

## Comments and Addenda

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### Diagonalization of the Weyl tensor

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The gravitational superenergy flux can usually be reduced to zero by a suitable Lorentz transformation. The latter is derived here explicitly.

The following problem was recently posed by Wheeler<sup>1</sup>: Find the Lorentz transformation diagonalizing  $\mathcal{S}$  and  $\mathcal{B}$  (the two  $3 \times 3$  symmetric traceless matrices representing the Weyl tensor). This problem can easily be solved by noting that the complex matrix  $\mathcal{S} + i\mathcal{B}$  transforms under the Lorentz group according to<sup>2,3</sup>

$$\mathcal{S}' + i\mathcal{B}' = M^{-1}(\mathcal{S} + i\mathcal{B})M,$$

where  $M$  is a complex orthogonal  $3 \times 3$  matrix. The latter is related in a well-known way<sup>4</sup> to the Lorentz-transformation matrix, e.g.,

$$M_1^1 = L_2^2, L_3^3, -L_2^3, L_3^2, +iL_2^0, L_3^1, -iL_2^1, L_3^0.$$

To diagonalize  $\mathcal{S} + i\mathcal{B}$ , we need only the following steps:

- (1) Solve the characteristic equation

$$|\mathcal{S} + i\mathcal{B} - \lambda I| = 0.$$

This is a cubic equation for  $\lambda$ ; see Ref. 2, Eqs. (9–11).

- (2) In the generic case, the three eigenvalues are different and the linear equations

$$(\mathcal{S} + i\mathcal{B})M = \lambda M$$

have a unique solution, up to a sign. [We want  $M$  to be orthogonal, so we must normalize it by  $M^T M = I$  and  $\text{Det}(M) = 1$ .]

- (3) Express  $M$  in terms of complex Euler angles.<sup>5</sup>

- (4) Reconstruct  $L$  from their real and imaginary parts.<sup>5</sup>

<sup>1</sup>J. A. Wheeler, *Phys. Rev. D* **16**, 3384 (1977).

<sup>2</sup>A. Peres, *Nuovo Cimento* **18**, 36 (1960).

<sup>3</sup>A. Peres, *J. Math. Mech.* **11**, 61 (1962).

<sup>4</sup>A. Einstein and W. Mayer, *Sitzungsber. Preuss. Akad.*

Wiss. 522 (1932).

<sup>5</sup>F. Halbwachs, P. Hillion, and J. P. Vigier, *Ann. Inst. H. Poincaré* **16**, 115 (1959).