

Errata

Erratum: Simple bound for K_{l3} decay parameters using the πK ($s=0, T=\frac{1}{2}$) phase shift [Phys. Rev. D 6, 1943 (1972)]

M. Micu and E. E. Radescu

The first relation of Eq. (13) should read

$$A = \frac{1 - \omega^2}{(1 + \gamma)(\sigma^2 - \omega^2)} \left[\left(\frac{1 - \omega^2}{\omega(\sigma^2 - \omega^2)} - \frac{1 + \omega^2}{2\omega^3} \right) \ln \frac{y + \omega}{y - \omega} - \frac{1 - \sigma^2}{\sigma(\sigma^2 - \omega^2)} \ln \frac{y + \sigma}{y - \sigma} + \frac{y(1 - \omega^2)}{\omega^2(y^2 - \omega^2)} \right].$$

Also, the right-hand side of the second inequality of (15) should read 0.117 instead of 0.17.

Erratum: Møller scattering and weak neutral currents [Phys. Rev. D 10, 3629 (1974)]

R. Gastmans and Y. Van Ham

L. DeRaad, Jr. pointed out that some of the polarization-dependent terms in Eq. (12) are wrong. The correct α^3 cross section for Møller scattering in the extreme relativistic limit is obtained by replacing in the coefficient of u the factor $(1 - z)^3/(1 + z)$ by $(1 - z)^3$, in the coefficient of w the factor $(5 + 6z)$ by $(4 + 7z)$, and in the coefficient of w^2 the factor $(3 - 2z)$ by $(3 - 2z - z^2)$. Complete agreement with L. DeRaad, Jr. is then obtained for this cross section [see L. DeRaad, Jr., Phys. Rev. D (to be published)].

As a result, the values for δ_α in Table III are wrong: They are now given within 0.1% by the values of δ_α in Table II. Consequently, conclusion (i) should be modified to read "and the effect of initial polarization is negligible."

Erratum: Single soft-pion emission in proton-antiproton annihilation [Phys. Rev. D 10, 3653 (1974)]

Gerald W. Intemann and Gary K. Greenhut

We have become aware of $\bar{p}p$ annihilation data at vanishing lab momentum [N. Barash *et al.*, Phys. Rev. 139, B1659 (1965)] that can be compared with our theoretical results. These data give an experimental branching ratio

$$\frac{\sigma(\bar{p}p \rightarrow K_{1,2}^0 K_{1,2}^0 \pi^0)}{2\sigma(\bar{p}p \rightarrow K_1^0 K_2^0)} = 1.0 \pm 0.3$$

which agrees fairly well with the theoretical ratio for $\sigma(\bar{p}p \rightarrow K^0 \bar{K}^0 \pi^0)/\sigma(\bar{p}p \rightarrow K^0 \bar{K}^0)$ in Eq. (2.18). After subtracting the K^* resonance contribution, the experimental data yield the branching ratio

$$\frac{\sigma(\bar{p}p \rightarrow K^0 K^+ \pi^-)}{2\sigma(\bar{p}p \rightarrow K^+ K^-)} = 1.3 \pm 0.5,$$

which agrees within two standard deviations with our result for $\sigma(\bar{p}p \rightarrow K^0 K^- \pi^+)/\sigma(\bar{p}p \rightarrow K^+ K^-)$ in Eq. (3.11).

We wish to thank Dr. Noel Yeh for bringing these additional $\bar{p}p$ data to our attention.