Deep learning network to distinguish binary black hole signals from short-duration noise transients

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Blip glitches, a type of short-duration noise transient in the LIGO-Virgo data, are a nuisance for the binary black hole (BBH) searches. They affect the BBH search sensitivity significantly because their timedomain morphologies are very similar, and that creates difficulty in vetoing them. In this work, we construct a deep-learning neural network to efficiently distinguish BBH signals from blip glitches. We introduce sine-Gaussian projection (SGP) maps, which are projections of gravitational wave (GW) frequency-domain data snippets on a basis of sine-Gaussians defined by the quality factor and central frequency. We feed the SGP maps to our deep-learning neural network, which classifies the BBH signals and blips. Whereas only simulated BBH signals are used for training, both simulated and real BBH signals are used for testing. For glitches only blips from real LIGO data are used for both testing and training. We show that our network significantly improves the identification of the BBH signals in comparison to the results obtained using traditional- χ^2 and sine-Gaussian χ^2 . For example, our network improves the sensitivity by 75% at a false-positive probability of 10^{-2} for BBHs with total mass in the range $[80, 140]M_{\odot}$ and SNR in the range [3, 8]. When tested on real GW events, it correctly identifies 95% of the events in GWTC-3. The computation time for classification is a few minutes for thousands of SGP maps on a single core. With further optimization in the next version of our algorithm, we expect a further reduction in the computational cost. Our proposed method can potentially improve the veto process in the LIGO-Virgo GW data analysis and conceivably support identifying GW signals in lowlatency pipelines.

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I. INTRODUCTION

The age of gravitational wave (GW) astronomy has accelerated since the detection of the first GW event, GW150914 [1,2]. As of now, the LIGO [3] and Virgo [4] collaborations have detected over 90 compact binary coalescence (CBC) signals [5–10], which include binary black boles (BBHs), neutron-star-black-hole binaries and binary neutron stars. Now, with KAGRA [11] joining the global network and with LIGO-India expected to come online later this decade [12], the rate of detection of CBC events is expected to increase quite significantly. One challenge that detection efforts face is erosion of some sensitivity of our detection pipelines caused by noise transients that share some characteristics with CBC waveforms [13]. Some of these noise artifacts are difficult to differentiate from CBC signals, especially when the latter are from binaries with large total mass [14,15]. Considerable efforts have been invested in guarding against misclassification of those transients as CBC signals. This work contributes to that effort.

One specific type of short-duration noise transient, known as a blip glitch, is one of the major contributors to degrading the search sensitivity of CBC signals as well as other short-duration GW signals [2,15]. Recent studies on blip glitches found that both LIGO detectors (located in Livingston and Hanford) report around 2 blip glitches per hour [16]. Blips are found in Virgo and GEO600 detectors as well [16]. The origin of a majority of blip glitches is still a mystery [16]. When it comes to the morphology of blip glitches, they are of short duration, $\mathcal{O}(10)$ ms in the time domain, and have a bandwidth of more than 100 Hz in the frequency domain. The shape of a blip glitch in the time domain resembles GW signals from BBHs either with a large total mass, with highly asymmetric component masses, or with component spins and orbital angular momentum antialigned. As a result, blips affect the efficiency of search of such BBH signals. Studies conducted to diagnose the source of blip glitches show that they have very little correlation with auxiliary channels. Only a few blips show correlations with the laser intensity

stabilization, computer errors, power recycling control signals, and relative humidity inside the detector. In total, only around 8% of LIGO Hanford (H1) and 2% of LIGO Livingston (L1) blips in the second observation run (O2) run have shown correlations with any of the auxiliary channels [16]. It is possible that more than one physical mechanism is responsible for blips. Hence, it is important to understand blip glitches and mitigate their adverse effects on the search sensitivity to CBCs and other short-duration signals. In this work, we have mainly used blip glitches identified by the Gravity Spy [17–20] project from H1 and L1 LIGO O2 run and first part of third observational (O3a) run. Gravity Spy is a state-of-the-art tool to classify the glitches present in LIGO data. Up to now, it has identified 23 different types of glitches including blips in the LIGO data by combining efforts of human volunteers (citizen science) and deep learning networks.

In the past, there have been efforts to understand blips (e.g., [21–23]) and to mitigate their effect on CBC searches with a combination of modeling and statistical methods (e.g., [14,24,25]). Even though some of these methods are effective at vetoing blips, there is still a lot of room for improvement. In recent years, machine learning (ML) algorithms have been proposed as being useful for GW data analysis. There have been several studies on classification, characterization or parameter estimation of GW data with the help of ML algorithms such as the convolutional neural network (CNN), the recurrent neural network, and the variational autoencoders (e.g., [26-37]). CNNs are better suited for spatial data such as images (e.g., [17,27]), whereas recurrent neural networks work well with temporal or sequential data such as text or video (e.g., [32,33]). Likewise, different ML algorithms can help in achieving various goals of GW data analysis. The purpose of this work is to exploit the tremendous image classification capabilities of CNNs in order to veto the blip glitches that appear in the LIGO-Virgo data. Better discrimination of glitches against BBH signals help in improving the significance of BBH detections. This helps in the identification of weak signals, which reduces the Malmquist bias and, thereby, aids the characterization of BH and BBH populations more accurately.

This paper is arranged as follows. In Sec. II, we introduce the sine-Gaussian projection maps and explain how to generate them. In Sec. III, we give details of the simulations and data preparation. Section IV outlines the neural network that we constructed in this work. Section V describes the existing methods like traditional χ^2 and sine-Gaussian χ^2 to veto blip like glitches. Finally, in Sec. VI, we present the results and performance of the network followed by Sec. VII on conclusions and future prospects.

II. SINE-GAUSSIAN PROJECTION MAP

A critical aspect of ML is presenting the data to the network in the most appropriate way. The data representation used in ML networks can have a significant impact on their performance. There are machine learning studies where the two-dimensional spectrogram of GW data like omega Q-scan or omicron time-frequency (TF) maps [38] have been used to categorize glitches (e.g., [17]), and continuous wavelet transform-based TF maps are used to classify signals from glitches (e.g., [27]). In this section, we introduce the sine-Gaussian projection (SGP) maps and explain how they are a useful way to represent GW data when it comes to distinguishing blips from BBH signals using a deep-learning image classifier.

Blips are one of the classes of glitches that resemble high mass-ratio or high component-mass BBH signals in the time domain and show similar TF morphology [14,16], as illustrated in Fig. 1. As a result, machine-learning algorithms that make use of TF maps may not be as efficient in distinguishing blips and high-mass BBH signals. Alternatively, showing a contrast between the blips and high-mass BBH in the two-dimensional data by projecting these signals on sine-Gaussian parameter space may provide certain advantages to a machine-learning network for classifying them better. As we show below, the blips and high-mass BBH signals have strong projections in different regions of the sine-Gaussian parameter space, which comprises the quality factor (Q) and the central frequency (f_0) , and present easily distinguishable features. In this way, the SGP maps offer complementary information for exploitation in the classification problem. The projected GW data snippet on the normalized sine-Gaussian waveforms in the $Q - f_0$ parameter space is called an SGP map. These SGP maps constitute our input data to the CNN. We call our neural network a sine Gaussian projection map-convolution neural network (SIGMA-NET).

Mathematically, a sine-Gaussian waveform can be expressed as

$$g(t) \equiv A e^{-4\pi f_0^2 \frac{(t-t_0)^2}{Q^2}} \cos(2\pi f_0 t + \phi_0), \qquad (1)$$

where A is constant amplitude, Q is the quality factor, f_0 is the central frequency, ϕ_0 is the phase, and t_0 is the central time of the sine-Gaussian waveform. The projection of a data train **x** on the sine-Gaussian **g** is defined as [39]

$$(\mathbf{x}, \mathbf{g}) = 4\Re \int_{f_{\text{lower}}}^{f_{\text{upper}}} \frac{\tilde{x}^*(f)\tilde{g}(f)}{S_n(f)} df, \qquad (2)$$

where $S_n(f)$ is the power spectral density of the noise, and f_{lower} and f_{upper} are determined, respectively, by the seismic cutoff frequency and the Nyquist frequency.

To construct an SGP map, we first choose the appropriate region of the sine-Gaussian parameter-space onto which the data snippet is to be projected. For our study, the parameter ranges are $f_0 \in [20, 520]$ Hz and $Q \in [2, 80]$. The choice of these ranges is made after extensive study of the

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FIG. 1. Examples of time-frequency maps of some of the blips (right panels) and BBH events (left panels) showing the similarities and differences of the two in this two-dimensional representation.

projection of various blips and CBC signals with the objective of capturing most of their power. We then sample points in that chosen region uniformly and the data is projected over the sine-Gaussian waveform corresponding to each sampled point using Eq. (2). The projections so computed are presented using a color map, where the color represents how strong or weak the projection is, as shown in Fig. 2. The pixels are smoothed using spline interpolation.

As we can see, there are clear distinguishing features between blips and BBH signals in the SGP maps. Blip glitches show high projection at low Q (in the range from 2 to 10) and have a projection over a wider range of frequencies than the BBH signals, as shown in Fig. 2. The BBH signals, on the other hand, do not show much projection above 200 Hz. The extended projection along Qdepends on the component masses of the binary, although it is always more extended than blips along the Q values. The projections of BBH signals contrast well with the projections of blips on the SGP maps, especially, for high signalto-noise ratio (SNR) transients (e.g., above match-filter SNR of 7). For low SNRs (<7), this contrast diminishes and identifying the BBH projections correctly becomes challenging. To overcome this issue impacting low SNR signals, we use multiview learning, which is explained in Sec. III.

III. DATA SIMULATION AND GENERATION

In this study, we work with real blips and simulated nonspinning BBH signals injected in real data. Limiting our attention to nonspinning BBH signals is of some relevance in realistic searches since significant spin has been found only in a small subset of the BBHs detected so far [8]. This aspect notwithstanding, it is highly desirable to unearth more spinning BBHs since those systems can provide some fascinating exhibition of two-body dynamics predicted by general relativity. They will also be important in informing more fully our survey of the BBH and black hole populations. We plan to extend the present work to the case of spinning black holes in the future. Here we note that one of the important differences we anticipate in that study is the bandwidth of the signals, arising from the fact that the frequency of the innermost stable circular orbit (ISCO) is affected by black hole spin [40]. Signals from BBHs with higher ISCO frequencies will adversely effect their



FIG. 2. The sine-Gaussian projection map of an H1 blip from O2 run with SNR 16 (left) and a simulated BBH event injected in real noise from H1 O2 run with SNR 16 (right). Here the SNR for blips is calculated in the same way as for BBH signals. The color bar represents value of projections as calculated using Eq. (2).

distinguishability from blips owing to an increased overlap in the frequency bands. Nevertheless, the fact that the projection of the BBH signals in SGP maps exhibit a larger spread in the quality factor compared to blips is expected to provide some discriminatory power. We will explore in the future how useful this feature proves to be in searches for spinning BBHs. We next describe the procedure followed here for fetching real data and generating simulated data in the form of time series, which are then used to prepare the SGP maps.

We begin by identifying GPS times of the blip glitches from Gravity Spy and then fetch the corresponding H1 and L1 data from O2 and O3a [41]. While most of our blips are taken from O2, we choose 30% of the 6000 test sample blips from O3a. We use the "cleaned CO2" and "cleaned CO1" frame types for the O2 and O3 data, respectively [23]. For each blip, we choose a 16-sec-long data segment out of which a 2-sec-long segment is chopped such that the GPS time falls at the center of the segment while the whole data segment is used to calculate the power spectral density (PSD), as described later in this section.

We simulate nonspinning BBH signals using the IMRPhenomPv2 [42] template such that their masses $m1, m2 \in [10, 100] M_{\odot}$ and their sky locations follow a uniform distribution. The frequency cutoff on the lower end (f_{lower}) is set to 20 Hz and the upper cutoff (f_{upper}) is set to 2048 Hz. The match-filtering SNR for BBHs is also drawn uniformly from the range [4,10] for the H1 detector. This fixes the SNR for the BBH signal in the L1 detector, accordingly.

We inject the simulated BBH signals into the real O2 noise as explained below. First, we identify a 64-sec-long data segment from the O2 data that has no known triggers.

We further verify by plotting a Q-transform map and ensure that this data segment has no GW signals and is artifact free. Next, we randomly draw 2-sec-long data segments from this main segment such that the peak of the BBH signal is located at the center of the segment.

Subsequently, we apply a windowing function to the BBH and blip time series data samples in order to reduce the spectral leakage due to discontinuity at the edges of the chopped segment. Windowing ensures a smooth and gradual transition of time series amplitude to zero at the edges. As the input to our network is the SGP maps, it is crucial to find an optimal combination of various factors that will maximize the projection of BBH signals and blips. As a result, the length of the data segment, the location of the signal or blip within the segment and the choice of windowing had to be tuned by performing several tests. For example, longer data segments result in increased noise in the SGP maps and shorter data segments also adversely affect the SGP maps due to the windowing process. Hence, we find that an intermediate length of two seconds is the most appropriate choice for the signal parameter ranges chosen for this work.

Once the time series data for blips and BBH are prepared, we proceed to generate the SGP maps for our sample. The first step is to calculate the PSDs. For each blip, we use the 16-sec-long data segment that contains the blip and apply the standard Welch's method to calculate the PSD. This method ensures that the presence of the blip in the segment does not affect the PSD calculation. For the BBH sample, we use the same 64-sec-long data segment that was used as the background noise for the BBH signals. We apply the same Welch's method to get the PSD required for the BBH sample.



FIG. 3. Here we show an example of a simulated BBH input sample fed to the network. The H1 (left) and L1 (right) SGP maps are placed adjacent to each other. Axis labels and color bars are omitted. Here, the SNR \approx 10 for H1 and \approx 9 for L1.

In the next step, we calculate the projection of the data (BBH and blips) on the sine-Gaussian waveform [see Eqs. (1) and (2)] and represent it using the 2D parameter space of $Q - f_0$ (see Sec. II for details) for each detector separately. The color in each pixel of an SGP map represents the strength of the projection.

Initially, we experimented with analyzing SGP maps of individual detectors separately for each blip and BBH signal. While the network performance was satisfactory at high SNRs, there was room for further improvement in the sensitivity of low-SNR GW events. We then considered analyzing maps from both detectors simultaneously. There is a clear advantage of working with multidetector data because a true astrophysical signal, such as that from a BBH, if sufficiently strong, will appear in both the detectors whereas nonastrophysical noise transients such as blips rarely coincide for two independent detectors. This is well evident for an example data of BBH (Fig. 3) and a blip (Fig. 4) showing both H1 and L1 detectors. One could consider including data from more than two detectors



FIG. 4. Similar to Fig. 3, except that a blip input sample is fed to the network here. The H1 (left) and L1 (right) SGP maps are kept adjacent to each other. We can see a blip is present in H1 (left) with an SNR ≈ 11 .



FIG. 5. Schematic diagram of our neural network. Each layer uses a relu activation function except for the last one, which uses a sigmoid function. The final output is the probability that the input image corresponds to a BBH signal. Here, "Conv k" stands for kth convolutional layer, "MaxP k" stands for kth max-pooling layer, "Flatten" denotes the flattening of 2D matrix, "Dropout(x)" layer is to tackle overfitting by dropping some of the neurons, and "Dense(k)" is fully connected layer with k neurons [47].

which may or may not be efficient in our end goal. However, we leave this exercise for future work and stick to using two-detector data.

In order to analyze SGP maps from both detectors simultaneously, we consider a multiview learning model of the CNN, e.g., [43]. In this form of learning, the network learns using multiple image representations of a given class of objects. The multiview model of learning increases the distinguishable features in the input images and the learning capability of the network as a whole. Multiview learning could be done in mainly two ways, parallel view and merged view. In parallel-view learning, images are fed to the network through more than one input channel and combined subsequently after passing through a few hidden layers. On the contrary, in merged-view learning, the images are placed adjacent to each other in a grid and then fed to the network through a single input channel.

In this study, we adopt the merged-view learning where the SGP maps, for each BBH or blip, created using the H1 and L1 detectors, are placed laterally as a single input image for the network. For example, see Figs. 3 and 4 where the images, without any labels and tick marks, are fed to the network.

We choose 10 000 SGP maps for training and 3000 SGP maps for validation where the blips and BBH signals are in equal proportion. For the test sample containing BBH signals, we produce two distinct samples as a function of mass and mass ratios. Each BBH sample comprises 12 000 SGP maps. We produce 6000 SGP maps of blips for the test sample. Additionally, we use the

sample of 49 real GW events taken from O1, O2, and O3a [5,6].

IV. DEEP LEARNING NETWORK

The goal of our study is to distinguish BBH from blip glitches. For this classification problem, we use a deep learning algorithm known as convolutional neural network [44–46]. We follow the supervised learning approach where we train the network on SGP maps belonging to both the BBH and blip classes. The trained network then gives us predictions whether a test SGP map contains a BBH or a blip.

We describe below the structure of the network developed in this work and summarize the training process leading to the final model that will make robust predictions. As mentioned in the previous section, the input data for the network are two SGP maps corresponding to H1 and L1 detectors for each BBH or blip. These maps are placed laterally such that their final dimension is 150×300 along with three color channels as seen by the network. Our network has four convolution blocks followed by a flattened layer and two fully connected blocks that are connected to an output layer (see Fig. 5). Each convolution block has a convolution layer and a pooling layer. We choose the relu (rectified linear unit) activation function in the convolution layers and a max-pooling option in the pooling layer. Max pooling has been shown to perform better in classification tasks where one has to deal with sharp features in the images in comparison to other options such as average pooling and min pooling. Subsequently, we introduce a flattened layer followed by the relu activation function. Next, in each of the fully connected blocks, we have a dropout and a dense layer. The dropout layer helps in preventing overfitting issues. Finally, in the output layer, we apply the sigmoid activation function to predict the classification probabilities. This layer has a single neuron that gives the probability that the input image belongs to BBH class. In order to determine the probability for the class of blips, we simply subtract the predicted probability from 1.

We choose binary cross-entropy to define the loss function of the network as it is appropriate for binary classification. We use the "adam" optimizer [48] with a learning rate set to 0.001 and other parameters set to their default values as they happened to be an adequate choice based on the tests we conducted. We use the metric "accuracy" to measure the network performance.

Our training sample includes 5000 maps of blips and BBHs each whereas the validation sample includes 1500 maps of blips and BBHs each. We train our network for 10 epochs and each epoch consists of the training samples divided into 300 batches. The validation sample is divided into 100 batches. During the training process, we use the "fit generator" method available in KERAS [49], which loads the data into primary memory in batches and feeds it to the network for training. As we do not load the whole training data at once into the memory, it reduces the load on the computing node and makes the training process quite fast.

V. EXISTING METHODS TO TACKLE BLIPS

Before testing the performance of our machine learning model and comparing it with existing methods for vetoing blips in GW data, we briefly summarize two of the main statistics in the GW data analysis that are currently employed for vetoing the blips, namely, traditional χ^2 [50], and sine-Gaussian χ^2 [14].

The traditional χ^2 is constructed by subdividing the template waveform triggered by the match-filtering process into *p* nonoverlapping frequency bins. Each bin contributes equally to the SNR of the best matching CBC template [50]. If the data segment **s** is adequately described as a Gaussian noise with an added CBC signal that shows a large correlation with the template **h**, this will follow a χ^2 distribution with 2p - 2 degrees of freedom. Most noise transients present in the GW data show a higher χ^2 in comparison to CBC signals, making it easy to distinguish signals vs noise artifacts. The definition for the traditional reduced χ^2 is

$$\chi_{\rm r}^2 = \frac{1}{2p - 2} \sum_{i=1}^p \|(\mathbf{s}|\mathbf{h}_i) - (\mathbf{h}_i|\mathbf{h}_i)\|^2,$$
(3)

where p is the number of bins, \mathbf{h}_i is the time domain representation of the waveform corresponding to the *i*th bin, and for any two data segments the inner product is defined by Eq. (2). The traditional χ^2 is then combined with the SNR to create a ranking statistic called reweighted SNR, which, as used in the PyCBC [51–54] analysis, is

$$\tilde{\rho} = \begin{cases} \rho & \text{for } \chi_{\rm r}^2 \le 1\\ \rho [\frac{1}{2} (1 + (\chi_{\rm r}^2)^3)]^{-\frac{1}{2}} & \text{for } \chi_{\rm r}^2 > 1 \end{cases}$$
(4)

The sine-Gaussian (SG) χ^2 takes advantage of the excess power that blips typically possess in higher frequency regions in comparison to CBC signals. It computes this excess power by utilizing a set of sine-Gaussian waveforms whose central frequency ranges from 30–120 Hz above the final frequency of the triggered template [14]. It is defined as

$$\chi^2_{\mathrm{r,sg}} = \frac{1}{2N} \sum_{i=1}^{N} \langle \mathbf{s} | \mathbf{g}_i \rangle^2, \qquad (5)$$

where \mathbf{g}_i is the *i*th sine-Gaussian waveform.

The reweighted SNR used in the case of SG χ^2 is given as

$$\rho_{sg}^{\sim} = \begin{cases} \tilde{\rho} & \text{for } \chi_{r,sg}^{2} \le 4 \\ \tilde{\rho}(\chi_{r,sg}^{2}/4)^{-\frac{1}{2}} & \text{for } \chi_{r,sg}^{2} > 4 \end{cases}.$$
 (6)

The number of bins p in the case of the traditional χ^2 and the number of sine-Gaussian waveforms in the case of the SG χ^2 are tuned empirically based on the distribution of single detector background triggers in engineering data.

VI. PERFORMANCE AND RESULTS

After training our network on a sample of real blips and simulated BBH signals, we obtain a final model that can make predictions on a test sample. In this section, we evaluate our model's performance using some of the standard metrics, for example, the receiver-operating characteristic (ROC) curves and F1 score. We also present model predictions for the real GW events.

TABLE I. Details of the test sample. Whereas blips are divided based on their SNRs alone, the BBH samples are divided based on their SNRs and a mass parameter, which is taken to be either their total mass or their mass ratio. The total number of unique samples for each type is given in the last column.

	SNR		
Туре	[3,8]	[8,15]	Total
Blips	3000	3000	6000
BBH1			12 000
$M_{\text{total}} \in [20, 80] M_{\odot}$	3000	3000	
$M_{\text{total}} \in [80, 140] M_{\odot}$	3000	3000	
BBH2			12 000
$q \in [1, 2]$	3000	3000	
$q \in [2, 10]$	3000	3000	
GWTC-2 and GWTC-3			81
Other marginal and OGC4 events			34



FIG. 6. Comparison of ROC curves for CNN with those for two other methods, namely the traditional χ^2 [50] and the sine-Gaussian χ^2 [14]. The curves are shown for low- and high-SNR bins (solid and dotted, respectively). On the left, we show the performance at lowmass BBH signals, which is similar to that of the high-mass BBH signals on the right. The gray shaded region has more uncertainty for the low SNRs (dotted curve) due to limited sample size. The solid circle marks the TPP-FPP for a threshold of 0.8.

First, we describe how the test sample is prepared. We select a sample of 6000 blips that are divided equally into two SNR bins. For the BBHs, we generate two distinct samples such that one has a uniform distribution in total mass M_{total} and the other is uniform in mass-ratio q (which is defined as the ratio of the primary mass to the secondary mass). Each of the BBH samples is divided equally into two bins as per the total mass or the mass-ratio. These subsamples are further divided into two SNR bins. The breakdown of the full test sample according to the bins, how the bins are defined and the unique number of test samples used in our analysis are given in Table I. We also list the number of GW events tested in this work.

Next, we test the performance of our network in each bin for BBH versus blips. To generate the ROC curves, we calculate the true-positive probability (TPP) and falsepositive probability (FPP) for varying thresholds of detection.

$$TPP = \frac{True \ Positive}{True \ Positive + False \ Negative}$$
(7)

$$FPP = \frac{False Positive}{False Positive + True Negative}$$
(8)

The resulting ROC curves for BBH1 sample split by their total mass are shown in Fig. 6 and for BBH2 sample split by their mass ratio are shown in Fig. 7. As expected, the network performs better at high SNRs (solid curves) than low SNRs for both BBH1 and BBH2 samples. The gray band roughly shows the region where we have limited samples in low SNR bins that increase the uncertainties on the ROC. Hence, for the low SNRs, we focus on FPP > 0.01. Interestingly, the ROC curves are only marginally different across the bins in total mass or the mass ratio. It is encouraging to see that the network is robust and performs equally well in the total mass or mass-ratio ranges we explore here.

We also compare the performance of the neural network classifier with existing methods to veto blips mentioned in Sec. V, namely traditional χ^2 and SG χ^2 . We can see from the ROC curves that our network performs better than the traditional χ^2 and SG χ^2 , which are used here in their network form calculated as quadrature sum of their value in H1 and L1 detectors. Although, in current LIGO search pipelines these statistics are applied in modified form with improved modeling of background distribution [53,55,56]. The performance of our neural network classifier is better at higher total masses and higher mass ratios in comparison to the lower bins. In all cases, the CNN performs either better or comparable to traditional χ^2 and SG χ^2 . It is important to highlight the improvement in sensitivity with our network at low SNRs for high total mass and/or high mass-ratio BBHs, where the existing methods are known to have a poor performance. At high masses, our network shows 75% increment in the TPP than the traditional χ^2 and SG χ^2 at an FPP of 10^{-2} and for high mass ratios, the neural network shows around 50% improvement in TPP at an FPP of 10^{-2} compared to other methods.

We use another metric called F1 score to assess the efficiency of our network. F1 score is defined as the harmonic mean of precision and recall of a classifier.

$$F1 \text{ score} = \frac{2 * \text{ precision} * \text{ recall}}{\text{ precision} + \text{ recall}}, \qquad (9)$$



FIG. 7. Comparison of ROC curves for CNN with those for two other methods. Same as Fig. 6 except the performance is shown for low and high mass ratios of BBH signals (left and right, respectively).

where

$$precision = \frac{TP}{TP + FP},$$
 (10)

and

$$\operatorname{recall} = \frac{\mathrm{TP}}{\mathrm{FN} + \mathrm{TP}}.$$
 (11)

Here, TP is true positive, FP is false positive, and FN is false negative. The F1 score is calculated for the two SNR and two total mass bins. The F1 score values are shown in Fig. 8. As we can see, the network performs well both at low and high SNRs. The performance seems to be almost independent of the total mass of the binary.

We also test our network on the real events from GWTC-3 and the results are shown in Fig. 9. Four events



FIG. 8. F1 score of our neural network for a threshold of 0.8. Scores are shown for low and high total mass of the BBH signals (x axis) further split by their SNRs (y axis).

are not included in our sample, as one of the detector (H1 or L1) data is not available for these events. The network predicts 95% GW events correctly as BBH signals. Among the high SNRs, the event GW170817 is misclassified, which is a binary neutron star merger signal with a glitch overlapping in the L1 detector. Among the low SNR events, GW191219_163120, GW200308_173609, and GW200220_124850 are not classified as BBH. Upon inspection of the SGP maps, we find that there is a hint of projection in one of the detectors, but the projection in the other detector is not very clean due to low SNR. This is consistently seen in the maps of all of the three events, which is the most likely explanation for the misclassification.

In Table II, we show the computation times for analyzing GW data with our network. The training and classification of several thousand samples by the network



FIG. 9. Detection efficiency of our network on the real BBH events from GWTC-3. The format is the same as in Fig. 8 and the threshold is also set to 0.8.

TABLE II. Computation time for various sections of SIGMA-NET. All times are computed for a single core of Intel Xeon Gold 6142 CPU. Multiprocessing and parallelizing further reduces the computation times. For example, parallelizing on k processors reduces all computational costs by a factor of k.

Operation	Time
Generating 1 SGP map	30 sec
Training of the network on 10 000 samples	20 min
Classification of 3000 SGP maps	90 sec

is quite fast even on a single core. Although the generation of SGP maps is done using multiple cores, which makes it reasonably fast, there is further room for improvement in the process of generating a single SGP map. The latter process executes on a single core currently. We anticipate that through some changes, for example, parallellization of certain loops and by reducing the data sampling rate from 4096 to 2048 Hz without affecting our results, we can speed up the generation of an SGP map by factor of a few.

VII. CONCLUSION AND FUTURE PROSPECTS

Blip glitches are known to have similar time-frequency characteristics to those of BBH signals of high total mass. Thus, the standard pipelines find it difficult to distinguish high-mass BBH signals from these glitches. We use sine-Gaussian projection (SGP) maps, a new way of visualizing GW data, to discriminate BBH signals from blips. We create the SGP maps by projecting the BBH signals and blips on the two-dimensional parameter space of sine-Gaussian waveforms. These maps are complementary to the continuous-wave transform maps commonly used in GW data analyses.

In order to demonstrate the usefulness of SGP maps we develop a deep learning framework that uses convolutional neural network to classify BBH signals and blips. It can help in reducing the ambiguity between CBC and blip triggers generated from LIGO-Virgo runs. We compare our method with those used in the standard pipelines for identifying BBH signals in the LIGO-Virgo GW data. We find that our network performs consistently much better than traditional χ^2 and sine-Gaussian χ^2 for BBH signals with SNR > 8 and a total mass $\in [20, 140]M_{\odot}$ where the mass ratio goes from 1 to 10. For low SNRs ($\in [3, 8]$) too, the network performs significantly better when FPP > 0.01. However, we notice that below FPP of ≈ 0.01 , the limited size of our training sample does not allow characterization of our network robustly as we do not have sufficient numbers of blips in real data from first two observing runs. This will not be a limiting factor in our subsequent work where we plan to include data from third observing run as well. We also find that our network is able to correctly identify 77 out of 83 CBC events from LIGO-Virgo's O1, O2, and O3 observing runs. Our network correctly identifies 15 out of the 19 LIGO marginal events [8] and all of the 15 GW events from the fourth open gravity catalog (4-OGC) [57]. An immediate application where our network can contribute to BBH searches is one where it scrutinizes triggers judged as blips, e.g., by other tools such as the one in Refs. [18,43]. In a subset of those cases, SIGMA-NET would lend support to such a categorization. On the other hand, if its support is weak for the hypothesis that a trigger is a blip (as opposed to a BBH signal), then it would be flagged for further analysis by signal-specific search pipelines (e.g., those employing matched filtering of the corresponding data snippet). The network output can also be used to weight the SNR of the triggers. It will reduce the background due to blips and hence improve the search sensitivity.

Broader varieties of glitches have been targeted in ML networks in Refs. [27,58]. The recent work in Ref. [58] shows how a χ^2 statistic can be tuned with the help of an ML network for a wide variety of glitches to perform the vetoing task optimally. In this instance, the ML network helps in choosing the time and frequency shifts that, when applied to the triggered CBC template, yield the basis vectors for the χ^2 [22]. That basis apply collects the power of glitches in the region where the CBC signals are expected to have negligible projections. The improvement in performance of that network, however, is not as high as what was achieved in this work (for blips). Nevertheless, in future it will be interesting to study if ideas from that work can be utilized to broaden the target glitch categories without hurting the performance improvement much. We could also consider including data from additional detectors (e.g., Virgo and KAGRA). Inclusion of multiple SGP maps for the same BBH will most likely improve the sensitivity of identifying them but at the cost of increased computational times, which results from the additional data that needs to be processed. However, more detector data also increases the probability of finding accidentally coincident glitches. This could then adversely affect the performance of the neural network. Thus, one will have to perform a riskbenefit exercise when deciding on using data from additional detectors.

In this work we focused on demonstrating the usefulness of the SGP maps by using them to distinguish blip glitches from BBH signals. In the future, with the inclusion of various artifact categories—including the no-signal and noglitch categories—in the training, SIGMA-NET could be developed to potentially analyze GW data in low-latency searches in order to generate an initial set of triggers as candidate BBH events where the speed would be significantly higher than the match-filtering based template-bank search.

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- B. Abbott, R. Abbott, T. Abbott, M. Abernathy, F. Acernese, K. Ackley, C. Adams, T. Adams, P. Addesso, R. Adhikari *et al.*, Phys. Rev. Lett. **116**, 061102 (2016).
- [2] B. P. Abbott, R. Abbott, T. D. Abbott, M. R. Abernathy, F. Acernese, K. Ackley, M. Adamo, C. Adams, T. Adams, P. Addesso *et al.*, Classical Quantum Gravity **33**, 134001 (2016).
- [3] J. Aasi, B. P. Abbott, R. Abbott, T. Abbott, M. R. Abernathy, K. Ackley, C. Adams, T. Adams, P. Addesso *et al.*, Classical Quantum Gravity **32**, 074001 (2015).
- [4] F. Acernese et al., J. Phys. Conf. Ser. 610, 012014 (2015).
- [5] B. Abbott, R. Abbott, T. Abbott, S. Abraham, F. Acernese, K. Ackley, C. Adams, R. Adhikari, V. Adya, C. Affeldt *et al.*, Phys. Rev. X 9, 031040 (2019).
- [6] R. Abbott, T. Abbott, S. Abraham, F. Acernese, K. Ackley, A. Adams, C. Adams, R. Adhikari, V. Adya, C. Affeldt *et al.*, Phys. Rev. X **11**, 021053 (2021).
- [7] R. Abbott, T. D. Abbott, F. Acernese, K. Ackley, C. Adams, N. Adhikari, R. X. Adhikari, V. B. Adyan *et al.* (T. L. S. and Virgo Collaborations), arXiv:2108.01045.
- [8] R. Abbott, T. Abbott, S. Abraham, F. Acernese, K. Ackley, A. Adams, C. Adams, R. Adhikari, V. Adya, C. Affeldt *et al.*, arXiv:2111.03606.
- [9] T. Venumadhav, B. Zackay, J. Roulet, L. Dai, and M. Zaldarriaga, Phys. Rev. D 101, 083030 (2020).
- [10] A. H. Nitz, T. Dent, G. S. Davies, S. Kumar, C. D. Capano, I. Harry, S. Mozzon, L. Nuttall, A. Lundgren, and M. Tápai, Astrophys. J. 891, 123 (2020).
- [11] KAGRA Collaboration, Nat. Astron. 3, 35 (2019).
- [12] M. Saleem, J. Rana, V. Gayathri, A. Vijaykumar, S. Goyal, S. Sachdev, J. Suresh, S. Sudhagar, A. Mukherjee, G. Gaur, B. Sathyaprakash, A. Pai, R. X. Adhikari, P. Ajith, and S. Bose, Classical Quantum Gravity **39**, 025004 (2022).
- [13] B. P. Abbott, R. Abbott, T. D. Abbott, S. Abraham, F. Acernese, K. Ackley, C. Adams, V. B. Adya, C. Affeldt, M. Agathos *et al.*, Classical Quantum Gravity **37**, 055002 (2020).
- [14] A. H. Nitz, Classical Quantum Gravity 35, 035016 (2018).
- [15] B. P. Abbott, R. Abbott, T. D. Abbott, M. R. Abernathy, F. Acernese, K. Ackley, C. Adams, T. Adams, P. Addesso, R. X. Adhikari *et al.*, Classical Quantum Gravity 35, 065010 (2018).

- [16] M. Cabero, A. Lundgren, A. H. Nitz, T. Dent, D. Barker, E. Goetz, J. S. Kissel, L. K. Nuttall, P. Schale, R. Schofield, and D. Davis, Classical Quantum Gravity 36, 155010 (2019).
- [17] M. Zevin, S. Coughlin, S. Bahaadini, E. Besler, N. Rohani, S. Allen, M. Cabero, K. Crowston, A. K. Katsaggelos, S. L. Larson, T. K. Lee, C. Lintott, T. B. Littenberg, A. Lundgren, C. Østerlund, J. R. Smith, L. Trouille, and V. Kalogera, Classical Quantum Gravity 34, 064003 (2017).
- [18] S. Bahaadini, V. Noroozi, N. Rohani, S. Coughlin, M. Zevin, J. Smith, V. Kalogera, and A. Katsaggelos, Inf. Sci. 444, 172 (2018).
- [19] S. Coughlin, S. Bahaadini, N. Rohani, M. Zevin, O. Patane, M. Harandi, C. Jackson, V. Noroozi, S. Allen, J. Areeda *et al.*, Phys. Rev. D **99**, 082002 (2019).
- [20] S. Soni, C. P. L. Berry, S. B. Coughlin, M. Harandi, C. B. Jackson, K. Crowston, C. Østerlund, O. Patane, A. K. Katsaggelos, L. Trouille *et al.*, Classical Quantum Gravity 38, 195016 (2021).
- [21] T.D. Canton, S. Bhagwat, S.V. Dhurandhar, and A. Lundgren, Classical Quantum Gravity 31, 015016 (2014).
- [22] S. Bose, S. Dhurandhar, A. Gupta, and A. Lundgren, Phys. Rev. D 94, 122004 (2016).
- [23] D. Davis, J. S. Areeda, B. K. Berger, R. Bruntz, A. Effler, R. C. Essick, R. P. Fisher, P. Godwin, E. Goetz, and A. F. Helmling-Cornell, Classical Quantum Gravity 38, 135014 (2021).
- [24] P. Joshi, R. Dhurkunde, S. Dhurandhar, and S. Bose, Phys. Rev. D 103, 044035 (2021).
- [25] D. Davis, L. V. White, and P. R. Saulson, Classical Quantum Gravity 37, 145001 (2020).
- [26] H. Gabbard, M. Williams, F. Hayes, and C. Messenger, Phys. Rev. Lett. **120**, 141103 (2018).
- [27] S. Jadhav, N. Mukund, B. Gadre, S. Mitra, and S. Abraham, Phys. Rev. D 104, 064051 (2021).
- [28] S. Kapadia, T. Dent, and T. Dal Canton, Phys. Rev. D 96, 104015 (2017).
- [29] T. D. Gebhard, N. Kilbertus, I. Harry, and B. Schölkopf, Phys. Rev. D 100, 063015 (2019).
- [30] D. George and E. A. Huerta, Phys. Rev. D 97, 044039 (2018).
- [31] N. Mukund, S. Abraham, S. Kandhasamy, S. Mitra, and N. S. Philip, Phys. Rev. D 95, 104059 (2017).

- [32] Z. Que, E. Wang, U. Marikar, E. Moreno, J. Ngadiuba, H. Javed, B. Borzyszkowski, T. Aarrestad, V. Loncar, S. Summers et al., in 2021 IEEE 32nd International Conference on Application-specific Systems, Architectures and Processors (ASAP) (2021), 10.1109/asap52443.2021.00025.
- [33] A. Schmitt, K. Fu, S. Fan, and Y. Luo, in Proceedings of the 2nd International Conference on Computer Science and Software Engineering, CSSE 2019, (2019), p. 73, 10.1145/ 3339363.3339377.
- [34] M. B. Schäfer, F. Ohme, and A. H. Nitz, Phys. Rev. D 102, 063015 (2020).
- [35] D. George and E. A. Huerta, Phys. Rev. D 97, 044039 (2018).
- [36] J. Yan, M. Avagyan, R. E. Colgan, D. Veske, I. Bartos, J. Wright, Z. Márka, and S. Márka, Phys. Rev. D 105, 043006 (2022).
- [37] H. Gabbard, M. Williams, F. Hayes, and C. Messenger, Phys. Rev. Lett. **120**, 141103 (2018).
- [38] F. Robinet, N. Arnaud, N. Leroy, A. Lundgren, D. Macleod, and J. McIver, SoftwareX 12, 100620 (2020).
- [39] S. Dhurandhar, A. Gupta, B. Gadre, and S. Bose, Phys. Rev. D 96, 103018 (2017).
- [40] K. G. Arun, A. Buonanno, G. Faye, and E. Ochsner, Phys. Rev. D 79, 104023 (2009).
- [41] Rich Abbott, Thomas D. Abbott, Sheelu Abraham, Fausto Acernese, Kendall Ackley, Carl Adams, Rana X. Adhikari, Vaishali B. Adya, Christoph Affeldt, Michalis Agathos *et al.*, SoftwareX 13, 100658 (2021).
- [42] M. Hannam, P. Schmidt, A. Bohé, L. Haegel, S. Husa, F. Ohme, G. Pratten, and M. Pürrer, Phys. Rev. Lett. 113, 151101 (2014).

- [43] S. Bahaadini, N. Rohani, S. Coughlin, M. Zevin, V. Kalogera, and A. K. Katsaggelos, arXiv:1705.00034.
- [44] Y. LeCun and Y. Bengio, in *Handbook of Brain Theory and Neural Networks*, edited by M. A. Arbib (MIT Press, Cambridge, MA, 1995), p. 3361.
- [45] Z. Li, W. Yang, S. Peng, and F. Liu, arXiv:2004.02806.
- [46] J. Schmidhuber, Neural Netw. 61, 85 (2015).
- [47] Basic CNN architecture, https://www.upgrad.com/blog/ basic-cnn-architecture/ (2022).
- [48] D. P. Kingma and J. Ba, arXiv:1412.6980.
- [49] F. Chollet et al., KERAS (2015), https://keras.io.
- [50] B. Allen, Phys. Rev. D 71, 062001 (2005).
- [51] A. Nitz *et al.*, GWASTRO/PyCBC, Zenodo, 10.5281/zenodo.5347736 (2021).
- [52] S. A. Usman, A. H. Nitz, I. W. Harry, C. M. Biwer, D. A. Brown, M. Cabero, C. D. Capano, T. D. Canton, T. Dent, S. Fairhurst *et al.*, Classical Quantum Gravity **33**, 215004 (2016).
- [53] G. S. Davies, T. Dent, M. Tápai, I. Harry, C. McIsaac, and A. H. Nitz, Phys. Rev. D 102, 022004 (2020).
- [54] B. Allen, W. G. Anderson, P. R. Brady, D. A. Brown, and J. D. E. Creighton, Phys. Rev. D 85, 122006 (2012).
- [55] A. H. Nitz, T. Dent, T. Dal Canton, S. Fairhurst, and D. A. Brown, Astrophys. J. 849, 118 (2017).
- [56] A. H. Nitz, T. Dal Canton, D. Davis, and S. Reyes, Phys. Rev. D 98, 024050 (2018).
- [57] A. H. Nitz, S. Kumar, Y.-F. Wang, S. Kastha, S. Wu, M. Schäfer, R. Dhurkunde, and C. D. Capano, arXiv:2112 .06878.
- [58] C. McIsaac and I. Harry, Phys. Rev. D **105**, 104056 (2022).