Comments on "Anharmonic oscillator"

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It is pointed out that the quantization conditions derived by Lu, Wald, and Young in a paper on an anharmonic oscillator are the same as the well-known JWKB quantization condition, particularized to JWKB approximations of the third and fifth orders, which can be derived in a much simpler and more general way. A number of corrections to the above-mentioned paper by Lu, Wald, and Young are also pointed out.

The Miller and Good¹ method, as extended by Lu^2 to include higher-order terms, was used in a paper by Lu, Wald, and Young³ to calculate the energy levels of an anharmonic oscillator. In the present note we point out that, as particular cases of the well-known JWKB-quantization condition, one immediately obtains the quantization conditions derived in a more complicated way by Lu, Wald, and Young.³ We also point out a number of corrections to their paper.

Since Lu, Wald, and Young³ use the harmonic oscillator, corresponding to $p_2^2 = 2m[(n + \frac{1}{2})\hbar(k/m)^{1/2} - \frac{1}{2}kx^2]$, as the solved problem, the right-hand member of their Eq. (6) as well as that of their Eq. (16) is easily seen to be equal to $2\hbar(n + \frac{1}{2})\pi$. One therefore realizes that their quantization conditions (6) and (16) to the orders \hbar^2 and \hbar^4 , respectively, are the same as the usual JWKB-quantization condition of the third order and of the fifth order, respectively. Simple, transparent ways of deriving this quantization condition in arbitrary order have recently been reviewed by Fröman and Fröman.⁴

We also point out that the following corrections should be introduced in the paper by Lu, Wald, and

Young.³ In Eq. (4) the term $-(1-\alpha)S$ in the square brackets should be replaced by $-(1-\alpha)S^2$. On the right-hand side of Eq. (5b), the numerator $(a^2\hbar^2/m)^{1/2}$ should be replaced by $(a^2\hbar^2/m)^{1/3}$. On the right-hand side of Eq. (19) inside the square brackets, the last term in the parentheses multiplying $F(\frac{1}{2}\pi, K)$ should not be $-3a^4b^2$ but $-3a^2b^4$. On the right-hand side of Eq. (21), inside the square brackets, the next-to-last term in the parentheses multiplying $F(\frac{1}{2}\pi, K)$ should not be $-345a^2b^2$ but $-345a^2b^6$. On the right-hand side of Eq. (22) the factor $224(1 - \alpha^2)$ in front of the first integral should be replaced by $224(1-\alpha)^2$, in the integrand of the fifth integral $(a^2 + S^2)^{7/2}$ should be replaced by $(a^2 + S^2)^{5/2}$, and in the parentheses in front of the sixth (i.e., the last) integral the first term should not be $-224(1-\alpha^2)/\alpha^{21/4}$ but $-224(1-\alpha)^2/\alpha^{21/4}$.

From what was said in the beginning of this note it is clear that by calculating the energy levels by means of the JWKB-quantization condition in the third and in the fifth orders, we should reproduce the figures in columns (b) and (c), respectively, of Table I in the paper by Lu, Wald, and Young.³ Such a check shows that the values in

n	$\alpha = 0.2$	$\alpha = 0.4$	$\alpha = 0.6$	$\alpha = 0.8$	$\alpha = 1.0$
0	2.03946	2.14574	2.23348	2.31566	2.397 99
1	6,51073	7.34173	8,14573	8.89626	9.59651
2	11,62919	13.69182	15.55664	17.24233	18.78613
3	17,23316	20.838 88	23,98999	26.79553	29.34300
4	23,23906	28.62427	33.23445	37,30223	40.97724
5	29.59302	36.951 92	43.16390	48.61194	53.51735
6	36.256 68	45.75586	53.69261	60.62282	66.84756
7	43.200 97	54.987 39	64.75772	73.26043	80.88340
8	50.40292	64.608 72	76.31073	86.46749	95.55998
9	57.84384	74.589 53	88.31288	100,19830	110.82557

TABLE I. Corrected values in columns (c) of Table I in the paper by Lu, Wald, and Young.³

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columns (b) are correct (except for one unit in the last figure at two places), while the values in columns (c) are wrong. The reason for this may be due to some of the errors in the formulas pointed out above or to errors in the numerical calculations, but we have not investigated this question further. We only give the correct values which should replace the values in columns (c) of Table I in the paper by Lu, Wald, and Young.³

Finally we remark that the difference between the results to second order in \hbar , given by Lu, Wald, and Young³ in columns (b) of their Table I, and the JWKB results obtained by Handelsman and Lew⁵ to second order in \hbar , quoted by Lu, Wald, and Young³ in columns (d) of their Table I, is only due to the fact that Handelsman and Lew,⁵ starting from the JWKB-quantization condition, obtain the eigenvalue in the form of an explicit series expansion, which they truncate after a finite number of terms.

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