

Corrections to the sixth-order anomalous magnetic moment of the muon

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The contribution to the sixth-order muon anomaly from fourth-order electron vacuum polarization is determined analytically to order  $m_e/m_\mu$ . The result including the contribution from graphs containing two second-order lepton vacuum polarization subgraphs is  $(\alpha/\pi)^3\{(2/9)\ln^2(m_\mu/m_e) + (403/108 - 4\pi^2/9) \times \ln(m_\mu/m_e) + \zeta(3)/2 + 2\pi^2/27 + 5/27 + [383\pi^2/135 - 16\pi^2\ln(2)/9 - 13\pi^3/18]m_e/m_\mu\}$ .

In sixth order, the difference between the muon and electron magnetic moments can be expressed, for  $m_\mu/m_e \gg 1$ , as

$$a_\mu^{(6)} - a_e^{(6)} = \left(\frac{\alpha}{\pi}\right)^3 [A \ln^2(m_\mu/m_e) + B \ln(m_\mu/m_e) + C + D(m_e/m_\mu) + O((m_e/m_\mu)^2 \ln^2(m_\mu/m_e))]. \tag{1}$$

A and B are completely known analytically<sup>1,2</sup>:

$$A = \frac{2}{9},$$

$$B = \frac{31}{27} + 7\pi^2/9 - \frac{2\pi^2}{3} \ln 2 + \zeta(3). \tag{2}$$

All contributions to C, except the light-by-light contribution  $C^{(\gamma\gamma)}$ , are also known analytically<sup>3</sup> ( $C^{(\gamma\gamma)}$  is known numerically<sup>4,5</sup>):

$$C = \frac{1075}{216} - \frac{25}{18}\pi^2 + \frac{5}{3}\pi^2 \ln 2 - 3\zeta(3) + 3C_4 + C^{(\gamma\gamma)},$$

where

$$C_4 = \frac{11}{648}\pi^4 - \frac{2}{27}\pi^2 \ln^2 2 - \frac{1}{27} \ln^4 2 - \frac{8}{9}a_4 \tag{3}$$

and

$$a_4 = \sum_1^\infty \frac{1}{2^n n^4}.$$

The only contribution to D which is known analytically is the double-bubble contribution<sup>6</sup> [diagram (d) of Fig. 1]

$$D^{(d)} = -4\pi^2/45. \tag{4}$$

The contribution to D due to the other diagrams of Fig. 1 is known numerically:

$$D^{(a+b+c)} = -5.6776256, \tag{5}$$

$$D^{(e+f)} = 0.$$

In this paper we present an analytic calculation of  $D_{(a+b+c)}$ , the contribution to D from fourth-order electron vacuum polarization [the proper diagrams (a), (b), and (c) of Fig. 1].

This quantity is given by the following expression<sup>7,8</sup>:

$$D^{(a+b+c)} = \frac{\pi}{2} - 2\pi \int_0^1 \frac{dx x}{(1-x^2)^{3/2}} \left[ \frac{1}{\pi} \text{Im}\pi^{*(4)}(x) - \frac{1}{\pi} \text{Im}\pi^{*(4)}(1) \right] / \left(\frac{\alpha}{\pi}\right)^2, \tag{6}$$

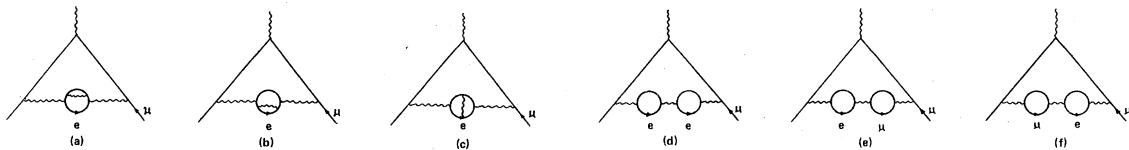


FIG. 1. Feynman diagrams representing the fourth-order vacuum polarization contribution to the sixth-order muon anomaly.

where

$$\begin{aligned} \frac{1}{\pi} \text{Im}\pi^{*(4)}(x) = & \left(\frac{\alpha}{\pi}\right)^2 \left\{ \left(-\frac{5}{8}x - \frac{3}{8}x^3\right) + x\left(-\frac{1}{2} + \frac{1}{6}x^2\right) \ln\left(\frac{64x^4}{(1-x^2)^3}\right) + \left(\frac{11}{16} + \frac{11}{24}x^2 - \frac{7}{48}x^4\right) \ln\left(\frac{1+x}{1-x}\right) \right. \\ & + \left(\frac{1}{2} + \frac{1}{3}x^2 - \frac{1}{6}x^4\right) \ln\left(\frac{(1+x)^3}{8x^2}\right) \ln\left(\frac{1+x}{1-x}\right) \\ & \left. - \left(\frac{1}{2} + \frac{1}{3}x^2 - \frac{1}{6}x^4\right) \left[4\Phi\left(-\frac{1-x}{1+x}\right) + 2\Phi\left(\frac{1-x}{1+x}\right) + \frac{\pi^2}{2}\right] \right\} \end{aligned} \quad (7)$$

and

$$\frac{1}{\pi} \text{Im}\pi^{*(4)}(1) = \frac{1}{4}(\alpha/\pi)^2. \quad (8)$$

So we can write

$$D^{(a+b+c)} = \frac{\pi}{2} - 2\pi \left[ R_1 + R_2 + R_3 + R_4 + R_5 - \frac{1}{4} \int_0^1 \frac{dx x}{(1-x^2)^{3/2}} \right], \quad (9)$$

where the five  $R_i$  correspond to the five terms in Eq. (7). It is easy to see that  $R_2 + R_3$ ,  $R_4$ , and  $R_5$  are finite and the combination

$$R'_1 = R_1 - \frac{1}{4} \int_0^1 \frac{x}{(1-x^2)^{3/2}} \quad (10)$$

is also finite.

We now evaluate the integrals. Our results are as follows:

$$R'_1 = \frac{1}{4} - \frac{\pi}{32}, \quad (11)$$

$$R_2 + R_3 = \pi \ln 2 - \frac{211}{288}\pi, \quad (12)$$

and

$$R_4 + R_5 = \frac{13}{36}\pi^2 - \frac{\pi}{9} \ln 2 - \frac{15}{216}\pi. \quad (13)$$

Adding the terms in Eqs. (11), (12), and (13) and substituting into Eq. (9), we obtain our result

$$\begin{aligned} D^{(a+b+c)} &= -\frac{13}{18}\pi^3 - \frac{16}{9}\pi^2 \ln 2 + \frac{79}{27}\pi^2 \\ &= -5.6776257. \end{aligned} \quad (14)$$

This is in excellent agreement with the numerical value in Eq. (5). It is interesting that, although the term proportional to  $\pi$  cancels out, there remains a  $\pi^3$  term. This is the first time an odd power of  $\pi$  occurs in a  $g-2$  contribution.

Using Eqs. (4), (5), and (14), the contribution to  $D$  from all the graphs of Fig. 1 can be written:

$$D^{(a+b+c)} + D^{(d)} + D^{(e+f)} = -\frac{13}{18}\pi^3 - \frac{16}{9}\pi^2 \ln 2 + \frac{383}{135}\pi^2. \quad (15)$$

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<sup>8</sup>B. E. Lautrup and E. de Rafael, Phys. Rev. 174, 1835 (1968).