

**Erratum: QCD description of backward vector meson
hard electroproduction**
[Phys. Rev. D **91**, 094006 (2015)]

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(Received 4 October 2022; published 3 November 2022)

DOI: 10.1103/PhysRevD.106.099901

We correct a sign mistake in our calculation of backward vector meson electroproduction amplitude within the transition distribution amplitude (TDA) framework presented in Table I. The same error has also propagated into Table 3 of Ref. [1]. We present the corrected version of this table here. The altered signs are marked in boldface.

A useful cross check for this result is provided by the calculation of the integral convolutions $\mathcal{I}^{(k)}$, $k = 1, \dots, 6$, defined in Eq. (24) for the $\gamma^* p \rightarrow V^0 p$ reaction within the cross channel nucleon exchange model for VN TDAs considered in Appendix B. In this model, the convolution integrals $\mathcal{I}^{(k)}$ turn out to be proportional to the pQCD expression for the proton e.m. form factor F_1^p (see discussion in Ref. [2]):

TABLE I. 14 of the 21 diagrams contributing to the hard-scattering amplitude with their associated coefficient $T_\alpha^{(k)} \equiv D_\alpha \times N_\alpha^{(k)}$ (no summation over α assumed). The seven first ones with u -quark lines inverted are not drawn. The crosses represent the virtual-photon vertex.

α	Diagram D_α	Numerators
1		$N_\alpha^{(1)}$
		$N_\alpha^{(2)}$
		$N_\alpha^{(3)}$
		$N_\alpha^{(4)}$
	$\frac{Q_u(2\xi)^2}{(2\xi-x_1-i\epsilon)^2(x_3-i\epsilon)(1-y_1)^2y_3}$	$N_\alpha^{(5)}$
	$\frac{Q_u(2\xi)^2}{(2\xi-x_1-i\epsilon)^2(x_3-i\epsilon)(1-y_1)^2y_3}$	$N_\alpha^{(6)}$
2		$N_\alpha^{(1)}$
		$N_\alpha^{(2)}$
		$N_\alpha^{(3)}$
		$N_\alpha^{(4)}$
		$N_\alpha^{(5)}$
		$N_\alpha^{(6)}$
3		$N_\alpha^{(1)}$
		$N_\alpha^{(2)}$
		$N_\alpha^{(3)}$
		$N_\alpha^{(4)}$
	$\frac{Q_u(2\xi)^2}{(x_1-i\epsilon)(2\xi-x_2-i\epsilon)(x_3-i\epsilon)y_1(1-y_1)y_3}$	$N_\alpha^{(5)}$
	$\frac{Q_u(2\xi)^2}{(x_1-i\epsilon)(2\xi-x_2-i\epsilon)(x_3-i\epsilon)y_1(1-y_1)y_3}$	$N_\alpha^{(6)}$

(Table continued)

TABLE I. (Continued)

α	Diagram D_α		Numerators
4		$N_\alpha^{(1)}$	$-(V^P - A^P)(V_{1\mathcal{E}}^{VN} - A_{1\mathcal{E}}^{VN})$
		$N_\alpha^{(2)}$	$-(V^P - A^P)(V_{1n}^{VN} - A_{1n}^{VN})$
		$N_\alpha^{(3)}$	$-(V^P - A^P)(V_{1T}^{VN} - A_{1T}^{VN} + V_{2\mathcal{E}}^{VN} - A_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(4)}$	$-(V^P - A^P)(V_{2T}^{VN} - A_{2T}^{VN})$
		$N_\alpha^{(5)}$	$(V^P - A^P)(V_{2\mathcal{E}}^{VN} - A_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(6)}$	$-(V^P - A^P)(V_{2n}^{VN} - A_{2n}^{VN})$
5		$N_\alpha^{(1)}$	$(V^P + A^P)(V_{1\mathcal{E}}^{VN} + A_{1\mathcal{E}}^{VN})$
		$N_\alpha^{(2)}$	$(V^P + A^P)(V_{1n}^{VN} + A_{1n}^{VN})$
		$N_\alpha^{(3)}$	$(V^P + A^P)(V_{1T}^{VN} + A_{1T}^{VN} + V_{2\mathcal{E}}^{VN} + A_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(4)}$	$(V^P + A^P)(V_{2T}^{VN} + A_{2T}^{VN})$
		$N_\alpha^{(5)}$	$-(V^P + A^P)(V_{2\mathcal{E}}^{VN} + A_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(6)}$	$+(V^P + A^P)(V_{2n}^{VN} + A_{2n}^{VN})$
6		$N_\alpha^{(1)}$	0
		$N_\alpha^{(2)}$	0
		$N_\alpha^{(3)}$	0
		$N_\alpha^{(4)}$	0
		$N_\alpha^{(5)}$	0
		$N_\alpha^{(6)}$	0
7		$N_\alpha^{(1)}$	$-2(V^P V_{1\mathcal{E}}^{VN} + A^P A_{1\mathcal{E}}^{VN})$
		$N_\alpha^{(2)}$	$-2(V^P V_{1n}^{VN} + A^P A_{1n}^{VN})$
		$N_\alpha^{(3)}$	$-2(V^P (V_{1T}^{VN} + V_{2\mathcal{E}}^{VN}) + A^P (A_{1T}^{VN} + A_{2\mathcal{E}}^{VN}))$
		$N_\alpha^{(4)}$	$-2(V^P V_{2T}^{VN} + A^P A_{2T}^{VN})$
		$N_\alpha^{(5)}$	$2(V^P V_{2\mathcal{E}}^{VN} + A^P A_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(6)}$	$-2(V^P V_{2n}^{VN} + A^P A_{2n}^{VN})$
8		$N_\alpha^{(1)}$	0
		$N_\alpha^{(2)}$	0
		$N_\alpha^{(3)}$	0
		$N_\alpha^{(4)}$	0
		$N_\alpha^{(5)}$	0
		$N_\alpha^{(6)}$	0
9		$N_\alpha^{(1)}$	$-(V^P - A^P)(V_{1\mathcal{E}}^{VN} - A_{1\mathcal{E}}^{VN}) + 2T^P(T_{1\mathcal{E}}^{VN} + T_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(2)}$	$-(V^P - A^P)(V_{1n}^{VN} - A_{1n}^{VN}) + 4T^P(T_{1n}^{VN} + \frac{\Delta_1^2}{2M^2} T_{4n}^{VN})$
		$N_\alpha^{(3)}$	$-(V^P - A^P)(V_{1T}^{VN} - A_{1T}^{VN} + V_{2\mathcal{E}}^{VN} - A_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(4)}$	$+4T^P(T_{1T}^{VN} + T_{3\mathcal{E}}^{VN} + \frac{\Delta_1^2}{2M^2} T_{4T}^{VN})$
		$N_\alpha^{(5)}$	$-(V^P - A^P)(V_{2T}^{VN} - A_{2T}^{VN}) + 2T^P(T_{2T}^{VN} + T_{3T}^{VN})$
		$N_\alpha^{(6)}$	$(V^P - A^P)(V_{2\mathcal{E}}^{VN} - A_{2\mathcal{E}}^{VN}) - 2T^P(T_{3\mathcal{E}}^{VN} - T_{4\mathcal{E}}^{VN})$
		$N_\alpha^{(6)}$	$-(V^P - A^P)(V_{2n}^{VN} - A_{2n}^{VN}) + 2T^P(T_{2n}^{VN} + T_{3n}^{VN})$

(Table continued)

TABLE I. (Continued)

α	Diagram D_α		Numerators
10	$\frac{Q_d(2\xi)^2}{(x_1 - i\epsilon)(2\xi - x_2 - i\epsilon)^2 y_1 (1 - y_2)^2}$	$N_\alpha^{(1)}$	$-(V^P - A^P)(V_{1\mathcal{E}}^{VN} - A_{1\mathcal{E}}^{VN}) + 2T^P(T_{1\mathcal{E}}^{VN} + T_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(2)}$	$-(V^P - A^P)(V_{1n}^{VN} - A_{1n}^{VN}) + 4T^P(T_{1n}^{VN} + \frac{\Delta_7^2}{2M^2} T_{4n}^{VN})$
		$N_\alpha^{(3)}$	$-(V^P - A^P)(V_{1T}^{VN} - A_{1T}^{VN} + V_{2\mathcal{E}}^{VN} - A_{2\mathcal{E}}^{VN})$
			$+4T^P(T_{1T}^{VN} + T_{3\mathcal{E}}^{VN} + \frac{\Delta_7^2}{2M^2} T_{4T}^{VN})$
		$N_\alpha^{(4)}$	$-(V^P - A^P)(V_{2T}^{VN} - A_{2T}^{VN}) + 2T^P(T_{2T}^{VN} + T_{3T}^{VN})$
		$N_\alpha^{(5)}$	$(V^P - A^P)(V_{2\mathcal{E}}^{VN} - A_{2\mathcal{E}}^{VN}) - 2T^P(T_{3\mathcal{E}}^{VN} - T_{4\mathcal{E}}^{VN})$
11		$N_\alpha^{(6)}$	$-(V^P - A^P)(V_{2n}^{VN} - A_{2n}^{VN}) + 2T^P(T_{2n}^{VN} + T_{3n}^{VN})$
		$N_\alpha^{(1)}$	0
		$N_\alpha^{(2)}$	0
		$N_\alpha^{(3)}$	0
		$N_\alpha^{(4)}$	0
		$N_\alpha^{(5)}$	0
12		$N_\alpha^{(6)}$	0
		$N_\alpha^{(1)}$	$(V^P + A^P)(V_{1\mathcal{E}}^{VN} + A_{1\mathcal{E}}^{VN})$
		$N_\alpha^{(2)}$	$(V^P + A^P)(V_{1n}^{VN} + A_{1n}^{VN})$
		$N_\alpha^{(3)}$	$(V^P + A^P)(V_{1T}^{VN} + A_{1T}^{VN} + V_{2\mathcal{E}}^{VN} + A_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(4)}$	$(V^P + A^P)(V_{2T}^{VN} + A_{2T}^{VN})$
		$N_\alpha^{(5)}$	$-(V^P + A^P)(V_{2\mathcal{E}}^{VN} + A_{2\mathcal{E}}^{VN})$
13	$\frac{Q_d(2\xi)^2}{(x_1 - i\epsilon)(2\xi - x_1 - i\epsilon)(x_2 - i\epsilon)y_1(1 - y_2)y_2}$	$N_\alpha^{(6)}$	$+(V^P + A^P)(V_{2n}^{VN} + A_{2n}^{VN})$
		$N_\alpha^{(1)}$	$2T^P(T_{1\mathcal{E}}^{VN} + T_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(2)}$	$4T^P(T_{1n}^{VN} + \frac{\Delta_7^2}{2M^2} T_{4n}^{VN})$
		$N_\alpha^{(3)}$	$4T^P(T_{1T}^{VN} + T_{3\mathcal{E}}^{VN} + \frac{\Delta_7^2}{2M^2} T_{4T}^{VN})$
		$N_\alpha^{(4)}$	$2T^P(T_{2T}^{VN} + T_{3T}^{VN})$
		$N_\alpha^{(5)}$	$-2T^P(T_{3\mathcal{E}}^{VN} - T_{4\mathcal{E}}^{VN})$
14	$\frac{Q_d(2\xi)^2}{(x_1 - i\epsilon)(2\xi - x_1 - i\epsilon)(x_2 - i\epsilon)y_1 y_2 (1 - y_3)}$	$N_\alpha^{(6)}$	$+2T^P(T_{2n}^{VN} + T_{3n}^{VN})$
		$N_\alpha^{(1)}$	$(V^P - A^P)(V_{1\mathcal{E}}^{VN} - A_{1\mathcal{E}}^{VN})$
		$N_\alpha^{(2)}$	$(V^P - A^P)(V_{1n}^{VN} - A_{1n}^{VN})$
		$N_\alpha^{(3)}$	$(V^P - A^P)(V_{1T}^{VN} - A_{1T}^{VN} + V_{2\mathcal{E}}^{VN} - A_{2\mathcal{E}}^{VN})$
		$N_\alpha^{(4)}$	$(V^P - A^P)(V_{2T}^{VN} - A_{2T}^{VN})$
		$N_\alpha^{(5)}$	$-(V^P - A^P)(V_{2\mathcal{E}}^{VN} - A_{2\mathcal{E}}^{VN})$
	$N_\alpha^{(6)}$	$+(V^P - A^P)(V_{2n}^{VN} - A_{2n}^{VN})$	

$$\begin{aligned}
\mathcal{I}^{(1)}(\xi, \Delta^2)|_{N(940)} &= -\mathcal{I}_0 \frac{K_{1\mathcal{E}}^{VN}(\xi, \Delta^2, G_{VNN}^V, G_{VNN}^T)}{2\xi}; \\
\mathcal{I}^{(2)}(\xi, \Delta^2)|_{N(940)} &= -\mathcal{I}_0 \frac{K_{1n}^{VN}(\xi, \Delta^2, G_{VNN}^V, G_{VNN}^T)}{2\xi}; \\
\mathcal{I}^{(3)}(\xi, \Delta^2)|_{N(940)} &= -\mathcal{I}_0 \frac{K_{1T}^{VN}(\xi, \Delta^2, G_{VNN}^V, G_{VNN}^T) + K_{2\mathcal{E}}^{VN}(\xi, \Delta^2, G_{VNN}^V, G_{VNN}^T)}{2\xi}; \\
\mathcal{I}^{(4)}(\xi, \Delta^2)|_{N(940)} &= -\mathcal{I}_0 \frac{K_{2T}^{VN}(\xi, \Delta^2, G_{VNN}^V, G_{VNN}^T)}{2\xi}; \\
\mathcal{I}^{(5)}(\xi, \Delta^2)|_{N(940)} &= +\mathcal{I}_0 \frac{K_{2\mathcal{E}}^{VN}(\xi, \Delta^2, G_{VNN}^V, G_{VNN}^T)}{2\xi}; \\
\mathcal{I}^{(6)}(\xi, \Delta^2)|_{N(940)} &= -\mathcal{I}_0 \frac{K_{2n}^{VN}(\xi, \Delta^2, G_{VNN}^V, G_{VNN}^T)}{2\xi}.
\end{aligned} \tag{1}$$

Here G_{VNN}^V and G_{VNN}^T stand for the vector and tensor couplings of vector mesons to nucleons, and \mathcal{I}_0 is a well-known convolution of nucleon DAs with hard scattering kernel occurring in the leading order perturbative QCD description of the proton electromagnetic form factor $F_1^p(Q^2)$ [3]:

$$Q^4 F_1^p(Q^2) = \frac{(4\pi\alpha_s)^2 f_N^2}{54} \mathcal{I}_0. \tag{2}$$

The explicit expressions for the functions $K_{1\mathcal{E}}^{VN}$, $K_{2\mathcal{E}}^{VN}$, K_{1T}^{VN} , K_{2T}^{VN} , are summarized in Appendix B.

The modifications of the convolution integrals do not considerably change our cross section estimates within the cross channel nucleon exchange model for ωN , ρN , and ϕN TDAs. Therefore, our predictions stay consistent with those employed in Ref. [4]. As an example we show the replacement of Fig. 3 as Fig. 1 here presenting the unpolarized cross section $\frac{d^2\sigma}{d\Omega_\omega}$ (in nb/sr) for backward $\gamma^* + p \rightarrow p + \omega$ for fixed $W = 3.20$ GeV as a function of Q^2 in the u -channel nucleon exchange model for ωN TDAs employing the Bonn 2000 set [6] for the $G_{\omega NN}^{V,T}$ couplings.

Let us also report a misprint in the last line of Eq. (23). The tensor structure $\mathcal{S}_{s_1 s_2}^{(6)\lambda_\gamma \lambda_V}$ must read

$$\mathcal{S}_{s_1 s_2}^{(6)\lambda_\gamma \lambda_V} = (\mathcal{E}^*(p_V, \lambda_V) \cdot n) \bar{U}(p_2, s_2) \hat{\epsilon}(q, \lambda_\gamma) \hat{\Delta}_T U(p_1, s_1). \tag{3}$$

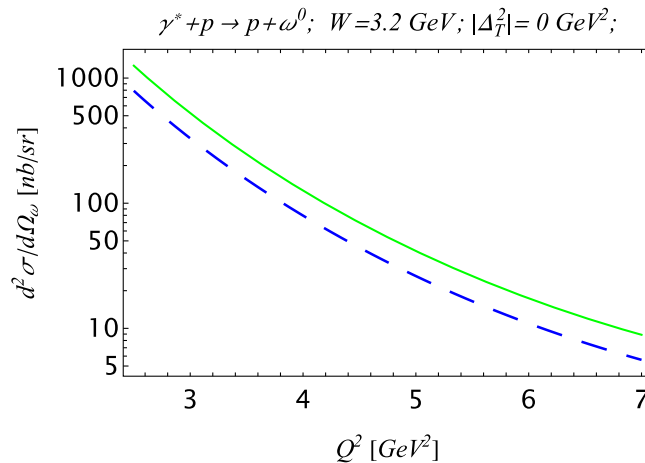


FIG. 1. Unpolarized cross section $\frac{d^2\sigma}{d\Omega_\omega}$ (in nb/sr) for backward $\gamma^* + p \rightarrow p + \omega$ for fixed $W = 3.20$ GeV as a function of Q^2 in the u -channel nucleon exchange model for ωN TDAs. COZ [3] (long-dashed blue line) and KS [5] (solid green line) solutions for the leading twist nucleon DA are used as the phenomenological input.

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