Without (4.26), (4.30) becomes

$$K'=k'\langle\mu'\rangle\ , \tag{2}$$

where the subscript $h'N$ on $\langle\mu'\rangle$ has been omitted. Combining (3.7) and (4.21) yields

$$K'=\alpha k'\langle\mu\rangle\ . \tag{3}$$

Thus from (2) and (3) we have

$$\alpha=\frac{\langle\mu'\rangle}{\langle\mu\rangle}=\frac{\int dR^2\Omega'(R)}{\left[\int dR^2\Omega(R)\right]\left[\int dR^2 g'(R)\right]}\ . \tag{4}$$

Now, the average multiplicity $\langle n\rangle_{hA}$ is given by (3.21), where $\beta=\alpha p$ is a parameter known empirically to be 0.5, as shown by extensive comparison with data in Sec. V. From (2.18), (4.3), (4.4), and (4.22), we have

$$p=\frac{\sigma_{\text{in}}^{h'N}}{\sigma_{\text{in}}^{hN}}=\int dR^2 g'(R)\ . \tag{5}$$

Consequently, it follows from (4) that

$$\beta=\frac{\int dR^2\Omega'(R)}{\int dR^2\Omega(R)}\ . \tag{6}$$

Adopting a simple form for $\Omega'(R)$ by assuming that it has the same R dependence as $\Omega(R)$ on the ground that the transverse dimension of h' is unchanged from h due to time dilation, we write

$$\Omega'(R)=c\Omega(R)\ . \tag{7}$$

Then c can be trivially determined by (6):

$$c=\beta=0.5\ . \tag{8}$$

The resultant form for $\Omega'(R)$ is slightly different from that given in (4.27), but the effect of the difference on the phenomenology done in Ref. 1 is negligible. More specifically, instead of $p=0.79$ as given in (4.29), we now have $p=0.6$, and thus instead of $\alpha=0.63$ as given in (4.31), we now have $\alpha=0.83$, the product being invariant: $\alpha p=\beta=0.5$. It is β, not α and p separately, that determines the dominant terms in the moments $\langle n^r\rangle$. This can be seen in Eqs. (A18)–(A20) in Ref. 1, where, for example, the $\langle\nu(\nu-1)(\nu-2)\rangle$ term is far more important than $\langle\nu(\nu-1)\rangle$ and is weighted by β^3. Hence, the effect of our change in $\Omega'(R)$ on the multiplicity distributions plotted in Figs. 7–15 in Ref. 1 is imperceptible.

Subsequent work on nucleus-nucleus collisions[2] is based on the analysis presented here.

We are grateful to Dr. N. N. Nikolaev for discussions that led us to notice the error in (4.25).

[1] X. N. Wang and R. C. Hwa, Phys. Rev. D **39**, 2573 (1989).
[2] R. C. Hwa and X. N. Wang, Report No. OITS-410, 1989 (unpublished).

Erratum: Weak-boson production at Fermilab Tevatron energies [Phys. Rev. D 40, 83 (1989)]

E. L. Berger, F. Halzen, C. S. Kim, and S. Willenbrock

The vertical scale along the right-hand side of Fig. 7 was in error. The correct version is reproduced below.

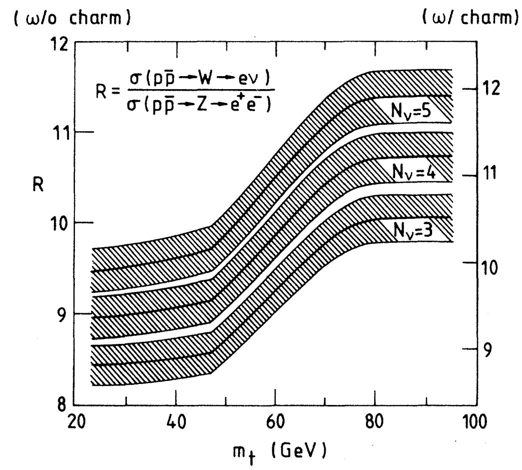

FIG. 7. The prediction of the ratio R of $W \rightarrow e\nu$ to $Z \rightarrow e^+e^-$ events at $\sqrt{s}=1.8$ TeV is plotted as a function of m_t for $N_\nu=3$, 4, and 5, with (right-hand scale) and without (left-hand scale) inclusion of contributions from the charm quark.

Erratum: Search for a neutral Higgs boson in B-meson decay [Phys. Rev. D 40, 712 (1989)]

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Equation (4) for the branching ratio $\Gamma(B \rightarrow H^0K)/\Gamma(B \rightarrow H^0X)$, plotted in Fig. 4, was incorrectly copied from Haber, Schwartz, and Snyder (Ref. 9). The factor $(1-M_H^2/\Lambda^2)$ in the denominator should be squared. This has the effect of making our Higgs-boson searches in the $B \rightarrow H^0K$ modes more sensitive than we assumed, especially at the higher values of M_H. Upper limits for $B \rightarrow H^0X$ from exclusive $B \rightarrow H^0K^-$ modes shown in Fig. 14 should be lowered by a factor of 0.89, 0.76, and 0.65 at $M_H=2$, 3, and 3.6 GeV, respectively. The corrected limits remain higher than the unchanged limit from inclusive $B \rightarrow (\mu^+\mu^-)X$ (solid curve in Fig. 14), except at $M_H=3.1$ GeV. The conclusions of our paper are therefore unchanged.

We are indebted to Gad Eilam for finding our mistake.