

## Comments and Addenda

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### $\Lambda p$ Resonance Observed in Nuclear Emulsion

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It is found that the reported enhancement in the  $\Lambda p$  mass distribution at 2110 MeV in an emulsion experiment of Jain is probably due to experimental biases in the data.

RECENTLY, an enhancement in the  $\Lambda p$  mass distribution has been reported<sup>1</sup> from an analysis of  $K^-$  capture at rest with nuclear emulsions giving rise to  $\Lambda$ ,  $p$ , and  $\pi^-$  final state particles. The enhancement is situated at 2110 MeV and has a full width at half-maximum of about 20 MeV. With these observations, a claim of a possible  $\Lambda p$  resonance having  $I = \frac{1}{2}$ ,  $J = 1$ ,  $l = 0$  has been made. Earlier several bubble-chamber experiments performed with  $K^-$  interactions in deuterium have been reported<sup>2-6</sup> in which an enhancement in the  $\Lambda p$  mass distribution from  $K^-d \rightarrow \Lambda p \pi^-$  reaction has been observed. The mass  $M$  and width  $\Gamma$  of the enhancement are observed as  $M \approx 2126$  MeV,  $\Gamma \sim 10$  MeV by Cline *et al.*<sup>3</sup>;  $M \approx 2130$  MeV,  $\Gamma \sim 10$  MeV by Alexander *et al.*<sup>4</sup>;  $M \approx 2129$  MeV,  $\Gamma \sim 7$  MeV by Tan<sup>5</sup>; and  $M \approx 2129$  MeV,  $\Gamma \sim 10$  MeV by the BEGI collaboration.<sup>6</sup> The only positive claim of a resonance in these experiments is made by Cline *et al.*,<sup>3</sup> whose  $\Lambda p$  mass peak is well below the  $\Sigma^+n$  threshold of  $\approx 2129$  MeV. The results of other experiments have been mostly explained as due to a kinematic effect.

The mass and width of the  $\Lambda p$  enhancement,  $M \approx 2110$  MeV,  $\Gamma \approx 20$  MeV, as observed in emulsion<sup>1</sup> is clearly inconsistent with the values  $M = 2126-2130$  MeV and  $\Gamma \leq 10$  MeV, quoted in deuterium experiments.<sup>3-6</sup> In what follows, we wish to point out that the emulsion

result<sup>1</sup> could be subject to experimental biases, and therefore the observed result is of a doubtful nature.

It is well known that the scanning efficiency for two-prong stars (due to  $V^0$  decays) in nuclear emulsions is very low. Further, once having found a  $V^0$ -decay event, its correlation with the parent star, which may be separated by distances  $\gtrsim 1$  cm, cannot be found reliably. With the exception of mass measurements (which do not require any knowledge of the parent star), properties like spin, lifetime, decay modes, etc. for the  $\Lambda$  hyperon have for these reasons been studied almost exclusively using bubble-chamber or cloud-chamber techniques. Several years ago, we reported<sup>7</sup> our studies of the  $\Lambda$  hyperon in nuclear emulsion. We scanned  $\Lambda$  events by three different methods.

(i) Following back the tracks of  $\pi^-$  mesons from  $\pi^-$ -capture stars: In this method, an area scan was first made for  $\pi^-$ -capture stars in the region of stopping  $K^-$  mesons. The pion tracks from these stars were then traced back to their points of origin or 25 mm, whichever distance was less.

(ii) Area scan for  $V$  events: In this method a direct search was made under low magnification ( $10 \times 15$ ) for two-prong stars in which one of the secondary tracks may be due to a  $\pi$  meson and the other due to a proton.

(iii) Area scan for "hanging tracks": In this method an area scan was made under low magnification ( $10 \times 15$ ) for grey or black tracks which appear to have originated within the emulsion pellicle. The point of origin of the track was then scrutinized under higher magnification ( $100 \times 15$ ) to detect the presence of an associated track which may be otherwise missed under low magnification.

<sup>7</sup> B. Bhowmik, D. P. Goyal, and N. K. Yamdagni, Nuovo Cimento **22**, 296 (1961); **21**, 1066 (1961); **23**, 108 (1962).

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<sup>1</sup> P. L. Jain, Phys. Rev. **187**, 1816 (1969).

<sup>2</sup> O. I. Dahl, N. Horwitz, D. H. Miller, J. J. Murray, and P. G. White, Phys. Rev. Letters **6**, 142 (1961).

<sup>3</sup> D. Cline, R. Laumann, and J. Mapp, Phys. Rev. Letters **20**, 1452 (1968).

<sup>4</sup> G. Alexander, B. H. Hall, N. Jew, G. Kalmus, and A. Kernan, Phys. Rev. Letters **22**, 483 (1969).

<sup>5</sup> Tai Ho Tan, Phys. Rev. Letters **23**, 395 (1969).

<sup>6</sup> Birmingham-Edinburgh-Glasgow-Imperial College collaboration (unpublished).

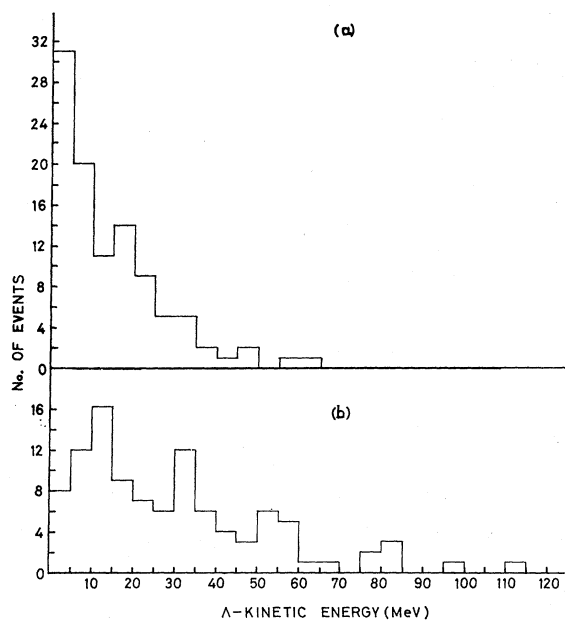


FIG. 1. Energy spectra of  $\Lambda$  events: (a) data of Ref. 1; (b) our data.

The first method gives events having a restricted range of the decay pion and thus the  $\Lambda$  sample obtained would be biased. Events found by method (ii) preferentially record those events in which both the decay  $\pi^-$  and  $p$  are rather slow, appearing as grey or black tracks, and having an opening angle fairly large,  $> 30^\circ$ . Most events missed by method (ii), however, could be recorded by method (iii), the decay proton always being expected to be "grey" for the highest energy of the  $\Lambda$ ,  $\sim 150$  MeV. Some events are, of course, still missed by both methods; in particular, those having a very short proton track and a rather fast pion track, either of which is difficult to observe under low magnification.

In Figs. 1(a) and 1(b) we reproduce the  $\Lambda$ -energy spectra for events reported in Ref. 1 and for 103 of our events found by scanning methods (ii) and (iii). A comparison of these figures reveals a general disagreement of shapes of the two distributions. The proportion

of slow  $\Lambda$ 's having energy  $\leq 20$  MeV, is 75% in Fig. 1(a), which is considerably higher than the 45% found in Fig. 1(b). The highest energy of a  $\Lambda$  in Fig. 1(a) is  $\sim 65$  MeV, while in Fig. 1(b) it is  $\sim 115$  MeV. The above disagreement is not surprising since events collected in Ref. 1 are those scanned by method (ii) only, which selects preferentially low-energy  $\Lambda$ 's.

We next examine the question of correlation of an observed  $\Lambda$  with the parent  $K^-$  star. The actual production of an observed  $\Lambda$  in emulsion could be either from a  $K^-$ -capture star, or from a  $\Sigma^\pm$  interaction or capture star, or from a  $\Sigma^0 \rightarrow \Lambda + \gamma$  decay. Further, a  $\Lambda$  after being produced could, in some cases, change its original direction because of nuclear scattering in the dense medium of the emulsion. A scan along the calculated direction of flight from the observed  $\Lambda$  decay could therefore lead to several choices of the parent stars. As reported in Ref. 1, in some cases more than one  $K^-$ -capture star corresponding to a particular  $\Lambda$  is indeed found. The possibility of the  $\Lambda$  being produced from a  $\Sigma^\pm$  interaction or  $\Sigma^0$  decay, however, seems to have been ignored. A distribution of the measured proper lifetime of the  $\Lambda$  events could have thrown some light on the reliability of the sample collected in Ref. 1. However, no such estimate is reported.

It may be further remarked that a  $K^-$  capture in emulsion would take place on one of its constituent nuclei, which are mainly C, N, O, Ag, and Br. The assumption, then, that an observed final state of  $\Lambda$ ,  $\pi^-$ , and  $p$  has arisen due to a  $K^-$  capture on two nucleons only, and the presence of other nucleons has not distorted the primary interaction, is not well substantiated in Ref. 1.

To conclude, the reported  $\Lambda p$  enhancement in nuclear emulsion<sup>1</sup> is open to considerable doubt because of experimental difficulties and biases discussed above. The  $\Lambda p$  enhancement probably occurs due to a preferential selection of low-energy  $\Lambda$ 's.

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