Comments and Addenda

The Comments and Addenda section is for short communications which are not of such urgency as to justify publication in Physical Review Letters and are not appropriate for regular Articles. It includes only the following types of communications: (1) comments on papers previously published in The Physical Review or Physical Review Letters; (2) addenda to papers previously published in The Physical Review or Physical Review Letters, in which the additional information can be presented without the need for writing a complete article. Manuscripts intended for this section should be accompanied by a brief abstract for information-retrieval purposes. Accepted manuscripts will follow the same publication schedule as articles in this journal, and galleys will be sent to authors.

Type-N gravitational field with twist – Analytical addendum to Hauser's solution*

Daniel E. Novoseller

California Institute of Technology, Pasadena, California 91125 (Received 30 December 1974)

Hauser has presented a type-N gravitational field with twist in terms of a function f(y), which he expressed as the solution of a differential equation. The exact solutions of this equation are presented here as Gauss hypergeometric functions.

Recently,¹ Hauser presented an exact type-N solution of the Einstein field equations with twisting rays. He presents the solution in terms of a function f(y) which satisfies

$$(1+y^2)\frac{d^2f}{dy^2} + \frac{3}{16}f = 0 \quad . \tag{1}$$

Hauser notes that the two solutions of Eq. (1) have the following properties:

(a) The even solution has zeros only at $y=\pm 5.5$. The odd solution has a zero only at y=0. We will denote these solutions as $f_e(y)$ and $f_o(y)$, with

$$f_e(y) = f_e(-y) , \qquad (2)$$

$$f_o(y) = -f_o(-y) \ .$$

.

(b) As $y \rightarrow \infty$, two independent solutions behave like

$$f_{3/4} - C_1 y^{3/4} ,$$

$$f_{1/4} - C_2 y^{1/4} .$$
(3)

We present here the exact solutions of Eq. (1)

in terms of the hypergeometric function.^{2,3} If we define

$$z = -y^2 ,$$

$$\dot{f} = \frac{df}{dz}$$
(4)

then Eq. (1) becomes

$$z(1-z)\ddot{f} + (\frac{1}{2} - \frac{1}{2}z)\dot{f} - \frac{3}{64}f = 0 .$$
 (5)

The two solutions in the neighborhood of zero are given by $Erdélyi^2$ [Eq. 2.3.1-(1)]:

$$u_{1(0)}(z) = {}_{2}F_{1}(-\frac{1}{8}, -\frac{3}{8}, \frac{1}{2}, z) ,$$

$$iu_{2(0)}(z) = (-z)^{1/2} {}_{2}F_{1}(\frac{3}{8}, \frac{1}{8}, \frac{3}{2}, z) .$$
(6)

These solutions are proportional to the even and odd solutions of Eq. (2), with

$$f_{e}(y) = a_{1}u_{1(0)}(-y^{2}),$$

$$f_{o}(y) = a_{2}u_{2(0)}(-y^{2}),$$
(7)

where a_1 , a_2 are constants. Note that Eq. 15.4.25 of Abramowitz and Stegun³ implies

$$u_{1(0)}(-y^2) = \frac{\Gamma\left(\frac{3}{8}\right)\Gamma\left(\frac{11}{8}\right)}{\sqrt{\pi} \, 2^{3/4}} \left(1 + y^2\right)^{1/4} \left\{ P_{-3/2}^{-1/4} \left(\left[y^2/(1+y^2) \right]^{1/2} \right) + P_{-3/2}^{-1/4} \left(- \left[y^2/(1+y^2) \right]^{1/2} \right) \right\}$$
(8)

for $0 < y < \infty$.

Expressions for $u_{1(0)}(-y^2)$ and $u_{2(0)}(-y^2)$ for y = 1 are given by Eq. (6), since these hypergeometric series are absolutely convergent at $|-y^2| = 1$.

As $z \to \infty$, these solutions are better understood by analytic continuation [Erdélyi², Eq. 2.10-(2)], which gives

$$u_{1(0)}(z) = B_{1}(-z)^{1/8} {}_{2}F_{1}(-\frac{1}{8}, \frac{3}{8}, \frac{5}{4}, 1/z) + B_{2}(-z)^{3/8} {}_{2}F_{1}(-\frac{3}{8}, \frac{1}{8}, \frac{3}{4}, 1/z) ,$$

$$(9)$$

$$iu_{2(0)}(z) = B_{3}(-z)^{1/8} {}_{2}F_{1}(\frac{3}{8}, -\frac{1}{8}, \frac{5}{4}, 1/z)$$

 $+B_4(-z)^{3/8} {}_2F_1(\frac{1}{8},-\frac{3}{8},\frac{3}{4},1/z),$

2331

with

$$B_{1} = \frac{\Gamma(\frac{1}{2})\Gamma(-\frac{1}{4})}{\Gamma(-\frac{3}{8})\Gamma(\frac{5}{8})}, \quad B_{2} = \frac{\Gamma(\frac{1}{2})\Gamma(\frac{1}{4})}{\Gamma(-\frac{1}{8})\Gamma(\frac{7}{8})},$$

$$B_{3} = \frac{\Gamma(\frac{3}{2})\Gamma(-\frac{1}{4})}{\Gamma(\frac{1}{8})\Gamma(\frac{6}{8})}, \quad B_{4} = \frac{\Gamma(\frac{3}{2})\Gamma(\frac{1}{4})}{\Gamma(\frac{3}{8})\Gamma(\frac{1}{8})}.$$
(10)

Thus the solutions of interest as $y \rightarrow \infty$ are

$$V_{1(\infty)}(-y^2) = y^{1/4} {}_2F_1(-\frac{1}{8}, \frac{3}{8}, \frac{5}{4}, -1/y^2) ,$$

$$V_{2(\infty)}(-y^2) = y^{3/4} {}_2F_1(-\frac{3}{8}, \frac{1}{8}, \frac{3}{4}, -1/y^2) ,$$
(11)

which give the behavior of Eq. (3). Note that Eq.

15.4.11 of Abramowitz and Stegun³ tells us that

$$V_{1(\infty)}(-y^2) = 2^{1/4} \Gamma\left(\frac{5}{4}\right) (1+y^2)^{1/4} P_{-3/2}^{-1/4} \left[y^2/(1+y^2) \right]^{1/2} ,$$
(12)

$$V_{2(\infty)}(-y^2) = 2^{-1/4} \Gamma(\frac{3}{4})(1+y^2)^{1/4} P_{-3/2}^{1/4} \left[\left[\frac{y^2}{(1+y^2)} \right]^{1/2} \right]$$

for $0 < y < \infty$.

A preliminary computer search⁴ for zeros of the even and odd solutions found no zeros in $f_0(y)$ other than at y = 0, and found one zero in $f_e(y)$, at $y \approx \pm 5.467$. (This was also found in Ref. 1.)

York, 1953), Vol. 1.

- ³Handbook of Mathematical Functions, edited by M. Abramowitz and I. Stegun, National Bureau of Standards Applied Mathematics Series, No. 55 (U. S. G. P. O., Washington, D. C., 1964).
- ⁴The computer program for the hypergeometric function was written by Dr. M. Griss.
- *Work supported in part by the U. S. Atomic Energy Commission. Prepared under Contract AT(11-1)-68 for the San Francisco Operations Office, U. S. Atomic Energy Commission.
- ¹I. Hauser, Phys. Rev. Lett. <u>33</u>, 1112 (1974); <u>33</u>, 1525(E) (1974). We use the notation of this reference.
- ²Higher Transcendental Functions (Bateman Manuscript Project), edited by A. Erdélyi (McGraw-Hill, New