



FIG. 2. (a) Charged-pion pair production by bremsstrahlung in e^+e^- annihilation diagrams. The radiated virtual photon can be from any of the four charged lines. (b) Charged-pion pair production involving Compton amplitude with timelike photons in e^+e^- annihilation channel.

damping when s_0 is beyond the resonance region ($\sqrt{s_0} \approx 1.5$ GeV). We note that all the other diagrams [Fig. 1(b) and Fig. 2(a)] are described by the pion electromagnetic form factor and produce a charged-pion pair with opposite charge conjugation from that produced in Fig. 1(a). Thus, if the charge of the individual pions is not observed, there is no interference in the cross section between the two types of amplitudes. Furthermore, since the pion pair produced by bremsstrahlung is in a p -wave state, these diagrams give no contribution¹⁶ to (15) in the limit $s \rightarrow 4m_\pi^2$.

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¹⁶If s is not close to $4m_\pi^2$, a standard calculation shows that the p -wave contribution in our special kinematics is proportional to the electron mass squared and therefore negligible for the e^-e^- case [cf., F. Calogero and C. Zemach, Phys. Rev. **120**, 1860 (1960)]. These authors claim that this contribution vanishes not merely in the relativistic limit. Their arguments seem to have neglected the spin dependence which exists when the electron mass is not zero; it can be a significant background to (13) when $Q^2 \gtrsim$ for the e^-e^+ case, however.

Erratum

Theory of Ω^- Decay, D. N. Goswami and J. Schechter [Phys. Rev. D **1**, 290 (1970)]. Equation (A1) should read

$$\Gamma(\Omega \rightarrow B\pi) = \frac{(E'^2 - m'^2)^{3/2}}{12\pi m(\Omega)} [A^2(E' + m') + B^2(E' - m')]. \quad (\text{A1})$$

Also a parenthesis is missing in Eq. (A2), which should read

$$\Gamma(\Omega \rightarrow \Xi^*\pi) = \frac{(E'^2 - m'^2)^{1/2}}{4\pi m(\Omega)} \{ \dots \}. \quad (\text{A2})$$