The 13 possible assignments of U' charges to the three generations, along with the fractional changes in Γ_{had} , R_b , and R_c (in units of δ). For $\delta=0.01$ we also display the shift in Γ_{had} in MeV, and the actual shifts in R_b and R_c .

	c_1	c_2	c_3	$\Delta\Gamma_{\text{had}}/\Gamma_{\text{had}}$	$\Delta R_b/R_b$	$\Delta R_c/R_c$	$\Delta\Gamma_{\text{had}}$	ΔR_b	ΔR_c
1	0	$-\frac{1}{2}$	$-\frac{1}{2}$	1.00	0.87	0.06	17.5	0.0019	0.0001
2	0	$-\frac{1}{2}$	1	-0.23	-3.52	1.30	-4.0	-0.0076	0.0022
3	0	1	$-\frac{1}{2}$	-0.78	2.65	-1.36	-13.5	0.0057	-0.0023
4	$-\frac{1}{2}$	0	$-\frac{1}{2}$	1.00	0.87	-1.00	17.5	0.0019	-0.0017
5	$-\frac{1}{2}$	0	1	-0.23	-3.52	0.23	-4.0	-0.0076	0.0004
6	1	0	$-\frac{1}{2}$	-0.78	2.65	0.78	-13.5	0.0057	0.0013
7	$-\frac{1}{2}$	$-\frac{1}{2}$	0	1.19	-1.19	-0.12	20.7	-0.0026	-0.0002
8	$-\frac{1}{2}$	1	0	-0.59	0.59	-1.54	-10.3	0.0013	-0.0027
9	1	$-\frac{1}{2}$	0	-0.59	0.59	1.66	-10.3	0.0013	0.0029
10	1	$-\frac{1}{2}$	$-\frac{1}{2}$	-0.18	2.05	1.25	-3.2	0.0044	0.0021
11	$-\frac{1}{2}$	1	$-\frac{1}{2}$	-0.18	2.05	-1.95	-3.2	0.0044	-0.0034
12	$-\frac{1}{2}$	$-\frac{1}{2}$	1	0.36	-4.11	0.70	6.3	-0.0089	0.0012
13	$-\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$	1.60	0.27	-0.53	27.9	0.0006	-0.0009

$=0.2179 \pm 0.0012$ [1]. From Fig. 1 we conclude that only charge assignments 2,5,10,11,12 satisfy the experimental requirements of R_b and Γ_{had} .

To consider the effect of R_c , in Fig. 2 we plot the correlated values of ΔR_c versus ΔR_b . The latest experimental value of R_c ($R_c^{\text{expt}}=0.1715 \pm 0.0056$ [1]) is now in good agreement with the standard model prediction ($R_c^{\text{SM}}=0.172$). Only cases 2,5,11,12 satisfy the experimental requirements from R_b , R_c , and Γ_{had} .³ It is important to note that these charge assignments predict modest shifts in R_c : $|\Delta R_c| < 0.005$, which are below the present experimental sensitivity and, therefore, naturally “protect” the standard model prediction.

We are then left with four experimentally preferred charge assignments:

	c_1	c_2	c_3
2	0	$-\frac{1}{2}$	1
5	$-\frac{1}{2}$	0	1
11	$-\frac{1}{2}$	1	$-\frac{1}{2}$
12	$-\frac{1}{2}$	$-\frac{1}{2}$	1

(6)

We note that these charge assignments are unlike any that have been so far considered in the literature. In fact, the popular assumption of a generation-independent U' charge is violated explicitly in these four successful cases. The constraints imposed by flipped SU(5) make that case (13 in Table I) disfavored by experimental data. Note also that the first two assignments imply $\Delta\Gamma_{u\bar{u}}=\Delta\Gamma_{d\bar{d}}=0$ or $\Delta\Gamma_{c\bar{c}}=\Delta\Gamma_{s\bar{s}}=0$, respectively. These perhaps “unnatural” charge assignments occur quite naturally in the string model. For future reference, we point out that in the string model F_4 is required to contain the third-generation quarks in order

³As the data shows preference for $\Delta R_c < 0$ and $\Delta R_b > 0$, we exclude from further consideration charge assignments that entail shifts in the same direction, i.e., cases 1,6,7,9,10.

to generate the top-quark Yukawa coupling. From Eq. (5) we see that we must have $c_3 = -\frac{1}{2}$, which uniquely selects charge assignment (11) in Eq. (6). In this stringy preferred case we obtain the following relation:

$$\Delta R_b \approx 0.0042 \left(\frac{\Delta\Gamma_{\text{had}}}{-3 \text{ MeV}} \right) \quad (7)$$

(and also $\Delta R_c \approx -0.76 \Delta R_b$).

Before proceeding, we should address the question of flavor-changing neutral currents (FCNCs) that necessarily arise in our model because the generation nonuniversal couplings of the Z' to quarks spoil the Glashow-Iliopoulos-Maiani (GIM) cancellation. For a leptophobic Z' the tree-level Z' -mediated contributions to Δm_K , Δm_D , and Δm_B should provide the strongest constraints. As usual, in obtaining the CKM matrix $V=U_L D_L^\dagger$ there is freedom of how to choose U_L and D_L such that V agrees with experimental

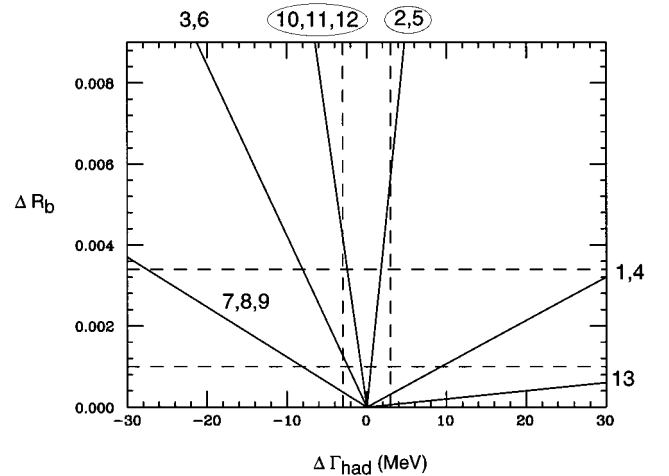


FIG. 1. The correlated shifts in R_b and Γ_{had} in flipped SU(5) for the various U' charge assignment combinations shown in Table I. The dashed lines delimit the experimental uncertainty in Γ_{had} and the required shift in R_b to fall within the experimental limits. The circled charge assignments (2,5,10,11,12) agree with experiment.

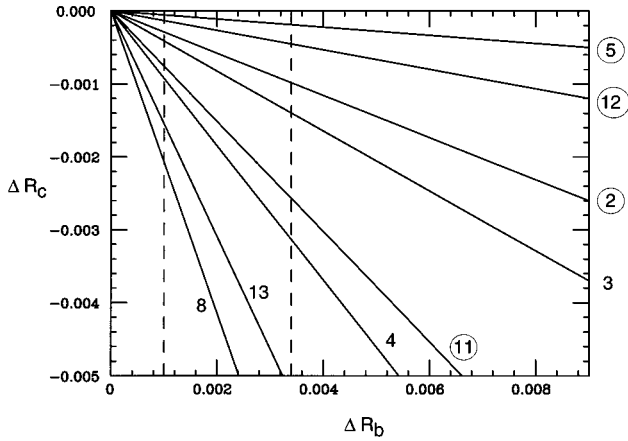


FIG. 2. The correlated shifts in R_c and R_b in flipped SU(5) for the various U' charge assignment combinations shown in Table I. (Combinations not shown entail shifts in R_c and R_b in the same direction.) The dashed lines delimit the required shift in R_b to fall within the experimental limits. The circled charge assignments (2,5,11,12) agree with experimental data on Γ_{had} (see Fig. 1).

observations. In Ref. [8] it is shown that taking $U_L=1$, the experimental value of Δm_B forces our parameter δ to uninterestingly small values. Taking instead $D_L=1$ automatically satisfies the Δm_K and Δm_B constraints, leaving only $\Delta m_D < 10^{-13}$ GeV to contend with. Following Ref. [8] we find $\delta(\Delta m_D) \approx (7 \times 10^{-6} \text{ GeV}) |\delta|^2 |(VgV^\dagger)_{12}|^2 [f_D / (0.22 \text{ GeV})]^2$, where we have used $V=U_L$. In our preferred charge assignment (11) we have $g = \text{diag}(-\frac{1}{2}, 1, -\frac{1}{2})$, and we obtain $|(VgV^\dagger)_{12}| = \frac{1}{2} |V_{12} V_{22}^*| \approx \frac{1}{10}$. Given the inherent uncertainties in this kind of calculations, we conclude that our above choice of $|\delta| \sim 10^{-2}$ is consistent with FCNC constraints at the present time.

III. STRINGY FLIPPED SU(5) LEPTOPHOBIA

In the context of stringy flipped SU(5), we will consider four criteria in the search for suitable U(1) gauge symmetries: they should be anomaly-free, leptophobic, unbroken, and mixed. The anomaly-free requirement can only be addressed in the context of a fundamental theory (such as string theory), where anomaly-free symmetries are automatic (although they may need to be broken). The (low-energy) leptophobia requirement is dictated by phenomenology and can be enforced by symmetry, or achieved dynamically via mixing in the gauge kinetic functions [21]. The new U(1) should remain unbroken after the vacuum shifting at the Planck scale that is required to cancel the anomalous $U_A(1)$. Finally, since the new gauge boson must be leptophobic, it cannot be produced directly at LEP, and thus can affect LEP physics only via mixing with the regular Z.

We seek the above properties in the context of a string-derived version of the $SU(5) \times U(1)$ model [22] that has been shown to possess various desirable properties regarding the vacuum energy, string unification, the dynamical generation of all mass scales, the top-quark mass, and the strong coupling [23]. The complete gauge group of the model breaks up into three identifiable pieces $G = G_{\text{obs}} \times G_{\text{hidden}} \times G_{U(1)}$, where $G_{\text{obs}} = SU(5) \times U(1)$, $G_{\text{hidden}} = SU(4) \times SO(10)$, and $G_{U(1)} = U_1(1) \times U_2(1) \times U_3(1) \times U_4(1) \times U_5(1)$. The

massless matter fields in the model are listed in Tables II–V, along with their charges under $G_{U(1)}$. Of note is the fact that $\text{Tr} U_{1,2,3,5} \neq 0$, whereas $\text{Tr} U_4 = 0$. These seemingly anomalous symmetries are artifacts of the truncation of the full string spectrum down to the massless sector. The low-energy effective theory is correctly specified by first “rotating” all the anomaly into a single anomalous $U_A \propto \sum_{i=1,2,3,5} [\text{Tr} U_i] U_i$ [24], and then adding a one-loop correction to the D term corresponding to U_A : $D_A \rightarrow D_A + \epsilon M^2$, where $\epsilon = g^2 \text{Tr} U_A / 192 \pi^2$ [25].

The mass spectrum of all the states in the tables can be obtained in a complicated procedure that takes into account the trilinear and nonrenormalizable contributions to the superpotential and, therefore, to the masses and interactions [22]. Such procedure does not have a unique outcome because the values of the vevs of the singlet fields in Table III are unknown (although constrained by the anomalous U_A cancellation conditions). The objective of this exercise is to obtain an electroweak-scale spectrum that resembles closely that of the MSSM. For instance, in Table II only one pair of Higgs pentaplets (their doublet components only) should survive to low energies. Below we will find constraints on which states must belong to the low-energy sector of the theory. We also note that the possible U' symmetries that we identify below do not change the predictions of the string model for the superpotential Yukawa couplings, as these symmetries are derived along with the rest of the model. The problem of hierarchical fermion mass generation finds a possible resolution in string models. As is well known, several internal (world-sheet) symmetries severely restrict the allowed Yukawa couplings, typically predicting nonzero trilinear couplings for only the third family. Nonrenormalizable terms usually contain the remaining Yukawa couplings, but these are naturally suppressed by powers of a small ratio that arises in the cancellation of the anomalous U_A . This program has met with some success in the present model [22].

We find three scenarios containing possibly new light neutral gauge bosons.

A. First scenario

The $U_4(1)$ gauge symmetry is traceless (anomaly free), the leptons ($\bar{f}_{2,3,5}, \mathcal{L}_{2,3,5}^c$) are not charged under it (leptophobic, see Table II), and it does not participate in the $U_A(1)$ cancellation mechanism (unbroken). Moreover, some of the quarks are charged under $U_4(1)$. The problem is that $U_4(1)$ and $U_Y(1)$ do not mix: the Higgs doublets that break the electroweak symmetry are uncharged under U_4 (see h_i, \bar{h}_i in Table II). The mixing via the gauge kinetic functions is not operative, as one can easily verify that $\text{Tr}(YU_4) = 0$. This fact “protects” the leptophobia, as otherwise the leptons would get their U_4 charges shifted away from zero.

The only representations charged under U_4 that may also contain quark fields are F_0, F_1 . Making this assignment also keeps U_4 unbroken through the $SU(5) \times U(1)$ symmetry-breaking process at the grand unified theory (GUT) scale. The U_4 symmetry may be broken radiatively at low energies if the singlet fields $\eta_{1,2}, \bar{\eta}_{1,2}$, which solely carry the U_4 charge (see Table III), acquire suitable dynamical vacuum expectation values.

As such this Z' will not affect LEP physics, as it does not obviously mix with the Z . However direct detection at hadron colliders via its hadronic decays or parity-violating spin asymmetries is possible. If we assume that F_0, F_1 contain first- and second-generation fields (F_4 is expected to contain the third-generation fields [22]) then the Z' couplings to quarks ($C'_{V,A}$) are as given in Eq. (3), but only for the first and second generations. Consulting Table II one finds $c_1 = -c_2 = \pm \frac{1}{2}$ and $c_3 = 0$ which, in the event that some mechanism for Z - Z' mixing were to be found, would imply $\Delta\Gamma_{\text{had}} = 0$ and $\Delta R_b = 0$, leaving R_b unshifted.

B. Second scenario

There are three linear combinations of $U_{1,2,3,5}$ that are orthogonal to $U_A = U_1 - 3U_2 + U_3 + 2U_5$, and therefore traceless. A suitable basis is provided by: $U'_1 = U_3 + 2U_5$, $U'_2 = U_1 - 3U_2$, $U'_3 = 3U_1 + U_2 + 4U_3 - 2U_5$. Is there a linear combination of $U'_{1,2,3}$ that is leptophobic? The leptons transform as $\bar{f}_{2,5}, \mathcal{L}_{2,5}^c: (0, \frac{3}{2}, -\frac{1}{2})$; $\bar{f}_3, \mathcal{L}_3^c: (\frac{3}{2}, 0, 1)$, from which it follows that there is a unique leptophobic linear combination: $U' \propto 2U'_1 - U'_2 - 3U'_3 \propto U_1 + U_3 - U_5$. This symmetry is by construction anomaly free and leptophobic, and some of the Higgs pentaplets are charged under it (i.e., mixed). The charges of all fields under U' are listed in the tables, along with the two additional traceless combinations that can be chosen to be $U'' = U_1 - 3U_2 + U_3 + 2U_5$ and $U''' = U_2 + U_3 + U_5$. From the tables we see that only a very limited set of fields is uncharged under U' ,

$$F_2, \bar{F}_5, \Phi_{31}, \bar{\Phi}_{31}, T_1, T_2, T_3, D_3, D_7, \quad (8)$$

and therefore their vevs will leave U' unbroken. The crucial question is whether the usual D - and F -flatness conditions may be satisfied with such a limited set of vevs. Moreover, this set of vevs generally breaks the hidden sector gauge group. The feasibility of this scenario will be addressed in more detail in Sec. IV below.

Let us assume that U' can indeed remain unbroken down to low energies. One can verify that $\text{Tr}(YU') = 0$ (at least for an appropriately chosen subset of light fields) and thus the U' charges are not shifted, and the leptophobia is protected. Previous studies [22] show that F_4 should contain the third-generation quarks. Also, F_2, \bar{F}_5 should be engaged in $SU(5) \times U(1)$ breaking, as they are neutral under U' , and therefore \bar{F}_4 should contain the additional representations required for string unification [23]. It remains to determine the fate of F_0 , F_1 , and F_3 , two of which should be assigned to contain the first- and second-generation quarks, whereas the third one will accompany \bar{F}_4 . Fortunately, the R_b, R_c analysis of Sec. II provides an important clue. Our assignment of F_4 entails $c_3 = -\frac{1}{2}$, and Eq. (6) then implies $c_1 = -\frac{1}{2}$ and $c_2 = 1$. This means that F_3 (which has $c = 1$) should contain the second generation, whereas F_0 or F_1 (or a linear combination thereof) should contain the first generation.

C. Third scenario

Since the unique leptophobic U' in the second scenario (in Sec. III B) requires the hidden sector gauge group

G_{hidden} to be broken in order to preserve D and F flatness, one may consider giving up *a priori* leptophobia, hoping to later be able to generate dynamic leptophobia via $U(1)$ mixing. This scenario may leave G_{hidden} and the desired $U(1)$ unbroken, and still preserve flatness. To this end we determine all possible linear combinations of U'_1, U'_2, U'_3 (or U', U'', U''') that are uncharged under more than one singlet field (in Table III), with the idea that vevs from this subset of fields should suffice to preserve flatness. We find four such linear combinations:

$$U_a = U_1 + U_2 + 2U_3,$$

$$\mathcal{S}_a = \{\Phi_{12}, \bar{\Phi}_{12}, \phi^+, \bar{\phi}^+, \phi^-, \bar{\phi}^-, \phi_{3,4}, \bar{\phi}_{3,4}\},$$

$$U_b = U_1 + U_2 + U_3 - 2U_5,$$

$$\mathcal{S}_b = \{\Phi_{12}, \bar{\Phi}_{12}, \Phi_{23}, \bar{\Phi}_{23}, \phi_{3,4}, \bar{\phi}_{3,4}\},$$

$$U_c = U_1 + U_2 - U_3 - 6U_5, \quad \mathcal{S}_c = \{\Phi_{12}, \bar{\Phi}_{12}, \phi_{45}, \bar{\phi}_{45}, \phi_{3,4}, \bar{\phi}_{3,4}\},$$

$$U_d = U_1 - 3U_2 + U_3 + 2U_5,$$

$$\mathcal{S}_d = \{\Phi_{31}, \bar{\Phi}_{31}, \phi^-, \bar{\phi}^-, \phi_{45}, \bar{\phi}_{45}\},$$

listed along with the singlet fields that are uncharged under them. It is easy to show that the D -flatness conditions specific to this model (as given in Ref. [22]) cannot be satisfied by vevs belonging to the subsets \mathcal{S}_a , \mathcal{S}_b , or \mathcal{S}_c . Subset \mathcal{S}_d may work. As one can verify, the leptons ($\bar{f}_{2,3,5}, \mathcal{L}_{2,3,5}^c$) couple universally to U_d , via a charge that would need to be shifted down to zero by the dynamical mixing mechanism [which is available as $\text{Tr}(YU_d) \neq 0$]. Unfortunately, universal leptophobia is not achievable because the dynamical shifts would be proportional to the hypercharge of the leptons, and thus would not be universal.

IV. SECOND SCENARIO REVISITED

There are several questions that need to be addressed in order to insure that a $U(1)$ symmetry at the string scale indeed becomes the U' symmetry with the desired properties at the electroweak scale. One must insure that this symmetry remains unbroken in the anomalous U_A cancellation mechanism, and upon $SU(5) \times U(1)$ symmetry breaking. After these hurdles have been cleared, on the way down to low energies the $g_{Z'}$ gauge coupling must evolve properly, and the gauge symmetry must be dynamically broken near the electroweak scale. These two effects give rise to the phenomenological parameter δ in Eq. (4), which must obtain a value in the range of interest ($|\delta| \sim 10^{-2}$). Addressing all these questions requires a very detailed study of the string model, which does not appear warranted at this point, given the uncertainty in the experimental measurements. Nonetheless, we will attempt to outline the main issues to be confronted in such a future study.

First of all, in the anomalous U_A cancellation mechanism three gauge symmetries must be broken U_A, U'', U''' , but U' must remain unbroken. Moreover, it does not help to have some linear combination involving U' to remain unbroken, as U' will be mixed and leptophobia would likely be

TABLE II. The observable sector massless matter fields and their transformation properties under $G_{U(1)}$. Under $SU(5) \times U(1)$ these fields transform as $F = (10, \frac{1}{2})$, $\bar{F} = (\bar{5}, -\frac{3}{2})$, $\ell^c = (1, \frac{5}{2})$, $h = (5, -1)$, and $\bar{h} = (\bar{5}, 1)$. Also indicated are the charges under U_A and three orthogonal linear combinations of interest (U', U'', U''').

	U_1	U_2	U_3	U_4	U_5	U_A	U'	U''	U'''
F_0	$-\frac{1}{2}$	0	0	$-\frac{1}{2}$	0	$\frac{3}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$	0
F_1	$-\frac{1}{2}$	0	0	$\frac{1}{2}$	0	$\frac{3}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$	0
F_2	0	$-\frac{1}{2}$	0	0	0	$\frac{1}{2}$	0	$\frac{3}{2}$	$-\frac{1}{2}$
F_3	0	0	$\frac{1}{2}$	0	$-\frac{1}{2}$	$\frac{3}{2}$	1	$-\frac{1}{2}$	0
F_4	$-\frac{1}{2}$	0	0	0	0	$\frac{3}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$	0
\bar{F}_4	$\frac{1}{2}$	0	0	0	0	$-\frac{3}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	0
\bar{F}_5	0	$\frac{1}{2}$	0	0	0	$-\frac{1}{2}$	0	$-\frac{3}{2}$	$\frac{1}{2}$
\bar{f}_2, ℓ_2^c	0	$-\frac{1}{2}$	0	0	0	$\frac{1}{2}$	0	$\frac{3}{2}$	$-\frac{1}{2}$
\bar{f}_3, ℓ_3^c	0	0	$\frac{1}{2}$	0	$\frac{1}{2}$	$\frac{1}{2}$	0	$\frac{3}{2}$	1
\bar{f}_5, ℓ_5^c	0	$-\frac{1}{2}$	0	0	0	$\frac{1}{2}$	0	$\frac{3}{2}$	$-\frac{1}{2}$
h_1	1	0	0	0	0	-3	1	1	0
\bar{h}_1	-1	0	0	0	0	3	-1	-1	0
h_2	0	1	0	0	0	-1	0	-3	1
\bar{h}_2	0	-1	0	0	0	1	0	3	-1
h_3	0	0	1	0	0	2	1	1	1
\bar{h}_3	0	0	-1	0	0	-2	-1	-1	-1
h_{45}	$-\frac{1}{2}$	$-\frac{1}{2}$	0	0	0	2	$-\frac{1}{2}$	1	$-\frac{1}{2}$
\bar{h}_{45}	$\frac{1}{2}$	$\frac{1}{2}$	0	0	0	-2	$\frac{1}{2}$	-1	$\frac{1}{2}$

lost. The vector-boson mass matrix is given by $\mathcal{M}_{\alpha\beta}^2 \propto \sum_a (\partial D_\alpha / \partial \phi_a) (\partial D_\beta / \partial \phi_a)^*$, where $D_\alpha = g \sum_i q_i^\alpha \phi_i \phi_i^*$, and $\alpha, \beta = U_A, U', U'', U'''$. The resulting 4×4 matrix must have zeros for all entries involving U' :

$$\mathcal{M}_{U', U'}^2 \propto \sum_a |\partial D_{U'} / \partial \phi_a|^2 = g^2 \sum_a (q'_a)^2 |\phi_a|^2 = 0, \quad (9)$$

is a necessary and sufficient condition, which implies that only fields uncharged under U' [those in Eq. (8)] should be allowed to acquire vevs. In this case the D -flatness conditions become

$$D_A : \frac{1}{2} (|V_2|^2 - |\bar{V}_5|^2) + \frac{5}{2} (|T_1|^2 + |T_2|^2 + |T_3|^2 + |D_3|^2 - |D_7|^2) - x_{31} + \epsilon M^2 = 0, \quad (10)$$

$$D' : 0 = 0, \quad (11)$$

$$D'' : \frac{3}{2} (|V_2|^2 - |\bar{V}_5|^2) = 0, \quad (12)$$

$$D''' : -\frac{1}{2} (|V_2|^2 - |\bar{V}_5|^2) + \frac{1}{2} (|T_1|^2 + |T_3|^2 + |D_3|^2 - |D_7|^2) - |T_2|^2 - x_{31} = 0, \quad (13)$$

TABLE III. The singlet fields and their transformation properties under $G_{U(1)}$, U_A , and three orthogonal linear combinations of interest (U', U'', U''').

	U_1	U_2	U_3	U_4	U_5	U_A	U'	U''	U'''
Φ_{12}	-1	1	0	0	0	2	-1	-4	1
$\bar{\Phi}_{12}$	1	-1	0	0	0	-2	1	4	-1
Φ_{23}	0	-1	1	0	0	3	1	4	0
$\bar{\Phi}_{23}$	0	1	-1	0	0	-3	-1	-4	0
Φ_{31}	1	0	-1	0	0	-5	0	0	-1
$\bar{\Phi}_{31}$	-1	0	1	0	0	5	0	0	1
ϕ_{45}	$\frac{1}{2}$	$\frac{1}{2}$	1	0	0	0	$\frac{3}{2}$	2	$\frac{3}{2}$
$\bar{\phi}_{45}$	$-\frac{1}{2}$	$-\frac{1}{2}$	-1	0	0	0	$-\frac{3}{2}$	-2	$-\frac{3}{2}$
ϕ^+	$\frac{1}{2}$	$-\frac{1}{2}$	0	0	1	-2	$-\frac{1}{2}$	4	$\frac{1}{2}$
$\bar{\phi}^+$	$-\frac{1}{2}$	$\frac{1}{2}$	0	0	-1	2	$\frac{1}{2}$	-4	$-\frac{1}{2}$
ϕ^-	$\frac{1}{2}$	$-\frac{1}{2}$	0	0	-1	0	$\frac{3}{2}$	0	$-\frac{3}{2}$
$\bar{\phi}^-$	$-\frac{1}{2}$	$\frac{1}{2}$	0	0	1	0	$-\frac{3}{2}$	0	$\frac{3}{2}$
$\phi_{3,4}$	$\frac{1}{2}$	$-\frac{1}{2}$	0	0	0	-1	$\frac{1}{2}$	2	$-\frac{1}{2}$
$\bar{\phi}_{3,4}$	$-\frac{1}{2}$	$\frac{1}{2}$	0	0	0	1	$-\frac{1}{2}$	-2	$\frac{1}{2}$
$\eta_{1,2}$	0	0	0	1	0	0	0	0	0
$\bar{\eta}_{1,2}$	0	0	0	-1	0	0	0	0	0
$\Phi_{0,1,3,5}$	0	0	0	0	0	0	0	0	0

where we have defined $x_{31} = |\Phi_{31}|^2 - |\bar{\Phi}_{31}|^2$, and $\epsilon = g^2 \text{Tr} U_A / 192 \pi^2 > 0$. These equations have the solution

$$|V_2| = |\bar{V}_5|, \quad (14)$$

$$x_{31} = -\frac{1}{4} \epsilon M^2 - \frac{15}{8} |T_2|^2 < 0, \quad (15)$$

$$|T_1|^2 + |T_3|^2 + |D_3|^2 - |D_7|^2 = -\frac{1}{2} \epsilon M^2 - \frac{7}{4} |T_2|^2 < 0. \quad (16)$$

Note that the last two equations imply $\langle \bar{\Phi}_{31} \rangle \neq 0$ and $\langle D_7 \rangle \neq 0$.

One also has to consider the $SU(5) \times U(1)$, $SU(4)$, and $SO(10)$ flatness conditions. The first one implies $|V_2| = |\bar{V}_5|$, which is satisfied automatically when the D'' condition is satisfied [Eq. (12)]. The $SU(4)$ and $SO(10)$ conditions are given by [26]

$$D_{SU(4)} : \sum_{a=3,7} D_a^* \tau^a D_a = 0;$$

$$D_{SO(10)} : \sum_{a=1,2,3} T_a^* \lambda^A T_a = 0, \quad (17)$$

where τ^a, λ^A are the generators of $SU(4)$ and $SO(10)$, respectively. In the basis in which the generators are given by the imaginary antisymmetric matrices, the D_a (T_a) are represented by 6-vectors (10-vectors). Specific choices for the

vevs of these representations will determine the pattern of SU(4) and SO(10) symmetry breaking. It is worth noticing and easy to verify that these flatness conditions are automatically satisfied for D_a, T_a vevs that are real [26]. However such restriction is not necessary; relaxing it allows situations where, e.g., $|D|^2 = D \cdot D^* \neq 0$, while $D \cdot D = 0$.

The standard F -flatness conditions require that $\partial W / \partial \phi_a = 0$ at the minimum of the scalar potential. Considering only the cubic contributions to W [22], we obtain

$$\frac{\partial W}{\partial \Phi_{31}} = T_1 \cdot T_1 + T_3 \cdot T_3 + D_3 \cdot D_3, \quad \frac{\partial W}{\partial T_1} = T_1 \Phi_{31},$$

$$\frac{\partial W}{\partial D_3} = D_3 \Phi_{31},$$

$$\frac{\partial W}{\partial \bar{\Phi}_{31}} = D_7 \cdot D_7, \quad \frac{\partial W}{\partial T_3} = T_3 \Phi_{31}, \quad \frac{\partial W}{\partial D_7} = D_7 \bar{\Phi}_{31}. \quad (18)$$

As such, the $\partial W / \partial D_7$ condition appears to be in conflict with the D -flatness requirement of $\bar{\Phi}_{31}$, $D_7 \neq 0$. Inclusion of nonrenormalizable contributions to the superpotential may resolve this impasse.

In the standard string unification picture, one would expect the new gauge coupling $g_{Z'}$ to evolve from low energies up to the string scale, where it will meet the other gauge couplings. In our case the proper normalization of U' is well defined from the requirement that all massless fields have conformal dimension 1: $U' \rightarrow U' / \sqrt{3}$. The gauge coupling evolves according to the beta function $b' = \frac{1}{3} \text{Tr} (Q')^2$, where Q' are the charges that appear in Tables II–V. Defining the MSSM matter content as consisting of the quark and lepton fields in $F_{0,3,4}, \bar{F}_{2,3,5}, \mathcal{L}_{2,3,5}^c$ and the Higgs doublets in h_1, \bar{h}_{45} [22], we find $b' = \frac{16}{3}$, which is smaller than the traditional $b_Y = \frac{33}{5}$. At the scale where the intermediate F_1, \bar{F}_4 states become excited, we find $b' = \frac{41}{6}$, which is very close to b_Y . Therefore at least up to the unification scale the evolution of U' is reasonable. The evolution up to the string scale requires a detailed understanding of the spectrum at such mass scales.

One also needs to worry about the U' symmetry breaking mechanism. Simply assuming that the usual Higgs doublets, which are charged under U' , will effect the symmetry breaking will lead to a much-too-large mixing angle θ [4]. One must resort to a singlet field acquiring a vev radiatively [18], much in the same spirit as in the original flipped SU(5) picture [27]. This possibility looks promising, as in this scenario there are many singlet fields that do not acquire vevs at the string scale. In this case one expects $\delta \sim \theta \sim M_Z^2 / M_{Z'}^2$, and the phenomenological requirement (from fits to electroweak data, including R_b) of $\delta \sim 10^{-2}$ naively implies $M_{Z'} \sim 10 M_Z$, although order 1 factors may alter this relation either way.

V. EXPERIMENTAL PROSPECTS

In this section we address some of the experimental consequences of the Z' gauge bosons proposed above. We concentrate on the second scenario in Sec. III [case (11) in Eq. (6)], as this is the one that may shift R_b in an interesting way. [The predicted correlation between $\Delta \Gamma_{\text{had}}$ and ΔR_b has been

given in Eq. (7) above.] This analysis will be sketchy. More precise calculations will be warranted once some as-yet-unknown parameters (i.e., $\theta, g_{Z'}, M_{Z'}$) become available. In this spirit we neglect the effects of Z - Z' mixing and concentrate on the direct production and decay of our leptophobic Z' at hadron colliders. The only variables in this case will be the strength of the U' gauge coupling ($g_{Z'}$) and the mass of the Z' .

A. Z' width and branching ratios

The Z' partial width into quarks can be expressed as [9]

$$\Gamma(Z' \rightarrow q\bar{q}) = \frac{\Gamma_0}{M_Z} \left(\frac{g_{Z'}}{g_Z} \right)^2 \left(\frac{1}{3} \right) [C_V'^2 + C_A'^2] M_{Z'}, \quad (19)$$

where $\Gamma_0 / M_Z = G_F M_Z^2 / (2\sqrt{2}\pi) \approx 0.011$, and the C_V', C_A' couplings can be obtained from Eq. (3) for the preferred charge assignment in Eq. (6) [i.e., case (11)]. The explicit factor of $1/3$ in Eq. (19) has been inserted to normalize the U' charges properly, as discussed in Sec. IV. The relevant couplings for the case of interest are:

	C_V'	C_A'	$B(Z' \rightarrow q\bar{q})$
u	$-\frac{1}{2}$	$-\frac{1}{2}$	$\frac{1}{18}$
d	-1	0	$\frac{1}{9}$
c	1	1	$\frac{2}{9}$
s	2	0	$\frac{4}{9}$
t	$-\frac{1}{2}$	$-\frac{1}{2}$	$\frac{1}{18}$
b	1	0	$\frac{1}{9}$

(20)

The branching ratios have been calculated by neglecting the top-quark mass (i.e., for $M_{Z'} \gg 2m_t$); they are only slightly increased if the top-quark does not contribute fully. Note the strong preference for light-quark decays. The total Z' width, neglecting the top-quark mass, is given by

$$\frac{\Gamma_{Z'}}{M_{Z'}} \approx 0.033 \left(\frac{g_{Z'}}{g_Z} \right)^2, \quad (21)$$

which shows that our leptophobic Z' is expected to be narrow (assuming that $g_{Z'} \sim g_Z$), and, therefore, amenable to the standard searches for Z' to dijets at hadron colliders [14,15]. (If supersymmetric particle decays are allowed, the width will increase although it still will remain relatively narrow.)

B. Z' production and experimental limits

Direct channel production of a narrow Z' at hadron colliders has a parton-level cross section given by [9]

$$\hat{\sigma}(q\bar{q} \rightarrow Z') = K \frac{4\pi^2}{3} \frac{\Gamma_0}{M_Z} \left(\frac{g_{Z'}}{g_Z} \right)^2 \left(\frac{1}{3} \right) [C_V'^2 + C_A'^2] \times \delta(\hat{s} - M_{Z'}^2), \quad (22)$$

TABLE IV. The hidden SO(10) decaplets ($\mathbf{10}$) T_i fields and their transformation properties under $G_{U(1)}$. Also indicated are the charges under U_A and three orthogonal linear combinations of interest (U', U'', U''').

	U_1	U_2	U_3	U_4	U_5	U_A	U'	U''	U'''
T_1	$-\frac{1}{2}$	0	$\frac{1}{2}$	0	0	$\frac{5}{2}$	0	0	$\frac{1}{2}$
T_2	$-\frac{1}{2}$	$-\frac{1}{2}$	0	0	$-\frac{1}{2}$	$\frac{5}{2}$	0	0	-1
T_3	$-\frac{1}{2}$	0	$\frac{1}{2}$	0	0	$\frac{5}{2}$	0	0	$\frac{1}{2}$

with an estimated K factor of $K \approx 1.3$ [9]. For our present purposes it should suffice to determine the ratio of our Z' cross section to that obtained assuming a Z' with standard model couplings to quarks (as traditionally assumed in the experimental literature). These ratios are

$$\frac{\hat{\sigma}(u\bar{u} \rightarrow Z')}{\hat{\sigma}(u\bar{u} \rightarrow Z')_{SM}} \approx 0.58 \left(\frac{g_{Z'}}{g_Z} \right)^2,$$

$$\frac{\hat{\sigma}(d\bar{d} \rightarrow Z')}{\hat{\sigma}(d\bar{d} \rightarrow Z')_{SM}} \approx 0.90 \left(\frac{g_{Z'}}{g_Z} \right)^2. \quad (23)$$

To estimate the impact of present experimental searches for Z' into dijets, one could average the contributions from up and down quarks (which gives 0.75 average coefficient) and then multiply by $B(Z' \rightarrow jj)/B(Z' \rightarrow jj)_{SM} \approx 1.4$,

$$\sigma(p\bar{p} \rightarrow Z' \rightarrow jj) \approx \left(\frac{g_{Z'}}{g_Z} \right)^2 \sigma(p\bar{p} \rightarrow Z' \rightarrow jj)_{SM}. \quad (24)$$

Assuming $g_{Z'} = g_Z$, one can study the exclusion plots obtained by UA2 [14] and CDF [15] in their searches for $Z' \rightarrow jj$. We conclude that the UA2 lower bound of $M_{Z'} > 260$ GeV is applicable, whereas CDF does not impose any further constraints. We note that even the UA2 lower bound may be easily evaded should $g_{Z'} < g_Z$.

C. Z' effects on top-quark production

Our leptophobic Z' will contribute to the top-quark cross section at the Tevatron. The tree-level $q\bar{q}$ parton-level cross sections are given by [7]

$$\hat{\sigma}(q\bar{q} \rightarrow g \rightarrow t\bar{t}) = \frac{4\pi\alpha_s^2}{27\hat{s}} \beta (3 - \beta^2), \quad (25)$$

$$\hat{\sigma}(q\bar{q} \rightarrow Z' \rightarrow t\bar{t}) = \frac{4\pi}{3} \left(\frac{\Gamma_0}{M_{Z'}} \right)^2 \left(\frac{g_{Z'}}{g_Z} \right)^4$$

$$\times \frac{\hat{s}}{(\hat{s} - M_{Z'}^2)^2 + (\hat{s}\Gamma_{Z'}/M_{Z'})^2} \left(\frac{1}{3} \right)^2 [C_{V_q}'^2 + C_{A_q}'^2]$$

$$\times \left[\frac{\beta}{2} (3 - \beta^2) C_{V_t}'^2 + \beta^3 C_{A_t}'^2 \right], \quad (26)$$

where $\beta^2 = 1 - 4m_t^2/\hat{s}$. Since $\alpha_s/(\Gamma_0/M_{Z'}) \approx 10$, the Z' contribution is usually negligible compared with the QCD con-

TABLE V. The hidden SU(4) fields and their transformation properties under $G_{U(1)}$. D_i represent sextets ($\mathbf{6}$), whereas \tilde{F}_i, \tilde{F}_i represent tetraplets ($\mathbf{4}, \bar{\mathbf{4}}$). Also indicated are the charges under U_A and three orthogonal linear combinations of interest (U', U'', U''').

	U_1	U_2	U_3	U_4	U_5	U_A	U'	U''	U'''
D_1	0	$-\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	0	$\frac{3}{2}$	$\frac{1}{2}$	2	0
D_2	0	$-\frac{1}{2}$	$\frac{1}{2}$	$-\frac{1}{2}$	0	$\frac{3}{2}$	$\frac{1}{2}$	2	0
D_3	$-\frac{1}{2}$	0	$\frac{1}{2}$	0	0	$\frac{5}{2}$	0	0	$\frac{1}{2}$
D_4	$-\frac{1}{2}$	$-\frac{1}{2}$	0	0	$\frac{3}{2}$	$\frac{3}{2}$	-1	2	0
D_5	0	$-\frac{1}{2}$	$\frac{1}{2}$	0	0	$\frac{3}{2}$	$\frac{1}{2}$	2	0
D_6	0	$\frac{1}{2}$	$-\frac{1}{2}$	0	0	$-\frac{3}{2}$	$-\frac{1}{2}$	-2	0
D_7	$\frac{1}{2}$	0	$-\frac{1}{2}$	0	0	$-\frac{5}{2}$	0	0	$-\frac{1}{2}$
\tilde{F}_1	$-\frac{1}{4}$	$\frac{1}{4}$	$-\frac{1}{4}$	0	$-\frac{1}{2}$	$\frac{1}{2}$	0	$-\frac{9}{4}$	$-\frac{1}{2}$
\tilde{F}_2	$\frac{1}{4}$	$\frac{1}{4}$	$-\frac{1}{4}$	0	$\frac{1}{2}$	-2	$-\frac{1}{2}$	$\frac{1}{4}$	$\frac{1}{2}$
\tilde{F}_3	$\frac{1}{4}$	$-\frac{1}{4}$	$-\frac{1}{4}$	0	$\frac{1}{2}$	$-\frac{3}{2}$	$-\frac{1}{2}$	$\frac{7}{4}$	0
\tilde{F}_4	$-\frac{1}{4}$	$\frac{3}{4}$	$\frac{1}{4}$	0	0	$\frac{1}{2}$	0	$-\frac{9}{4}$	1
\tilde{F}_5	$-\frac{1}{4}$	$\frac{1}{4}$	$-\frac{1}{4}$	0	$\frac{1}{2}$	$-\frac{1}{2}$	-1	$-\frac{1}{4}$	$\frac{1}{2}$
\tilde{F}_6	$-\frac{1}{4}$	$\frac{1}{4}$	$-\frac{1}{4}$	0	$-\frac{1}{2}$	$\frac{1}{2}$	0	$-\frac{9}{4}$	$-\frac{1}{2}$
\tilde{F}_1	$-\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{2}$	$-\frac{1}{2}$	$\frac{3}{2}$	$\frac{1}{2}$	$-\frac{7}{4}$	0
\tilde{F}_2	$-\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{4}$	$-\frac{1}{2}$	$-\frac{1}{2}$	$\frac{3}{2}$	$\frac{1}{2}$	$-\frac{7}{4}$	0
\tilde{F}_3	$\frac{1}{4}$	$-\frac{1}{4}$	$\frac{1}{4}$	0	$-\frac{1}{2}$	$\frac{1}{2}$	1	$\frac{1}{4}$	$-\frac{1}{2}$
\tilde{F}_4	$-\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{4}$	0	$-\frac{1}{2}$	$\frac{3}{2}$	$\frac{1}{2}$	$-\frac{7}{4}$	0
\tilde{F}_5	$-\frac{1}{4}$	$-\frac{1}{4}$	$\frac{1}{4}$	0	$-\frac{1}{2}$	2	$\frac{1}{2}$	$-\frac{1}{4}$	$-\frac{1}{2}$
\tilde{F}_6	$-\frac{3}{4}$	$\frac{1}{4}$	$-\frac{1}{4}$	0	0	$\frac{3}{2}$	-1	$-\frac{7}{4}$	0

tribution. There are two exceptions: (i) either the $C_{V,A}$ charges are large, or (ii) resonant production (i.e., $\hat{s} \approx M_{Z'}^2$) is possible. The model of Ref. [3] satisfies the first exception but not the second one, as they fix $M_{Z'} = 1$ TeV, while at the Tevatron top-quark production occurs near threshold (i.e., $\sqrt{\hat{s}} \sim 2m_t$). Our model does not satisfy the first exception, but it may yield observable effects by satisfying the second exception if $M_{Z'} \sim 400$ GeV. This expectation has been verified explicitly by integrating $\hat{\sigma}$ over the parton distribution functions (taken from Ref. [28]). The result for $\sqrt{s} = 1.8$ TeV, $m_t = 175$ GeV, and $g_{Z'}/g_Z = 1$ is shown in Fig. 3 as a function of $M_{Z'}$ for the charge assignment of interest; a K factor of $K = 1.3$ has been applied [7]. Current experimental precision on $\sigma_{t\bar{t}}$ does not allow one as yet to exclude conclusively any such (~ 0.5 pb) shifts. For reference, the corresponding standard model cross section is $\sigma_{t\bar{t}}^{SM} = 4.75 \pm 0.6$ pb [29], showing that the Z' effects may exceed the present theoretical uncertainty in the standard model prediction and, therefore, be in principle discernible in

future runs at the Tevatron. Of course, the possible Z' effects on the top-quark cross section are constrained by the allowed values of $M_{Z'}$ and $g_{Z'}/g_Z$ that result from the usual $Z' \rightarrow jj$ searches discussed above.

D. Parity-violating spin asymmetries at RHIC

A novel way in which our leptophobic Z' may be detected takes advantage of its maximal parity-violating couplings to up-type quarks (and parity-conserving couplings to down-type quarks). This property ($C_L^u = c, C_R^u = 0; C_L^d = C_R^d = c$) depends solely on the underlying flipped SU(5) symmetry, and is not affected by the details of the string model. The RHIC Spin Collaboration [30] intends to measure various spin asymmetries when RHIC starts colliding polarized protons at center-of-mass energies as large as $\sqrt{s} = 500$ GeV, with an integrated luminosity in excess of 1 fb^{-1} . The spin asymmetry of greatest sensitivity in our case is defined by [31]

$$A_{\text{LL}}^{\text{PV}} = \frac{d\sigma_{--} - d\sigma_{++}}{d\sigma_{--} + d\sigma_{++}}, \quad (27)$$

where $d\sigma_{\lambda_1\lambda_2}$ represents the cross section for scattering of protons of given helicities producing a jet ($p^{\lambda_1}p^{\lambda_2} \rightarrow j + X$). It has been shown that at RHIC energies the dominant contribution to the asymmetry comes from the interference between gluon exchange and Z' exchange in the t -channel scattering of quarks of the same flavor. The parton-level asymmetry for a given quark flavor is then proportional to [16]

$$T_{gZ'}^{--} - T_{gZ'}^{++} = (C_L'^2 - C_R'^2) \frac{8}{9} \alpha_s \alpha_{Z'} \hat{s}^2 \times \text{Re} \left(\frac{1}{\hat{t} D_{Z'}(\hat{t})} + \frac{1}{\hat{u} D_{Z'}(\hat{u})} \right), \quad (28)$$

where $\alpha_{Z'} = g_{Z'}^2/4\pi$ and $D_{Z'}(x) = x - M_{Z'}^2 + iM_{Z'}\Gamma_{Z'}$. It is then clear that down-type quarks will not contribute to the asymmetry (as they have $C_L' = C_R'$), whereas up-type quarks will contribute maximally, yielding (as $\hat{t}, \hat{u} < 0$)

$$A_{\text{LL}}^{\text{PV}} > 0, \quad (29)$$

once the integration over polarized quark distribution functions is performed. The standard model QCD-electroweak contribution to $A_{\text{LL}}^{\text{PV}}$ is also positive. Examples of observable parity-violating spin asymmetries have been displayed in Ref. [16] for the case of $M_{Z'} = 1$ TeV, and require rather large couplings (as proposed in Ref. [3]). Such large couplings are not available in our present model, but an observable asymmetry may still be present for lighter Z' masses. Examining Eq. (28), it appears that for $M_{Z'} \gg \sqrt{\hat{s}} \sim 100$ GeV one obtains an $\sim 1/M_{Z'}^2$ dependence on the amplitude (and, therefore, on the asymmetry).

The above qualitative analysis should motivate detailed studies of $A_{\text{LL}}^{\text{PV}}$ at RHIC. Of particular interest should be the case of a leptophobic Z' in the context of flipped SU(5), where the only unknowns are $M_{Z'}$ and the product $(C_L^u g_{Z'})^2$ (which appears as an overall constant), and the

sign of the asymmetry is predicted to be positive. Note also that the Z - Z' mixing considered above plays no role in the prediction for $A_{\text{LL}}^{\text{PV}}$ and, therefore, even if such mixing is eventually found not to be relevant, parity-violating spin asymmetries at RHIC would still be of great interest in probing models of physics at very high energies.

VI. CONCLUSIONS

We have addressed the possible existence of leptophobic Z' gauge bosons in consistent unified theories, in particular in the context of an underlying flipped SU(5) gauge group. Leptophobia in this case is natural, as quarks are largely split from leptons in the SU(5) representations. In contrast, traditional unified gauge groups seem to require a dynamical mechanism for generating leptophobia at the electroweak scale. Our leptophobic Z' possesses distinct couplings to quarks, violating parity maximally in the up-quark sector, and not at all in the down-quark sector. Moreover, string-based Z' charge assignments lead to scenarios where R_b is shifted in the direction indicated experimentally, while keeping Γ_{had} and R_c essentially unchanged. We have considered the origins of such phenomenologically desirable Z' in the context of string-derived flipped SU(5), and identified three possible scenarios. One of them is particularly compelling, and has been studied further.

We have also determined some basic properties of our leptophobic Z' gauge boson, such as its total width, branching ratios, and production cross section. Current experimental limits from Z' to dijet searches may be applicable if the U' coupling is comparable to the weak coupling. We have also studied the effect of Z' exchange on the top-quark cross section. Also, the parity-violating couplings to up-type quarks have the potential of yielding observable spin asymmetries in polarized pp scattering at RHIC.

It is important to realize that even if Z - Z' mixing is found

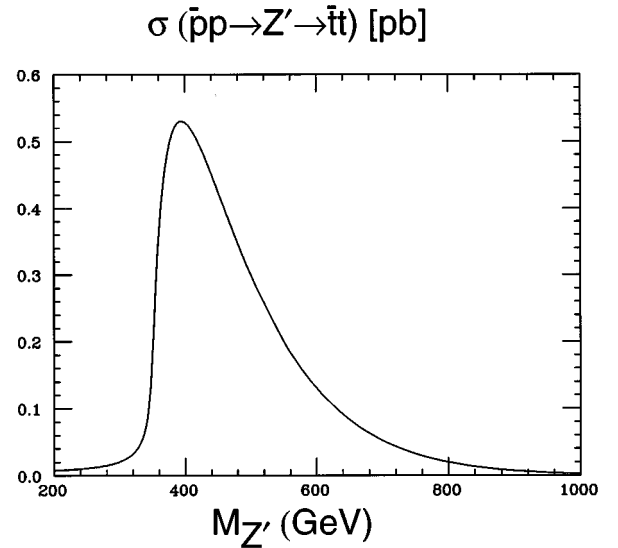


FIG. 3. The calculated cross section $\sigma(p\bar{p} \rightarrow Z' \rightarrow t\bar{t})$ versus $M_{Z'}$ (for $g_{Z'}/g_Z = 1$ and $m_t = 175$ GeV) at the Tevatron for the charge assignment of interest. [The cross-section scales with $(g_{Z'}/g_Z)^4$.] For reference $\sigma_{t\bar{t}}^{\text{SM}} = 4.75 \pm 0.6$ pb.

not to be relevant, leptophobic Z' gauge bosons may still be predicted by string models (unmixed or negligible mixed with the Z), and their existence should be probed experimentally in all possible ways. In the context of flipped SU(5), spin asymmetries at RHIC may be particularly sensitive probes.

The next step along these lines should include a more detailed study of the preferred string scenario, including non-normalizable terms in the superpotential and soft-supersymmetry-breaking terms. When a more clear picture of the effective theory below the string scale emerges, one should attempt a full dynamical evolution of the model down to low energies, paying particular attention to the radiative breaking of the U' symmetry via vevs of singlet fields. This crucial step will determine $M_{Z'}$, $g_{Z'}$, and θ , as well as the

spectrum of the new light degrees of freedom that accompany the singlet field.

Note added in proof. The calculation of A_{LL}^{PV} advocated in Sec. V D above has since been performed in Ref. [32], where it is shown that such asymmetry may be experimentally discernible from the corresponding standard model contribution for $M_{Z'} \lesssim 1$ Tev.

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