

$b\bar{b}$ spectroscopy

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We comment on the significance of the experimental findings in $b\bar{b}$ spectroscopy at the Cornell Electron Storage Ring, and provide new theoretical results for $E1$ transitions in $b\bar{b}$ based upon our quantum-chromodynamic potential model for quarkonia.

The $2P$ and $1P$ states of $b\bar{b}$ and their fine-structure splittings have been observed in a series of experiments at the Cornell Electron Storage Ring by Han *et al.*,¹ Klopfenstein *et al.*,² and Haas *et al.*³ We would like to comment on the significance of the experimental findings of these authors, and provide new theoretical results for $E1$ transitions in $b\bar{b}$ based upon our quantum-chromodynamic potential model for quarkonia.⁴

As indicated by Haas *et al.*³ and Gilchriese,⁵ the Cornell experimental results for the center of gravity and fine-structure splittings of the P states are in excellent agreement with our theoretical predictions.⁴ It is also noteworthy that these experimental results rule out a non-quantum-chromodynamic potential model.⁶ It is generally recognized that the spectroscopy of heavy quarkonia offers a unique testing ground for quantum chromodynamics because it does not involve the use of prescriptions such as hadronization. Thus, the overall agreement between our theoretical predictions for $b\bar{b}$ and the experimental findings represents a clear *quantitative* triumph of quantum chromodynamics.

In order to further compare the theoretical and experimental results for $b\bar{b}$, we have calculated the widths for the

$E1$ transitions

$$\Gamma_{E1}(^3S_1 \rightarrow ^3P_J) = \frac{4}{9} \frac{2J+1}{3} \alpha e_Q^2 k_J^3 |r_{fl}|^2,$$

$$\Gamma_{E1}(^3P_J \rightarrow ^3S_1) = \frac{4}{9} \alpha e_Q^2 k_J^3 |r_{fl}|^2,$$

and our results are given in Tables I and II. These widths have been obtained purely theoretically without any input from the experimental results of Refs. 1-3, and they are in agreement with the experimental data for those transitions that have been observed so far.

Further experimental and theoretical investigations of the $b\bar{b}$ system are highly desirable. Indeed, it is hoped that $b\bar{b}$ will become the hydrogen atom of quantum chromodynamics.⁵

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TABLE I. Theoretical widths for the $E1$ transitions $^3S_1 \rightarrow ^3P_J$.

Transition	k_J (MeV)	$ r_{fl} $ (GeV ⁻¹)	Γ_{E1} (keV)
$^3S_1 \rightarrow ^2^3P_2$	89	2.60	2.9
$\rightarrow ^2^3P_1$	103		2.7
$\rightarrow ^2^3P_0$	123		1.5
Total			7.1
$^3S_1 \rightarrow ^1^3P_2$	445	0.01	4.6×10^{-3}
$\rightarrow ^1^3P_1$	462		3.1×10^{-3}
$\rightarrow ^1^3P_0$	487		1.2×10^{-3}
Total			8.9×10^{-3}
$^2^3S_1 \rightarrow ^1^3P_2$	103	1.64	1.8
$\rightarrow ^1^3P_1$	120		1.7
$\rightarrow ^1^3P_0$	145		1.0
Total			4.5

TABLE II. Theoretical widths for the $E1$ transitions $^3P_J \rightarrow ^3S_1$.

Transition	k_J (MeV)	$ r_{fl} $ (GeV ⁻¹)	Γ_{E1} (keV)
$^2^3P_2 \rightarrow ^3S_1$	253	1.99	23.0
$^2^3P_1 \rightarrow$	239		19.4
$^2^3P_0 \rightarrow$	219		14.9
Total			57.3
$^2^3P_2 \rightarrow ^1^3S_1$	804	0.23	9.7
$^2^3P_1 \rightarrow$	790		9.2
$^2^3P_0 \rightarrow$	770		8.5
Total			27.4
$^1^3P_2 \rightarrow ^1^3S_1$	448	1.18	45.4
$^1^3P_1 \rightarrow$	431		40.4
$^1^3P_0 \rightarrow$	406		33.8
Total			119.6

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⁵M. G. D. Gilchriese, in *Particles and Fields—1983*, proceedings of the Meeting of the APS Division of Particles and Fields, Blacksburg, Virginia, edited by A. Abashian (AIP, New York, 1984), p. 93.

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