

Comments and Addenda

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Pontryagin density in Yang's R gauge

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A simple and computationally useful expression for the Pontryagin density in Yang's R gauge for self-dual SU(2) gauge fields is exhibited.

Yang¹ has shown how, with a suitable choice of gauge (the R gauge), the self-dual SU(2) gauge field equations reduce to Laplace-type equations for one real variable (φ) and one complex variable (ρ). The familiar Corrigan-Fairlie-Wilczek-Hooft ansatz turns out to be an especially simple solution to these equations. The purpose of this note is to exhibit a simple and computationally useful expression for the Pontryagin density in terms of the variables φ and ρ . As in Ref. 1

all considerations are local in character, and do not refer to global properties.

For self-dual gauge fields, the action density $S = \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a$ and the Pontryagin density $*S = \frac{1}{4} *F_{\mu\nu}^a F_{\mu\nu}^a = \frac{1}{8} \epsilon_{\mu\nu\lambda\rho} F_{\mu\nu}^a F_{\lambda\rho}^a$ are equal to each other, with (using the notation of Ref. 1)

$$S = F_{yz}^a F_{yz}^a + \frac{1}{2} (F_{yz}^a F_{yz}^a + F_{zx}^a F_{zx}^a). \tag{1}$$

A direct substitution of the R-gauge potentials into (1) gives

$$S = 2\varphi^2 \left[\left(\frac{\rho_y}{\varphi^2} \right)_{\bar{y}} \left(\frac{\bar{\rho}_y}{\varphi^2} \right)_y + \left(\frac{\rho_z}{\varphi^2} \right)_{\bar{z}} \left(\frac{\bar{\rho}_z}{\varphi^2} \right)_z + \left(\frac{\rho_x}{\varphi^2} \right)_{\bar{x}} \left(\frac{\bar{\rho}_x}{\varphi^2} \right)_x + \left(\frac{\rho_y}{\varphi^2} \right)_{\bar{z}} \left(\frac{\bar{\rho}_y}{\varphi^2} \right)_z \right] + 2 \left\{ \left[\left(\frac{\varphi_y}{\varphi} \right)_{\bar{y}} + \frac{\rho_y \bar{\rho}_y}{\varphi^2} \right]^2 + \left[\left(\frac{\varphi_z}{\varphi} \right)_{\bar{z}} + \frac{\rho_z \bar{\rho}_z}{\varphi^2} \right]^2 + 2 \left[\left(\frac{\varphi_y}{\varphi} \right)_{\bar{z}} + \frac{\rho_y \bar{\rho}_z}{\varphi^2} \right] \left[\left(\frac{\varphi_y}{\varphi} \right)_z + \frac{\rho_z \bar{\rho}_y}{\varphi^2} \right] \right\}. \tag{2}$$

Equation (2), as it stands, does not make the topology of the gauge fields transparent. We attempt therefore to simplify (2) by repeated use of the self-duality equations [i.e., Eq. (27) of Ref. 1]. A slightly tedious amount of algebra leads one to the following result:

$$S = *S = -\frac{1}{2} \square \square \ln \varphi + 2 \left[\partial_y \partial_{\bar{y}} \left(\frac{\varphi_z \varphi_{\bar{z}} - \rho_y \bar{\rho}_y}{\varphi^2} \right) - \partial_y \partial_z \left(\frac{\varphi_z \varphi_{\bar{y}} + \rho_z \bar{\rho}_y}{\varphi^2} \right) + \partial_x \partial_{\bar{x}} \left(\frac{\varphi_y \varphi_{\bar{y}} - \rho_z \bar{\rho}_z}{\varphi^2} \right) - \partial_x \partial_{\bar{y}} \left(\frac{\varphi_y \varphi_{\bar{z}} + \rho_y \bar{\rho}_z}{\varphi^2} \right) \right]. \tag{3}$$

We were unable to reduce the expression (3) any further except to note that for all published² self-dual gauge fields the quantity in the square brackets vanishes identically so that $S = *S = -\frac{1}{2} \square \square \ln \varphi$. It is tempting to conjecture, though we have no proof, that $S = *S = -\frac{1}{2} \square \square \ln \varphi$ for all self-dual gauge fields. The expression $*S = -\frac{1}{2} \square \square \ln \varphi$ can be trivially integrated over four-dimensional

Euclidean space to yield a topological characteristic of the gauge field, i.e., the Pontryagin index.

Note added. Recently E. Corrigan *et al.* [Phys. Lett. **72B**, 354 (1978)] have constructed the sequence of Atiyah-Ward (AW) *Ansätze* A_l ($l = 1, 2, 3, \dots$) for self-dual SU(2) gauge fields. The Pontryagin density $*S_l$ for A_l satisfies the following

striking "superposition" principle:

$$*S_I = \sum_{k=1}^l \left(-\frac{1}{2} \square \square \ln {}^{(k)}\varphi \right),$$

where ${}^{(k)}\varphi$ is Yang's R -gauge φ function for the AW Ansatz A_k . This result follows from the very simple transformation properties of Eq. (3) under the Backlund transformation used in generating

A_l .

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¹C. N. Yang, Phys. Rev. Lett. **38**, 1377 (1977).

²R. Jackiw, C. Nohl, and C. Rebbi, Phys. Rev. D **15**,

1642 (1977).