

## Fitting correlated data

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We discuss fitting correlated data, with the example of hadron mass spectroscopy in mind. The main conclusion is that the method of minimizing the correlated  $\chi^2$  is unreliable if the data sample is too small.

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### I. SPECIFICATION OF THE PROBLEM

We have  $N$  samples of unbiased estimators of quantities  $x_i$  with  $i = 1, \dots, D$ . Thus the data set is  $x_i(n)$  where  $n = 1, \dots, N$ . We assume that the samples  $x_i(n)$  are statistically independent versus  $n$  for fixed  $i$  but may be correlated in  $i$ . Such a situation arises in lattice gauge theory calculations where there are  $N$  independent configurations and  $D$  Green functions (linear combinations of Wilson loops or propagators) are measured versus time separation  $i$ . An introduction to this topic in the context of lattice gauge theory is provided by Toussaint [1].

The aim is to fit a given function  $F_i$  which depends on  $P$  parameters  $a_p$ . This function is to be fitted to the data samples  $x_i$ . Thus we need to find the best values of the parameters  $a_p$ , the errors associated with these best fit parameters, and the probability that the fit represents the data sample.

### II. REPRESENTATION OF THE PROBABILITY DISTRIBUTION OF THE SAMPLE

The data samples themselves give a probability distribution

$$S(x) = \frac{1}{N} \sum_{n=1}^N \delta^D(x - x(n)) .$$

We shall be interested in estimates of the probability distribution of the averages  $X_i$  of the data  $x_i$ . The most general way to achieve this is to fold the above distribution  $N$  times:

$$Q_S(X) = \int \delta^D \left( X - \frac{1}{N} \sum_{r=1}^N x^{(r)} \right) \prod_{s=1}^N d^D x^{(s)} S(x^{(s)}) .$$

Evaluating this distribution by simulation corresponds to the bootstrap method: many samples of  $N$  data are obtained by choosing from the original  $N$  possibilities randomly (with repetition allowed). In the limit of many such samples this corresponds to the above distribution  $Q_S$ . Such a procedure is in general inadequate for determining best fit parameters since a smooth representation of  $Q_S(X)$  is needed. An even more difficult task is usually to estimate the acceptability of such a fit. Thus the best fit parameters will yield  $X(f)$  and one must estimate the probability of such a value arising stochastically. This

also needs a smooth model of the distribution  $Q_S$ .

The natural interpolation for  $Q_S$  is suggested by the central limit theorem. Provided the underlying distributions of  $x_i$  are sufficiently localised, then for large  $N$ ,  $X_i$  will be Gaussianly distributed. We are specifically interested in the case where the different components  $x_i$  are statistically correlated. Thus a general Gaussian surface will be needed:

$$Q_G(X) = H \exp \left[ -\frac{1}{2} (X_i - \bar{X}_i) M_{ij} (X_j - \bar{X}_j) \right] .$$

Equating the first and second moments of these two expressions leads to the well-known identifications

$$\bar{X}_i = \frac{1}{N} \sum_{n=1}^N x_i ,$$

$$M_{ij} = N C_{ij}^{-1} ,$$

where

$$C_{ij} = \frac{1}{N-1} \sum_{n=1}^N (x_i - \bar{X}_i)(x_j - \bar{X}_j) .$$

To find the best fit parameters then corresponds to maximizing

$$\exp(-\chi^2/2)$$

where

$$\chi^2 = [F_i(a) - \bar{X}_i] M_{ij} [F_j(a) - \bar{X}_j]$$

with respect to  $a_p$  for  $p = 1, \dots, P$ . This is the usual correlated  $\chi^2$  method.

Diagonalizing the real symmetric positive-definite matrix  $C$  then allows us to write

$$\begin{aligned} Q_G &= H \exp(-\chi^2/2) = H \exp \left( -\sum_k^D \chi_k^2/2 \right) \\ &= H \exp \left( -\sum_k^D (Y_k - \bar{Y}_k)^2/2 \right) \end{aligned}$$

where

$$C_{ij} = R_{ki}\lambda_k R_{kj}, \quad \bar{Y}_j = R_{ji}\bar{X}_i N^{\frac{1}{2}}/\lambda_i^{\frac{1}{2}},$$

and

$$Y_j = R_{ji}X_i N^{\frac{1}{2}}/\lambda_i^{\frac{1}{2}}.$$

Consider, for example, a zero-parameter fit to  $X_i = F_i$  where  $F_i$  is the true value (the average over many samples). Then for one sample, the maximum is at  $X_i = \bar{X}_i \neq F_i$  and one can estimate the expected value of  $\chi^2$ .

Now letting  $\mathcal{F}_i = R_{ij}F_j N^{\frac{1}{2}}/\lambda_j^{\frac{1}{2}}$ , the probability distribution  $Q_G$  is a sum of  $D$  independent variables ( $Y_k - \mathcal{F}_k$ ) with unit variance. Thus  $\langle \chi^2 \rangle = D$  is expected. As we shall see later, there may be quite large corrections to this estimate in practice.

For a general fit  $X_i = F_i(a)$  with  $P$  parameters  $a$ , it is useful to regard this as a constraint on the uncorrelated variables  $Y_k = \mathcal{F}_k(a)$ . The set of  $\mathcal{F}_k$  values as the parameters  $a$  vary will be a  $P$ -dimensional surface and the fit minimizes  $\chi^2$  within this surface. So, effectively,  $\chi_k^2$  is zero for  $P$  components and thus the expected value of  $\chi^2$  only comes from the remaining  $D - P$  dimensions or “degrees of freedom.” Thus in the simplest case one expects  $\chi^2$  to be  $D - P$ .

What has become well known in the lattice gauge theory community is that the correlated  $\chi^2$  approach is not a very stable method [2]. Examples can also be constructed which give counterintuitive results [3]. I shall illustrate this with a simple simulation.

### III. A SPHERICAL GAUSSIAN DISTRIBUTION

For illustration, consider a true probability distribution of data points

$$S_T(x) = J \exp(-x_i x_i / 2),$$

where  $i = 1, \dots, D$ . This is actually quite general since after a change of variables (translation, rotation, and scaling), a general Gaussian surface can be brought to this form. Then we take  $N$  samples from this distribution giving data  $x_i(n)$  for  $n = 1, \dots, N$ . We obtain the distribution  $Q_G$  for this sample. We can easily maximize this distribution and clearly the sample maxima will be at  $X_i = \bar{X}_i$ . However, the true maximum is at  $X_i = 0$ . Thus we can estimate the probability of a particular sample having a maximum value of  $X_i$ . As well as using the above “correlated  $\chi^2$ ” approach, we also use a diagonal approximation for comparison. Thus we can calibrate standard statistical tests (such as  $\chi^2$  per degree of freedom) with this simulation in which the exact result is known.

Consider first the diagonal  $\chi^2$  for orientation. Then we evaluate by simulation the average value of  $\chi^2$  for the fit to  $X_i = F_i = 0$ . This corresponds to a zero parameter fit and so it is only the goodness of fit that can be addressed. This is usually quantified in terms of the value of the exponent  $\chi^2$ . The model and analysis then treat each of the  $D$  directions independently, so we obtain a value of  $\chi^2/D$  independent of  $D$ . The value of  $\chi^2 = N\bar{X}_i^2/C_{ii}$

for each dimension. Now the averages of  $\bar{X}_i^2$  and  $C_{ii}$  are  $1/N$  and 1, respectively, with both these estimates being unbiased. Thus if fluctuations among samples for both these quantities are neglected, we obtain an average of  $\chi^2/D$  of 1 as expected. However, their ratio will be a biased estimator, in general. For small  $N$  values, because the sample values fluctuate quite widely [the variance of  $C_{ii}$  is  $2/(N - 1)$ ], the value of  $\chi^2 = \bar{X}_i^2/C_{ii}$  for each dimension is enhanced on average by the contribution from low values of  $C_{ii}$ . Indeed, for large  $N$ , analytic evaluation gives  $\chi^2/D = 1 + 2/N$ . Our simulation values are shown in Table I as the  $D = 1$  column and are consistent with this analytic estimate. Thus there can be a small bias in the expected value of  $\chi^2/D$  even for uncorrelated evaluations. As we shall see such biases become much stronger for correlated  $\chi^2$ .

The way to illustrate the problem is to look at the shape of the Gaussian surface given from the sample data. One way to visualize the shape of the sample distribution  $Q_G$  is to obtain the eigenvalues  $\lambda$  and eigenvectors of  $C$ . Small eigenvalues correspond to a narrow distribution in the direction of the appropriate eigenvector. The width in any direction of the true distribution is 1. What we find by simulation is that the smallest eigenvalue of  $C$  from a sample of  $N$  data points can be very much smaller than 1 for small  $N$ . Of course, for  $N = D$ , the points will be linearly degenerate and thus  $\lambda_{\min} = 0$ . The results are shown in Fig. 1 and Table II. One sees that for  $N = 55$  (a typical number of configurations in a hadron spectrum calculation) and for  $D = 7$  the smallest eigenvalue is around 0.5 on average. Thus the distribution is 70% narrower in the direction corresponding to that eigenvalue.

The consequences of such a deformation (i.e., that the underlying spherical distribution appears squashed in the sample) are considerable in attempting to fit to the sample distribution. For example, let us consider again the fit of the true value  $X_i = 0$  to the sample distribution  $Q_G$ . This corresponds to a zero-parameter fit and so it is only the goodness of fit that can be addressed:  $\chi^2$  is expected to have an average value of  $D$ . Our results are shown in Table I. For example, we find, for the correlated fit,  $\chi^2/D$  is 3.03 with  $D = 5$  and  $N = 10$ . Using the usual criterion of  $\chi^2$  per degree of freedom would then imply that there was a very small probability (about 1%) that  $X_i = 0$  was the maximum of the distribution, which is quite wrong. This large  $\chi^2$  value arises since in this case the average value of  $\chi^2$  involves the sum of the inverses of the eigen-

TABLE I. Correlated  $\chi^2/D$  for  $N$  data with dimension  $D$  (averages from 10000 samples).

$N$	$D = 1$	$D = 3$	$D = 5$	$D = 7$	$D = 10$	$D = 15$
10	1.29	1.83	3.03	9.16	$\infty$	$\infty$
20	1.12	1.25	1.44	1.74	2.38	6.45
30	1.07	1.15	1.27	1.38	1.59	2.22
40	1.05	1.13	1.17	1.25	1.39	1.70
50	1.04	1.08	1.13	1.21	1.30	1.49
100	1.02	1.03	1.06	1.09	1.12	1.19

TABLE II. Data points  $N$  needed for smallest eigenvalue 0.5 with dimension  $D$ .

$D$	$N$
3	16
5	35
7	55
10	86
15	140
20	196

values of  $C$ . A very small eigenvalue will have the effect of increasing  $\chi^2$  more than that obtained from the estimated value of 1 for each eigenvalue. Again this effect is of  $O(1/N)$  as for the correlated case above. An analytic calculation for large  $N$  gives  $\chi^2/D = 1 + (D+1)/N$ . The difference is that the coefficient now increases with  $D$ . Moreover, the expression becomes infinite at  $N = D$ .

This effect can be even more pronounced in making a constrained fit since the distortion of the surface may shift the best fit value as well as affecting the probability of the fit. We simulated such a fit by finding the most probable point such that

$$X_i = F_i(a) = a\delta_{i1}.$$

This is a one parameter fit and the true value of  $a$  is zero. For each sample, we find the  $a$  value giving the maximum probability  $Q_G$ . We can also estimate the error on this best fit parameter within each sample by looking at the second derivative of  $\chi^2$  with respect to  $a$  in the usual way [1]. For  $N = 10$  and  $D = 5$  for a correlated fit, we find the error on the parameter  $a$  to be 0.056 on average over many samples. However, the average over many samples of the best fit value of  $a$  itself is  $\langle a^2 \rangle = 0.20$ . Thus, in this case, the error ascribed to the best fit value from the study of the quadratic sample surface is too small — about one-half of that actually present. This is a warning that the severe distortion of the sample Gaussian surface can upset the naive error analysis.

For the same situation with  $N = 10$  and  $D = 5$  but an uncorrelated fit, we get  $\langle a^2 \rangle = 0.10$  both from the sec-

ond derivative error analysis within each sample and also from the observed distribution of the best fit over many samples. Thus a diagonal  $\chi^2$  fit gives quite consistent results. Moreover, the true error on  $a$  is larger from a correlated fit. This can be understood since the diagonal  $\chi^2$  fit makes a stronger assumption about the nature of the data set and so allows a tighter fit.

Furthermore, the acceptability of these fits can be obtained from the value of  $\chi^2/(D-1)$  itself. We find a much bigger value for a correlated fit (2.26) than for a diagonal fit (1.28). At face value, this would suggest that the correlated fit was unacceptable. This latter effect is just the same as that described before and is due to the bias from the sample fluctuations at small  $N$ .

A standard method of finding the error on a sample value of fit parameters  $a$  is to repeat the whole fit with many bootstrap samples derived from the original sample. This corresponds to using samples of the point distribution  $Q_S(X)$  introduced earlier. A small bias can be introduced by this procedure and this can be corrected as discussed in the Appendix. With the correlated  $\chi^2$  method, this is inherently dangerous since the fluctuations in the shape of the Gaussian surface can be exaggerated. For example, if  $N = 10$  and  $D = 5$ , a bootstrap sample of  $N$  may only have contributions from five data sets. Thus the smallest eigenvalue will be zero with drastic consequences. In contrast, the bootstrap method is a completely satisfactory method for finding the overall errors with a diagonal  $\chi^2$  fit. A jackknife method can also be used instead of the bootstrap method but it has no obvious advantages except less computation.

Thus we learn that a representation of the probability surface as a general Gaussian is rather unstable if the number of data points is insufficient. In general, we need  $N > D^2$  to have a reasonable description with no big distortion of the shape in any direction. This can be understood roughly from the fact that the  $D$ -dimensional surface is determined by  $D(D+1)/2$  real numbers. The distortion averaged over all directions is less than that in the worst direction. Thus the  $\chi^2$  estimates made above need  $N > 10(D+1)$  to be less than 10%. To avoid serious fluctuations in sample estimates, it is prudent to take  $N > \max(D^2, 10(D+1))$ .

If one has fewer data points than this, then either one can estimate the expected probability distributions by simulation as in the example above, or one can assume a simpler model of the surface. The simplest model is to treat the data as uncorrelated. Then one has effectively  $D$  one-dimensional situations. Put another way, the minimum width is now to be taken along the prescribed  $D$  axes not along any arbitrary direction. For example, we use the same distribution introduced above but now with the different dimensions treated as uncorrelated. Then with  $D = 7$ ,  $N = 13$  data points is sufficient to have the smallest eigenvalue greater than 0.5 (to be compared with  $N = 55$  needed for correlated  $\chi^2$ ). Thus the shape is much more stable to fluctuations. Of course, this approach is only acceptable if there are indeed no statistically significant correlations in the data, as is the case in our example distribution.

A more realistic approach is to model such correla-

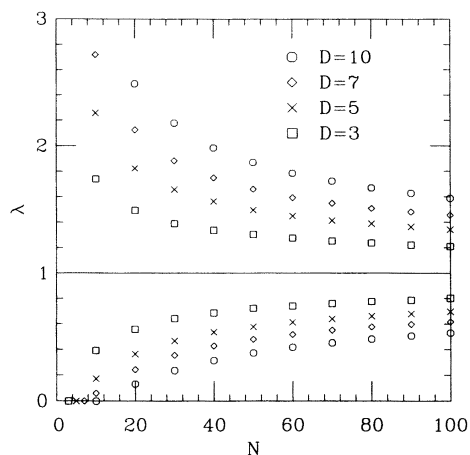


FIG. 1. The largest and smallest eigenvalues of  $C$  for  $N$  data samples in  $D$  dimensions.

tions by less parameters than a general quadratic — for instance, an overall common magnitude may represent the main effect. So putting  $y_i = x_i/x_1$  with  $y_1 = x_1$  may result in  $y$  being less correlated than  $x$ . Such an approach has the side effect that  $y_i$  are no longer unbiased estimators. To proceed, one makes a diagonal  $\chi^2$  fit to the data sample in the  $y$  variables. The error in the fit parameters is taken from a bootstrap analysis in which the original  $x$  data are bootstrapped. Then any bias can be accommodated as shown in the Appendix.

#### IV. CONCLUSION

Do not use correlated  $\chi^2$  with  $N$  data samples of  $D$  data unless  $N > \max(D^2, 10(D+1))$ . If you have insufficient data samples but there is statistically significant evidence for correlation among the data, then try to model any correlations among the data with less parameters. Even if some correlation among the data is suspected, it is reasonable to use an uncorrelated  $\chi^2$  fit but to estimate the errors on the parameters by an overall bootstrap of the fit rather than from the dependence of  $\chi^2$  on the parameters. The only drawback of such a procedure is that it may be difficult to estimate the goodness of fit since the correlations among the data may not have been adequately treated.

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#### APPENDIX: BOOTSTRAP BIASES

Here we discuss biases that can result from evaluating the average of some function of a quantity which is an unbiased estimator. We define an unbiased estimator as one for which the sample value  $\bar{x}$  averaged over many samples is the true average  $\hat{x}$ . So  $\langle \bar{x} \rangle = \hat{x}$ . We have in mind that  $x$  will be the sample average of some unbiased data and its distribution  $Q_S$  can be estimated by bootstrap as described in Sec. II.

Now if  $x$  is unbiased then  $f(x)$  will in general be biased. A simple illustration that explains the origin of this effect is that if  $x$  has a gaussian distribution with standard error  $\sigma$ , then  $(x \pm \sigma)^2 = x^2 + \sigma^2 \pm 2x\sigma$ . Hence the average of

$x^2$  is shifted systematically from the (average of  $x$ )<sup>2</sup>.

If the distribution of  $x$  is  $Q_S(x)$  with sample average  $\bar{x}$ , then the average of  $f$  over this distribution is

$$\bar{f} = \int dx f(x) Q_S(x) \approx f(\bar{x}) + \frac{1}{2} v f''(\bar{x}),$$

where  $v$  is the variance of the distribution  $Q_S$ . This is a rough estimate of the bias. Now the distribution  $Q_S(x)$  is itself only a sample and it has a sample mean  $\bar{x}$  which is distributed (in an unbiased way) around the true mean  $\hat{x}$  with distribution  $Q_T$  with variance  $v_T$ . To  $O(1/N)$ , the true variance  $v_T$  is equal to the sample variance  $v$ . The average over many samples gives

$$\langle f(\bar{x}) \rangle = \int dx f(\bar{x}) Q_T(\bar{x} - \hat{x}) \approx f(\hat{x}) + \frac{1}{2} v f''(\hat{x}).$$

Thus there is a further bias from the distribution of the sample mean. The key is that this bias is the same (to order  $1/N$ ) as that above: namely, the difference between the sample values of  $\bar{f}$  and  $f(\bar{x})$ . This suggests the unbiased estimator for  $f(\hat{x})$ :

$$f_U = 2f(\bar{x}) - \bar{f}.$$

Now this expression can be evaluated in a straightforward way by bootstrap simulation. The average of  $f(\bar{x})$  over the bootstrap samples gives  $\bar{f}$ , while  $f(\bar{x})$  is obtained by evaluating  $f$  with the sample average  $\bar{x}$ . Note that since  $x$  is an unbiased estimator, the average of  $\bar{x}$  over the bootstrap samples should be the same as the sample average in the limit of many bootstrap samples.

This technique is also valid for the parameters  $a$  of a fit to the data. The combination  $2a - \langle a \rangle_B$  will be unbiased where the angular brackets here imply an average of the  $a$  values obtained from fitting many bootstrap samples of the original data set.

The unbiased estimator can be thought of as the naive value  $f(\bar{x})$  with a systematic correction  $[f(\bar{x}) - \bar{f}]$ . This correction will itself have some error when evaluated from one data sample. An estimate of the error to be ascribed to this unbiased estimator can be obtained by a further bootstrap. One creates many first level bootstrap samples, and for each one calculates  $f_U$  by a further nested bootstrap. Then the variance of the values of  $f_U$  obtained is available. This is computationally demanding but straightforward in practice. This is analogous to the nested jackknife method proposed to deal with this same problem of bias [4].

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