Errata

Kinematics of Production Processes at High Energy and the Regge Pole Hypothesis, G. Luxton [Phys. Rev. D 2, 1926 (1970)]. Equation (3.1), which states that the dominant t-channel amplitude in the limit (2.2) is helicity conserving, does not follow from angular momentum conservation alone, but requires a separate assumption. Since the t-channel momenta tend to infinity, the transverse momentum components need not tend to zero, even in the limit of collinear momenta. Thus, the particles are permitted to have nonzero components of orbital angular momentum along the direction of collinearity.

I would like to thank Professor J. E. Mandula for several discussions that ultimately led to the discovery of this error.

Model for Electron Excitation of the Nucleon. II, P. L. Pritchett and J. D. Walecka [Phys. Rev. 168, 1638 (1968)]. In the discussion following Eq. (3.25), the values of the variational parameters should be $\rho_0 a = 6.1$ and $\epsilon_0 = 0.54$. The only appreciable effect is to change the values of f^2 in Table IV to 1.6, 0.3, and 0.3, respectively. The discussion of these results is unaltered.

Implications of Local Duality in a Set of Coupled Reactions, M. J. King and Kameshwar C. Wali [Phys. Rev. D 3, 1602 (1971)]. We wish to point out the following corrections. In both Eqs. (3.19) and (3.22), the minus signs between the two terms within the square brackets should be plus signs.

Electromagnetic Contribution to the Decays $K_S \to l\bar{l}$ and $K_L \to l\bar{l}$, L. M. Sehgal [Phys. Rev. 183, 1511 (1969)]. There is an error in Eq. (9) of this paper. This affects the result for $K_S \to l\bar{l}$ [Eq. (14b)] but not that for $K_L \to l\bar{l}$ [Eq. (14a)]. The final results of the paper should be

$$\frac{\text{Rate}(K_L \to l^+ l^-)}{\text{Rate}(K_L \to \gamma\gamma)} \geq \tfrac{1}{2}\alpha^2 \left(\frac{m}{M}\right)^2 \frac{1}{\beta}\left(\ln\frac{1+\beta}{1-\beta}\right)^2, \qquad (14a)$$

$$\frac{\text{Rate}(K_S \to l^+ l^-)}{\text{Rate}(K_S \to \gamma\gamma)} \geq \tfrac{1}{2}\alpha^2 \left(\frac{m}{M}\right)^2 \beta\left(\ln\frac{1+\beta}{1-\beta}\right)^2. \qquad (14b)$$

The second of these results is now in agreement with that of Martin, de Rafael, and Smith [Phys. Rev. D 2, 179 (1970)], provided that a missing factor of $\frac{1}{4}$ is supplied in Eq. (4.1) of their paper.

Some of the numbers in Table I are in error, and should be recomputed from the two equations given above.

I wish to thank Dr. C. Quigg and Dr. J. Smith for private communications.

Comments on Some Recent Muon Experiments, F. R. Hickey, P. J. McNulty, R. K. Shivpuri, and N. J. Wixon [Phys. Rev. D 3, 254 (1971)]. The symbol W on line 43, column 1, of p. 255, should have appeared as W^2.

Nonrelativistic Separable Potential Quark Model of the Hadrons, T. R. Mongan and Leon Kaufman [Phys. Rev. D 3, 1582 (1971)]. Equation (14) on page 1586 should read

$$\psi_{L=0}(\vec{x}_1, \vec{x}_2, \vec{x}_3) = N' \exp[-K^2(x_1{}^2 + x_2{}^2 + x_3{}^2)].$$

Study of $\pi^+ p$ Four-Prong Interactions from 2.95 to 4.08 GeV/c, David Brown, George Gidal, Robert W. Birge, Sun-Yiu Fung, Warren Jackson, and Robert T. Poe [Phys. Rev. D 1, 3053 (1970)]. A misprint occurred in the fifth and the sixth lines of the equation on p. 3064. The last three terms should each be multiplied by a factor $3\sqrt{3}$. The full equation should read

$$\begin{aligned} W(\theta_c, \varphi_c, \theta_d, \varphi_d) = (1/16\pi^2)\{&1 + \tfrac{1}{2}(1 - 3\rho_{00}^c)(1 - 3\cos^2\theta_c) - \tfrac{1}{2}(1 - 4\rho_{33}^d)(1 - 3\cos^2\theta_d) + R_8(1 - 3\cos^2\theta_c)(1 - 3\cos^2\theta_d) \\ &-3(\rho_{1,-1}^c \sin^2\theta_c \cos 2\varphi_c + \sqrt{2}\,\rho_{10}^c \sin 2\theta_c \cos\varphi_c) \\ &-2\sqrt{3}(\rho_{3,-1}^d \sin^2\theta_d \cos 2\varphi_d + \rho_{3,1}^d \sin 2\theta_d \cos\varphi_d) \\ &-3(1 - 3\cos^2\theta_d)[R_9 \sin^2\theta_c \cos 2\varphi_c + (1/\sqrt{2})R_{10} \sin 2\theta_c \cos\varphi_c] \\ &-\sqrt{3}(1 - 3\cos^2\theta_c)(R_{11} \sin^2\theta_d \cos 2\varphi_d + R_{12} \sin 2\theta_d \cos\varphi_d) \\ &+3\sqrt{3}(\sin^2\theta_c \sin^2\theta_d[R_{13}\cos(2\varphi_c + 2\varphi_d) + R_{14}\cos(2\varphi_c - 2\varphi_d)] \\ &\quad + \sin^2\theta_c \sin 2\theta_d[R_{15}\cos(2\varphi_c + \varphi_d) + R_{16}\cos(2\varphi_c - \varphi_d)] \\ &\quad + (1/\sqrt{2})\sin 2\theta_c \sin^2\theta_d[R_{17}\cos(\varphi_c + 2\varphi_d) + R_{18}\cos(\varphi_c - 2\varphi_d)] \\ &\quad + (1/\sqrt{2})\sin 2\theta_c \sin 2\theta_d[R_{19}\cos(\varphi_c + \varphi_d) + R_{20}\cos(\varphi_c - \varphi_d)])\}. \end{aligned}$$