

Two-jet invariant-mass distribution at $\sqrt{s}=1.8$ TeV

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(Received 2 November 1989)

We present the dijet invariant-mass distribution in the region between 60 and 500 GeV, measured in 1.8-TeV $\bar{p}p$ collisions in the Collider Detector at Fermilab. Jets are restricted to the pseudorapidity interval $|\eta| < 0.7$. Data are compared with QCD calculations; axigluons are excluded with 95% confidence in the region $120 < M_A < 210$ GeV for axigluon width $\Gamma_A = N\alpha_s M_A/6$, with $N = 5$.

The dijet invariant-mass spectrum is both a testing ground for QCD and the natural place to look for new massive objects which are strongly produced and decay into two jets. In this paper, we present the invariant-mass distribution of dijets produced in the process $\bar{p}p \rightarrow \text{jet}_1 + \text{jet}_2 + X$, and a limit on the production of axiglons,¹ based on a Collider Detector at Fermilab (CDF) exposure of 26 nb^{-1} .

The CDF is described in detail in Ref. 2. The components relevant to this analysis are the barrel-shaped central calorimeters which measure the electromagnetic and hadronic energy of jets. The inner electromagnetic calorimeter consists of layers of lead and scintillator, and has a resolution for electrons $\sigma/E = 13.5\%/\sqrt{E \sin \theta}$ (E in GeV). The outer hadron calorimeter is steel and scintillator, and has a resolution for pions $\sigma/E = 70\%/\sqrt{E}$. The calorimeters are segmented into projective towers, each subtending 0.1 units of η (pseudorapidity $\eta \equiv \ln[\cot(\theta/2)]$, with θ the polar angle with respect to the proton-beam direction) and 15° in ϕ (the azimuthal angle around the beam).

The data used in this analysis were taken with hardware triggers requiring the summed transverse energy ($E \sin \theta$) in the central calorimeters to be above 20, 30, 40, and 45 GeV, depending on the luminosity. To retain the projective geometry of the detector, and to ensure full containment of jets in the central calorimeter, the collision point for each event was required to lie within 50 cm of the detector center. About 17% of the data were rejected by this cut. Jets were identified according to the algorithm described in Ref. 3. Briefly, local depositions of energy are identified, and all calorimeter towers with transverse energy above 0.2 GeV in a cone of radius 0.7 in η - ϕ space are collected to form the jet. The observed jet energy and momentum are taken to be the scalar and vector sums of the calorimeter tower energies.

For the dijet event selection, the two highest- P_T jets in each event were required to be in the fiducial region $|\eta| < 0.7$. In addition, a cut was imposed on the azimuthal separation of the two jets, $\Delta\phi > 100^\circ$. Trigger bias was eliminated by using only events with dijet masses in the fully efficient region, which is above 60, 65, 80, and 90 GeV for the four different trigger samples. These requirements were arrived at by considering the E_T trigger thresholds and the allowed η interval, and are explained in Ref. 4.

Backgrounds from cosmic rays and accelerator losses were eliminated by requiring arrival times at the hadronic calorimeter within a 35-nsec window around the beam-beam interaction. Remaining backgrounds falling inside this timing window or deposits solely in the electromagnetic calorimeter, which has no timing information, were removed by the requirement that the electromagnetic fraction of the observed clusters lie in the region $0.05 < F_{EM} < 0.95$. The validity of these background suppression schemes was checked by scanning all events with central jets with $P_T > 100$ GeV. In 300 such events, one real jet event was lost, and two bad events (P_T ap-

proximately 100 GeV) were passed. These two background events were due to the *bremstrahlung* of cosmic rays inside the calorimeter.

The energy and momentum of each jet were corrected independently for losses due to cracks, leakage, energy out of the clustering cone, underlying event contribution, and nonlinear response of the calorimeter according to a detector simulation and the ISAJET Monte Carlo program.⁵ The ISAJET jet-fragmentation properties have been checked against charged-particle properties measured in the CDF jet data, and the calorimeter simulation has been tuned to the observed response for isolated charged pions in minimum-bias $\bar{p}p$ interactions in the CDF, and test beam data. As determined from the Monte Carlo program, the total loss of jet energy from all effects ranges from 30% for 30-GeV jets to 14% at 200 GeV. By varying the fragmentation parameters and the calorimeter response in the simulation, we estimate the uncertainty in the jet energy scale to be 9% at 30 GeV and 5% at 200 GeV. The dijet invariant mass M is defined as $\sqrt{(E_1 + E_2)^2 - (P_1 + P_2)^2}$, where E_i and P_i are the corrected energy and momenta of the two central jets. The dijet mass resolution $[\sigma(M)]$ as determined from the Monte Carlo simulation is shown in Fig. 1; the fractional resolution (σ/M) varies from 15% at 60 GeV to 10% at 400 GeV. Fits to two plausible functional forms for the resolution are also shown in Fig. 1. These are of the form $\sigma = \alpha\sqrt{M} + \beta M$ and $\sigma = \alpha\sqrt{M} + \beta$.

The resulting dijet invariant-mass spectrum was then corrected for the smearing (or feed-up) effect caused by the finite-mass resolution of the detector. The procedure was to convolve a test function of the form $Am^{-\alpha}e^{\beta m}$ with the dijet mass resolution and then fit to the data. For the best-fit parameters, the ratios of the unsmear-

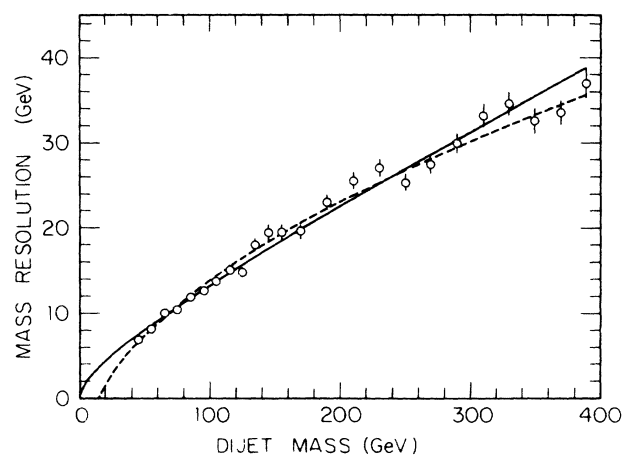


FIG. 1. The dijet mass resolution as determined from Monte Carlo simulation. A cluster cone size of 0.7 in η - ϕ and a tower E_T threshold of 0.2 GeV are used. Both jets were restricted to the fiducial region $|\eta| \leq 0.7$. Two fits are shown: $\sigma = 0.68\sqrt{M} + 0.065M$ (solid line) and $\sigma = 2.229\sqrt{M} - 8.40$ (dashed line).

to the smeared test functions, integrated over the data bins, were applied to the data. These correction factors range from 0.85 at 60 GeV to 0.9 above 200 GeV.

In Table I and Fig. 2 we present the corrected CDF dijet mass spectrum at $\sqrt{s} = 1800$ GeV, compared with a range of theoretical predictions.⁶ The result of UA2 at $\sqrt{s}=546$ GeV (Ref. 7) is also shown. In the figure, the uncertainties include the statistical uncertainties and the mass-dependent part of the systematic uncertainties. These systematic cross-section uncertainties, including the uncertainty in the jet energy scale (45% of the cross section at 60 GeV and 30% at 400 GeV), the resolution deconvolution (10%), and the integrated luminosity (15%), range from 48% at 60 GeV to 35% at 400 GeV. The solid lines are the envelopes of several QCD calculations with the structure functions DO 1, DO 2 (Ref. 8), EHLQ 1, EHLQ 2 (Ref. 9), and the momentum scales $Q = P_T$ and $P_T/2$. Our measurement and QCD are consistent within experimental and theoretical uncertainties. The effect of the contribution of hadronic decays of the W/Z bosons are not included in the calculation; their effect on the dijet mass spectrum is small at Fermilab Tevatron energies. The increasing contribution of gluons in this range of M_{jj} ; at higher energies tends to reduce the signal-to-noise ratio for detecting such decays, and with the limited statistics in this sample, a significant signal is not expected.

The dijet invariant-mass distribution, at both the theoretical and experimental level, is sensitive to the presence of additional jets in the event, through both event selection and definition of invariant mass. Restricting the sample to more “back-to-back” events ($\Delta\phi > 150^\circ$) lowered the measured cross-section roughly 20%. Merging

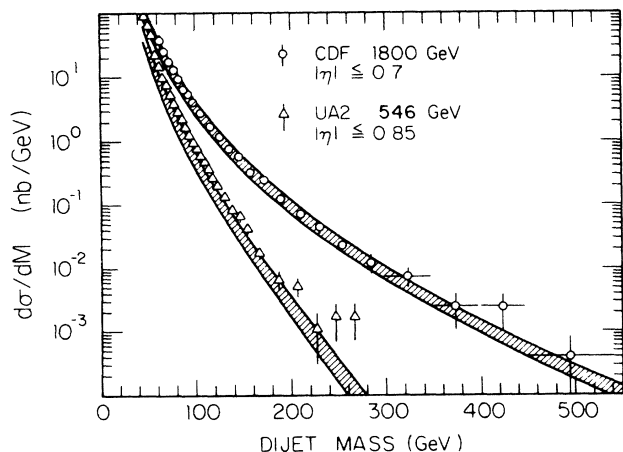


FIG. 2. The dijet mass spectra at $\sqrt{s} = 1800$ GeV (CDF) and $\sqrt{s} = 546$ GeV (UA2). The uncertainties include the statistical uncertainties and the mass-dependent systematic uncertainties. The additional CDF (UA2) mass-independent normalization uncertainty of 35% (45%) is shown in the key. The two pairs of solid lines are the envelopes of QCD calculations ($2 \rightarrow 2$) with the structure functions DO 1, DO 2, EHLQ 1, and EHLQ 2, and the momentum scales $Q = P_T/2$ and P_T .

TABLE I. The dijet mass spectrum. Jets are restricted to the pseudorapidity region $|\eta| < 0.7$.

| M_{JJ} (GeV) | $d\sigma/dM_{JJ}$ (nb/GeV) | Sys. error | Stat. error |
|-------------------|-------------------------------|---------------|------------------------|
| 62.5 | 37.8 | 18.5 | 4.4 |
| 67.5 | 25.5 | 12.0 | 0.62 |
| 72.5 | 17.7 | 8.0 | 0.52 |
| 77.5 | 13.2 | 5.8 | 0.45 |
| 82.5 | 9.39 | 4.0 | 0.31 |
| 87.5 | 6.46 | 2.70 | 0.26 |
| 92.5 | 5.44 | 2.22 | 0.21 |
| 97.5 | 4.17 | 1.68 | 0.18 |
| 105.5 | 2.83 | 1.11 | 0.11 |
| 115.5 | 1.72 | 0.65 | 0.083 |
| 125.5 | 1.17 | 0.43 | 0.068 |
| 135.5 | 0.756 | 0.275 | 0.055 |
| 145.5 | 0.569 | 0.204 | 0.048 |
| 157.5 | 0.321 | 0.113 | 0.029 |
| 172.5 | 0.247 | 0.085 | 0.026 |
| 189.5 | 0.121 | 0.041 | 0.016 |
| 210.5 | 0.0708 | 0.0237 | 0.0120 |
| 230.5 | 0.0445 | 0.0147 | +0.0117 -0.0076 |
| 254.5 | 0.0230 | 0.0075 | +0.0070 -0.0043 |
| 284.5 | 0.0122 | 0.0040 | +0.0056 -0.0029 |
| 323.5 | 0.00731 | 0.00237 | +0.00334 -0.00171 |
| 373.5 | 0.00244 | 0.00080 | +0.00237 -0.00074 |
| 423.5 | 0.00244 | 0.00081 | +0.00237 -0.00074 |
| 494.5 | 0.000404 | 0.000139 | +0.000929 -0.000118 |

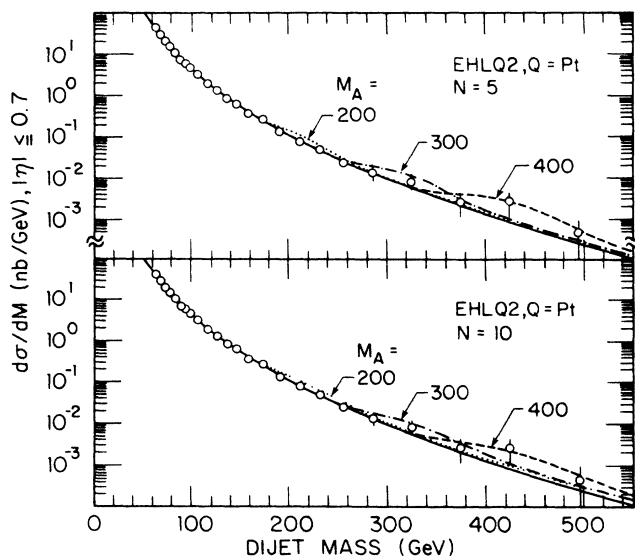


FIG. 3. The CDF dijet mass spectrum (uncorrected for mass resolution) compared with calculations involving the effect of axigluons, for two values of the axigluon width parameter, N . Predictions for axigluon masses of 200, 300, and 400 GeV are shown; these predictions have been normalized to the data in the region $60 < M_{JJ} < 120$ GeV.

nearby jets, within an η - ϕ radius of 1.2 (1.5) from the leading two jets, resulted in a cross section 30% (85%) higher. These effects, inherent to QCD,¹⁰ are not included in the systematic uncertainty reported for the cross section.

Recently, chiral models of QCD with the symmetry group $SU(3)_L \times SU(3)_R$ have been proposed¹ in which the symmetry is broken to $SU(3)_C$ and as a consequence the axial-vector gauge particles, axigluons, become massive. Bagger, Schmidt, and King¹¹ (BSK) have calculated the contribution of axigluons to the jet production cross sections assuming an axigluon width $\Gamma_A = N\alpha_s M_A/6$, with N being proportional to the number of decay channels (light quarks), α_s the strong coupling constant, and M_A the axigluon mass. They have ruled out the mass region $125 < M_A < 275$ GeV based on the UA1 jet P_T spectrum at $\sqrt{s} = 630$ GeV.¹² UA1 (Ref. 13) has excluded masses between 110 and 310 GeV for an axigluon width $\Gamma_A < 0.4 M_A$ (equivalent to $N = 24$ for $\alpha_s = 0.1$), using an incoherent sum of dijet and axigluon cross sections. We have repeated the BSK calculation at the Tevatron energy with EHLQ 2 and $Q = P_T$ using a coherent sum of amplitudes. EHLQ 2 was chosen because it gave the most conservative limits. For the axigluon analysis, we have compared to the data without the application of smearing corrections, and instead have added the effect of mass resolution to the theoretical predictions. The data and the calculations are shown in Fig. are excluded in the mass intervals $120 < M_A < 210$ GeV

3, where the predicted spectra have been normalized to the data in the region $60 < M_{JJ} < 120$ GeV.

The χ^2 for a fit including an axigluon is calculated for the three or four bins under the axigluon bump depending on the mass and width of the axigluon, after normalization to the low-mass region. To take into account the uncertainty in the jet energy scale, we have allowed both edges of each data bin to vary and have used the smallest χ^2 to set a limit. At the 95% confidence limit, axigluons for $N = 5$ ($\Gamma_A = 0.09M_A$), and $120 < M_A < 150$ GeV for $N = 10$ ($\Gamma_A = 0.18M_A$).

Although CDF data on dijet mass extends the range explored at the $Spp\bar{S}$, our sensitivity to axigluon masses below 300 GeV is lessened by the predominance of gluon-gluon and gluon-quark scattering in this region, and the fact that axigluon production proceeds via $q\bar{q}$ scattering. The higher statistics of the latest Tevatron run will extend our sensitivity to larger axigluon masses.

This work has been made possible by the construction and successful operation of the Tevatron collider by the Fermilab Accelerator Division. We thank J. Bagger, C. Schmidt, S. King, and E. Eichten for providing help with the theoretical calculations. This work was supported in part by the U.S. Department of Energy and National Science Foundation, the Italian Istituto Nazionale di Fisica Nucleare, the Ministry of Science, Culture, and Education of Japan, and the A. P. Sloan Foundation.

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