

Complex analysis of Askaryan radiation: A fully analytic model in the time domain

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The detection of ultra-high-energy (UHE, ≥ 10 PeV) neutrinos via detectors designed to utilize the Askaryan effect has been a long-time goal of the astroparticle physics community. The Askaryan effect describes radio-frequency radiation from high-energy cascades. When a UHE neutrino initiates a cascade, cascade properties are imprinted on the radiation. Thus, observed radiation properties must be used to reconstruct the UHE neutrino event. Analytic Askaryan models have three advantages when used for UHE neutrino reconstruction. First, cascade properties may be derived from the match between analytic function and observed data. Second, analytic models minimize computational intensity in simulation packages. Third, analytic models can be embedded in firmware to enhance the real-time sensitivity of detectors. We present a fully analytic Askaryan model in the time domain for UHE neutrino-induced cascades in dense media that builds upon prior models in the genre. We then show that our model matches semianalytic parametrizations used in Monte Carlo simulations for the design of IceCube-Gen2. We find correlation coefficients greater than 0.95 and fractional power differences $< 5\%$ between the fully analytic and semianalytic approaches.

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I. INTRODUCTION

The extrasolar flux of neutrinos with energies between [0.01–1] PeV has been measured by the IceCube collaboration [1]. Previous analyses have shown that the discovery of ultra-high-energy neutrinos (UHE- ν) will require an expansion in detector volume because the flux is expected to decrease with energy [2–6]. The UHE- ν flux could potentially explain the origin of UHE cosmic rays (UHECR), and provides the opportunity to study electroweak interactions at record-breaking energies [7,8]. Utilizing the Askaryan effect expands the effective volume of UHE- ν detector designs, because this effect offers a way to detect UHE- ν with radio pulses that travel more than 1 km in sufficiently RF-transparent media such as Antarctic and Greenlandic ice [9–11].

The Askaryan effect occurs within a dense medium with an index of refraction n . A relativistic particle with $v > c/n$ initiates a high-energy cascade with negative total charge. The charge radiates energy in the RF bandwidth, and the radiation may be detected if the medium does not significantly attenuate the signal [12,13]. The IceCube EHE analysis has constrained the UHE- ν flux to be $E_\nu^2 \phi_\nu \leq 2 \times 10^{-8} \text{ GeV cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$ between $[5 \times 10^{15} - 2 \times 10^{19}]$ eV [4]. Arrays of $\mathcal{O}(100)$ *in situ* detectors encompassing effective areas of $\approx 10^4 \text{ m}^2 \text{ sr}$ per station, spaced by $\mathcal{O}(1)$

RF attenuation length could discover a UHE- ν flux beyond the EHE limits. The most suitable ice formations exist in Antarctica and Greenland, and a group of prototype Askaryan-class detectors has been deployed. These detectors seek to probe unexplored UHE- ν flux parameter space from astrophysical and cosmogenic sources [5,6,14,15].

Askaryan radiation was first measured in the laboratory in silica sand, and later ice [16–18]. Cascade properties affect the amplitude and phase of the radiation. At RF wavelengths, cascade particles radiate coherently, and the radiation amplitude scales with the total track length of the excess negative charge. The RF pulse shape is influenced by the *longitudinal length* of the cascade, and the pulse is strongest when the viewing angle is close to the Cherenkov angle, θ_C . The *excess charge profile* describes the excess negative charge versus longitudinal position on the cascade axis. Radiation wavelengths shorter than the *lateral width* of the cascade, perpendicular to the cascade axis, are attenuated. At energies far above 10 PeV in ice, however, excess charge profiles generated by *electromagnetic cascades* experience the Landau-Pomeranchuk-Migdal (LPM) effect and can have multiple peaks [19,20]. This theoretical foundation has been constructed from a variety of experimental and simulation results.

The field of Askaryan-class detectors requires this foundation for at least two reasons. First, the theoretical form of the Askaryan RF pulse is used to optimize RF detector designs. Askaryan models are incorporated into

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simulations [21–23] in order to calculate expected signals and aid in detector design. For example, reconstruction tools for the radio component of IceCube-Gen2 combine machine learning and insights from Askaryan radiation physics [24–26]. Second, Askaryan models are used as templates to search large data sets for signal candidates [5,27]. The signal-to-noise ratios (SNRs) at RF channels are expected to be small ($\text{SNR} \approx 3$), because the amplitude of the radiated field decreases with the vertex distance ($1/r$), and the signal is attenuated by the ice [9,28,29]. Low SNR signals require correspondingly low RF trigger thresholds, but signals must be sampled for a bandwidth of [0.1–1] GHz. Thus, RF channels are triggered at high rates by thermal noise. UHE- ν signals will be hidden within millions of thermal triggers. Template-waveform matching between models and data is a powerful technique for isolating RF signals from high-energy particles [27,30].

Askaryan models fall into three categories: full Monte Carlo (MC), semianalytic, and fully analytic. The original work by Zas *et al.* (ZHS) [13] was a full MC model. The properties of cascades with total energy ≤ 1 PeV were examined. A parametrization for the Askaryan field below 1 GHz was offered, attenuating modes above 1 GHz via a frequency-dependent form factor tied to the lateral cascade width. The semianalytic approach was introduced by Alvarez-Muñiz *et al.* (ARVZ) [31]. This approach accounts for fluctuations in the charge excess profile, and provides an analytic vector potential observed at the Cherenkov angle. The vector potential at the Cherenkov angle is labeled the form factor, and observed fields are derived from the derivative of the vector potential once convolved with a charge excess profile from MC. Recent work also accounts for differences in fit parameters from electromagnetic and hadronic cascades, and other interaction channels, while matching full MC simulations [32].

Finally, fully analytic models of Askaryan radiation from first principles have been introduced. Ralston and Buniy (RB) gave a fully analytic model valid for observations of cascades in the near and far field, with the transition encapsulated by a parameter η [33]. The result was a complex frequency-domain model. Recently, a model and software implementation was given by Hanson and Connolly (JCH + AC) that built upon RB by providing an analytic form factor derived from GEANT4 simulations, and accounted for LPM elongation [34]. This work connected the locations of poles in the complex frequency plane to η and the form factor. The poles combine to form a low-pass filter for the Askaryan radiation. The JCH + AC results match the ZHS results while demonstrating the physical origins of model parameters. The RB and JCH + AC results are given in the Fourier domain, but most UHE- ν searches (like template matching) have taken place in the time domain. The goals of this work are to produce a fully analytic time-domain model accounting for complex poles, valid for all viewing

angles θ and $\eta < 1$, and to demonstrate that it matches semianalytic models.

In Sec. II, the cascade geometry, units, and vocabulary are defined. In Sec. III, we describe how the JCH + AC form factor fits into the current model [34]. In Sec. IV, the analytic Askaryan field, observed at $\theta = \theta_C$ (*on-cone*), is presented. In Sec. V, the analytic Askaryan field observed for $\theta \neq \theta_C$ (*off-cone*) is presented. In Sec. VI, fully analytic fields are matched to semianalytic fields generated with NuRadioMC [23] at 10 PeV (electromagnetic cascades) and 100 PeV (hadronic cascades). Though the LPM effect is activated in NuRadioMC, it has a negligible influence on the waveform comparison at these energies. In Sec. VII, the results are summarized and potential applications of the model are described.

II. UNITS, DEFINITIONS, AND CONVENTIONS

The coordinate system of the Askaryan radiation from a vector current density \vec{J} is shown in Figs. 1(a)–1(b). Primed cylindrical coordinates refer to \vec{J} , and the unprimed spherical coordinates refer to the observer. The zenith or *viewing angle* is measured with respect to the *longitudinal axis* (z'). The observer displacement is $r = |\vec{x} - \vec{x}'|$, in the \hat{r} direction. The origin is located where the cascade has the highest instantaneous charge density (ICD). The ICD is treated with cylindrical symmetry, so it has no ϕ' dependence. This assumption is based on the large number of cascade particles and momentum conservation. The lateral extent of the ICD is along the *lateral axis* (ρ'). The viewing angle is θ in spherical coordinates, and the Cherenkov angle occurs when θ satisfies $\cos(\theta_C) = 1/n_{\text{ice}}$ with $n_{\text{ice}} = 1.78 \pm 0.003$ [35].

In Fig. 1(c), an example excess charge profile $n(z')$ is shown with characteristic longitudinal length a . The individual ICDs represent the excess charge density for small windows of time, and $n(z')$ refers to the total excess charge as a function of z' . Approximating the central

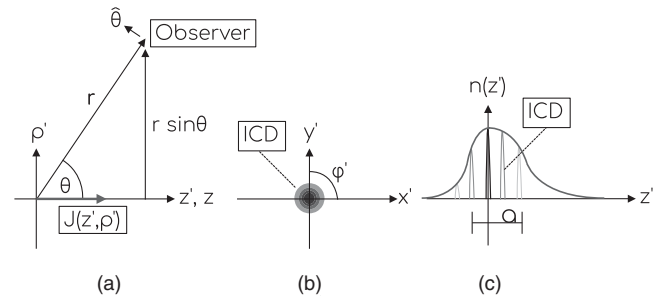


FIG. 1. (a) Side view of the coordinate systems used in the analysis. Spherical unprimed coordinates refer to the observer. Primed cylindrical coordinates refer to $\vec{J}(\rho', z')$. (b) Front view of the coordinate system. The ICD is assumed to have no ϕ' dependence. (c) The function $n(z')$ describes the total cascade excess charge, and it has a characteristic width a . The ICD has an instantaneous width much smaller than a [34].

portion of $n(z')$ as a Gaussian distribution $N(\mu, \sigma)$ corresponds to setting $a = 2\sigma$. Askaryan radiation occurs because $n(z')$ represents excess negative charge [13,34,36]. Cascades may be characterized as *electromagnetic*, initiated by charged outgoing leptons from UHE- ν interactions, or *hadronic*, initiated by the interaction between the UHE- ν and the nucleus. Electromagnetic cascades follow the Greisen distribution and hadronic cascades follow the Gaisser-Hillas distribution. An example of such an implementation via the ARVZ semianalytic parametrization is AraSim [11].

The units of the electromagnetic field in the Fourier domain are V/m/Hz, often converted in the literature to V/m/MHz. To make the distance dependence explicit, both sides of field equations are multiplied by r , as in $r\vec{E} = \dots$, making the units V/Hz. Throughout this work, an overall field normalization constant E_0 is used. E_0 may be linearly scaled with energy, as in other Askaryan models. We show that the on-cone field amplitude is proportional to E_0 times a characteristic frequency squared, so the units of E_0 are V/Hz². For off-cone results, we show that the field amplitude is proportional to E_0 times a characteristic frequency divided by a characteristic pulse width, and the units of E_0 remain V/Hz².

In Sec. III B, we review briefly the energy dependence of the longitudinal length a in both the electromagnetic and hadronic cases. For the Greisen distribution with critical energy E_{crit} , it can be shown that if $n_{\text{max}} = n(z_{\text{max}})$, where $z_{\text{max}} = \ln(E_C/E_{\text{crit}})$, then $n_{\text{max}}a \sim E_C/E_{\text{crit}}$. Thus, the area under the curve $n(z')$ scales with the total cascade energy E_C . RB demonstrated that the Askaryan radiation amplitude is proportional to $n_{\text{max}}a$ and therefore E_C . The cascade develops over a length $\approx a$, but the radiation is coherent over a length $\Delta z'_{\text{coh}}$ for which the displacement is constant to first order relative to a wavelength. The η parameter is the square of the ratio of a to $\Delta z'_{\text{coh}}$:

$$\eta = \left(\frac{a}{\Delta z'_{\text{coh}}} \right)^2 = \frac{k}{r} (a \sin \theta)^2. \quad (1)$$

In the far field, $\eta < 1$. In the first JCH + AC model, a limiting frequency ω_C was shown to filter the Askaryan radiation [34]:

$$\eta = \frac{\omega}{\omega_C}. \quad (2)$$

The effect of ω_C is described in Sec. IV. The Askaryan radiation is primarily polarized in the $\hat{\theta}$ direction, with a small amount along \hat{r} [31,34]. The wave vector is $k = (2\pi)/(n\lambda)$, where n is the index of refraction. A 3D wave vector was defined by RB, equivalent to $\vec{q} = nk(1, \vec{\rho}/R)$. The vector current density is treated by RB as a charge density times the velocity of the ICD: $\vec{J}(t, \vec{x}') = \rho(z' - vt, \rho')\vec{v}$. Further, the charge density is factored into $n(z')$ times the ICD:

$\rho(z' - vt, \rho') = n(z')f(z' - vt, \rho')$. The form factor \tilde{F} is the three-dimensional spatial Fourier transform of the ICD [33].

The result for \tilde{F} was derived analytically by JCH + AC [34], and that derivation is briefly described in Sec. III A. JCH + AC define a parameter σ , and \tilde{F} is a function of σ : $\tilde{F}(\sigma)$. The variable σ is related to the ratio of the lateral ICD width to the radiated wavelength. In the derivation of \tilde{F} , it is convenient to set σ equal to the ratio of the angular frequency to the low-pass cutoff frequency ω_{CF} of \tilde{F} :

$$\sigma = \frac{\omega}{\omega_{\text{CF}}}. \quad (3)$$

Armed with \tilde{F} , the longitudinal length a and the corresponding energy dependence on E_0 , the RB field equations $\vec{\mathcal{E}}$, and the displacement r , the Askaryan electromagnetic field may be assembled according to the following form [33]:

$$r\vec{E}(\omega, \theta) = E_0 \left(\frac{\omega}{2\pi} \right) \psi \vec{\mathcal{E}}(\omega, \theta) \tilde{F}(\omega, \theta). \quad (4)$$

The factor E_0 is proportional to the cascade energy. The factor ω is the angular frequency. The variable $\psi = -i \exp(ikr) \sin \theta$. The function $\vec{\mathcal{E}}(\omega, \theta)$ contains the vector and complex pole structure of the field (see Refs. [33,34]). The model represented by Eq. (4) is an *all- θ , all- ω* model. That is, Eq. (4) is valid at all frequencies and all viewing angles, provided one accepts the approximation of the central portion of $n(z')$ as Gaussian. The first goal of this work is to build an *all- θ , all- t* model in the time domain, derived from Eq. (4), and the second goal is to compare it to semianalytic parametrizations.

III. THE FORM FACTOR AND LONGITUDINAL LENGTH PARAMETER

To arrive at the main electromagnetic field in the time domain, the individual pieces of Eq. (4) must first be assembled. The first piece will be the form factor \tilde{F} that accounts for the 3D ICD, followed by some remarks about the energy dependence of the longitudinal length parameter a .

A. The form factor

The form factor is the 3D Fourier transform of the ICD $f(z', \vec{\rho}')$, with $\vec{q} = nk(1, \vec{\rho}/R)$ [33]:

$$F(\vec{q}) = \int d^3x' f(z', \rho') e^{-i\vec{q}\cdot\vec{x}'}. \quad (5)$$

The goal is to evaluate \tilde{F} in the Fourier domain for an ICD definition informed by cascade simulations. Simulations of the cascade induced by UHE- ν indicate a

thin wave of charge in z' spread uniformly in ϕ' , that decreases exponentially in ρ' . Using these observations JCH + AC completed the derivation in Ref. [34]. The final result was a simple analytic formula:

$$\tilde{F} = \frac{1}{(1 + (\omega/\omega_{\text{CF}})^2)^{3/2}} \quad (6)$$

The form factor acts as a low-pass filter with the cutoff frequency ω_{CF} :

$$\tilde{F} \approx \frac{\omega_0^2}{(\omega + i\omega_0)(\omega - i\omega_0)}. \quad (7)$$

The definition $\omega_0 = \sqrt{2/3}\omega_{\text{CF}}$ has been used. The effect of the approximation is illustrated in Fig. 2. Equation (7) matches the original ZHS parametrization (see Eq. 20 of Ref. [13]).

I. A note about the Molière radius

In Sec. VI B, the decay constant l of the lateral component of the ICD is inferred from best-fit values of ω_0 . The connection between the l parameter and ω_0 was described by JCH + AC [34]. Put simply, the ICD decays by a factor of $1/e$ a lateral distance l from the cascade axis. Note, however, that the l parameter is not the Molière radius. The Molière radius is the lateral radius which forms a cylinder containing 90% of the energy deposition of the cascade. For ice with a density of 0.917 g cm^{-3} , one can estimate $R_M \approx 9.2 \text{ cm}$ using standard formulas. Although it is tempting to compare l to R_M , these parameters have different definitions. Knowing that l is related to ω_0 , l may be estimated as $\lambda/2$ in ice at the cutoff frequency. At 3 GHz in ice, $\lambda/2 \approx 2.8 \text{ cm}$, and at 1 GHz in ice, $\lambda/2 \approx 8.4 \text{ cm}$. Although the results are at the same order of magnitude as R_M , there are three effects limiting the high-frequency spectrum of the radiation: ω_0 , ω_C , and the viewing angle. Thus, $l < R_M$ is possible for a radiation spectrum limited to $\lesssim 1 \text{ GHz}$.

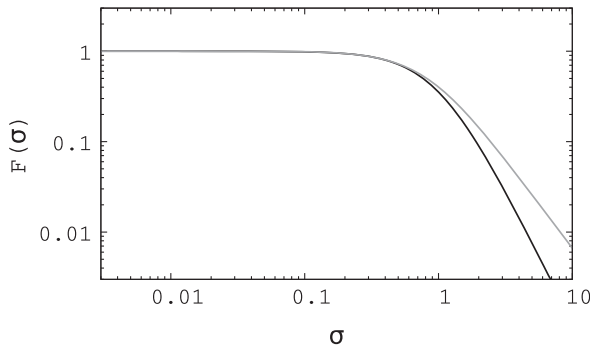


FIG. 2. Black: Eq. (6), graphed versus $\sigma = \omega/\omega_{\text{CF}}$. Gray: the two-pole approximation.

B. The longitudinal length parameter

The next piece required in the assembly of the main electromagnetic field is the energy dependence of the overall amplitude, and the energy dependence of the longitudinal length parameter, a , which is a part of $\vec{\mathcal{E}}$ in Eq. (4) [33]. What follows are two separate discussions, one for electromagnetic cascades, and one for hadronic cascades. Though we share these calculations for convenience, note that a variety of theoretical and experimental results on this topic are available [16,37,38].

1. Electromagnetic case

The number of charged particles versus distance in radiation lengths $n(z')$ in an electromagnetic cascade taking place in a dense medium with initial cascade energy E_C , critical energy E_{crit} , normalization parameter n_0 , and age s is [34]

$$n(z') = \frac{n_0}{\sqrt{\ln(E_C/E_{\text{crit}})}} \exp \left\{ z' \left(1 - \frac{3}{2} \ln(s) \right) \right\}. \quad (8)$$

To find the energy-dependent width of the Greisen distribution, four steps are necessary: (1) normalization of $n(z')$ as a fraction of the maximum excess charge, (2) conversion of $n(z')$ to $n(s)$, (3) determination of the width of $n(s)$ by approximating the central portion as a Gaussian distribution, and (4) conversion of the width from s units to radiation lengths z' , and then converting those results to a distance. Define the ratio $R = n(z_{\text{max}} \pm a/2)/n_{\text{max}}$, so the FWHM occurs when $R = 0.5$. The final result in radiation lengths is

$$a = \sqrt{\ln(E_C/E_{\text{crit}})} \sqrt{-6 \ln(R)} \quad (9)$$

Since $R < 1$, $\ln(R) < 0$ and a is real valued, and a in Eq. (9) is in radiation lengths. In solid ice the density is $\rho_{\text{ice}} = 0.917 \text{ g cm}^{-3}$, and the electromagnetic radiation length is $z_0 = 36.08 \text{ g cm}^{-2}$ [34]. Converting to distance gives

$$a = \frac{z_0}{\rho_{\text{ice}}} \sqrt{\ln(E_C/E_{\text{crit}})} \sqrt{-6 \ln(R)} \quad (10)$$

Note that $a \propto \sqrt{\ln(E_C)}$, as shown by RB and others. The product $n_{\text{max}} a$ is proportional to the energy E_C/E_{crit} . For this reason RB took $n_{\text{max}} a$ as the field normalization rather than E_C [33]. As an example, letting $R = 0.4$, and $E_{\text{crit}} \approx 10^8 \text{ eV}$, gives $a \approx 4$ meters for $E_C = 10^{16} \text{ eV}$. We show in Sec. VI that our fitted a values are close to 4 meters when matched to semianalytic parametrizations.

2. Hadronic case

The Gaisser-Hillas distribution describes hadronic cosmic-ray air showers, but has also been applied to hadronic cascades in dense media in codes like AraSim [11,22]. The original function reads

$$n(z') = n_{\max} \left(\frac{z' - z_0}{z_{\max} - z_0} \right)^{(z_{\max} - z_0)/\lambda} e^{\frac{z_{\max} - z'}{\lambda}}. \quad (11)$$

The variables are defined as follows: n_{\max} is the instantaneous maximum number of particles in the cascade, z' is the longitudinal distance in radiation lengths, z_0 is the initial starting point, λ is the interaction length, and z'_{\max} is the location of n_{\max} . Using the same steps as the electromagnetic case, we find

$$a = \sqrt{\lambda z'_{\max}} \sqrt{-8 \ln(R)} \quad (12)$$

The a parameter again goes as $\sqrt{z_{\max}} \propto \sqrt{\ln(E_C)}$ which produces similar lengths as the electromagnetic case when scaled by the appropriate interaction length and ice density.

IV. ON-CONE FIELD EQUATIONS

The $\hat{\theta}$ component of the electromagnetic field at $\theta = \theta_C$ will now be built in the time domain from Eq. (4). Setting $\theta = \theta_C$ in the general RB field equations (Appendix A), with Eq. (6) for \tilde{F} , $\sigma = \omega/\omega_{CF}$ and $\eta = \omega/\omega_{CF}$, and letting E_0 be proportional to the cascade energy E_C produces Eq. 45 from JCH + AC [34]:

$$r\tilde{E}(\omega, \theta_C) = \frac{(-i\omega)E_0 \sin(\theta_C) e^{i\omega r/c}}{(1 - i\omega/\omega_C)^{1/2} (1 + (\omega/\omega_{CF})^2)^{3/2}}. \quad (13)$$

More detail is provided in Appendix A. Let the retarded time be $t_r = t - r/c$ (with the appropriate value for c given the index versus depth), and let $\omega_0 = \sqrt{\frac{2}{3}}\omega_{CF}$ and $\hat{E}_0 = E_0 \sin \theta_C$. Finally, let $\epsilon = \omega_0/\omega_C$. The inverse Fourier transform of Eq. (13) is

$$rE(t, \theta_C) = \frac{\hat{E}_0 i \omega_C \omega_0^2}{\pi} \frac{d}{dt_r} \int_{-\infty}^{\infty} \frac{e^{-i\omega t_r}}{(2i\omega_C + \omega)(\omega + i\omega_0)(\omega - i\omega_0)} d\omega. \quad (14)$$

In Eq. (14), the derivative with respect to the retarded time d/dt_r is introduced to remove a factor of $(-i\omega)$ from the numerator. Accounting for the complex poles and the sign of t_r , complex integration and expansion to first order in ϵ yields

$$rE(t, \theta_C) = \frac{1}{3} \hat{E}_0 \omega_{CF}^2 \begin{cases} (1 - \frac{1}{2}\epsilon) e^{\omega_0 t_r} & t_r < 0 \\ (2e^{-2\omega_C t_r} - (1 + \frac{1}{2}\epsilon) e^{-\omega_0 t_r}) & t_r > 0 \end{cases} \quad (15)$$

Equation (15) represents the time-domain solution for the on-cone $\hat{\theta}$ component of the Askaryan electric field. The expansion to first order in ϵ is only performed so the final result resembles semianalytic results for $\vec{E} = -\partial \vec{A} / \partial t_r$ [31,32]. Table I summarizes the definitions of the parameters in Eq. (15). Fit results for the parameters of Table I are shown in Sec. VI.

Notice that the amplitude is asymmetric, and the parameter ϵ influences the asymmetry. The ϵ parameter was studied by JCH + AC in detail. For example, Fig. 10 of Ref. [34] shows that $\epsilon \approx [0.1 - 1]$ for an inverse lateral width $l^{-1} = \sqrt{2\pi}\rho_0 \approx 20 \text{ m}^{-1}$ and $a \approx 4 \text{ m}$. The best-fit results for ϵ and a are shown in Sec. VI. JCH + AC showed that the expression for ϵ is the product of the ratio of the lateral to longitudinal length, and the ratio of the longitudinal length to the observer displacement, making it a physical parameter connecting the event geometry to the cascade shape [34]. Figure 3 displays normalized examples of Eq. (15) for different values of ω_0 , ω_C , and ϵ .

A. Verification of the uncertainty principle

As a check on the procedures used to perform the inverse Fourier transform that produces Eq. (15), we verify below that the uncertainty principle holds, for $\Delta\theta \rightarrow 0$. JCH + AC provided the Gaussian width of the radiation in the Fourier domain: σ_ν , where ν represents the frequency in Hz. Generally speaking, Fourier transform pairs must obey $\sigma_\nu \sigma_t \geq 1/(2\pi)$. The following procedure is used to compute the width σ_t of the on-cone field. First, the $t_r < 0$ and $t_r > 0$ cases are each treated as probability distributions and normalized. Next, the average positive and negative retarded times, $\bar{t}_{r,+}$ and $\bar{t}_{r,-}$, are computed. Finally, subtracting the two averages yields σ_t :

TABLE I. Parameters used to build Eq. (15). Fitted values in comparison to semianalytic parametrizations are shown in Sec. VI.

Parameter	Definition
\hat{E}_0	$E_0 \sin(\theta_C)$
E_0	$\approx n_{\max} a$
ω_0	$\sqrt{\frac{2}{3}}\omega_{CF}$
ω_{CF}	$(c\sqrt{2\pi}\rho_0)/(n \sin \theta)$ (see Eqs. 22, 23, and 46 of Ref. [34])
ω_C	$(rc)/(na^2 \sin^2 \theta)$ (see Eq. 39 of Ref. [34])
ϵ	ω_0/ω_C
t_r	$t - r/c$

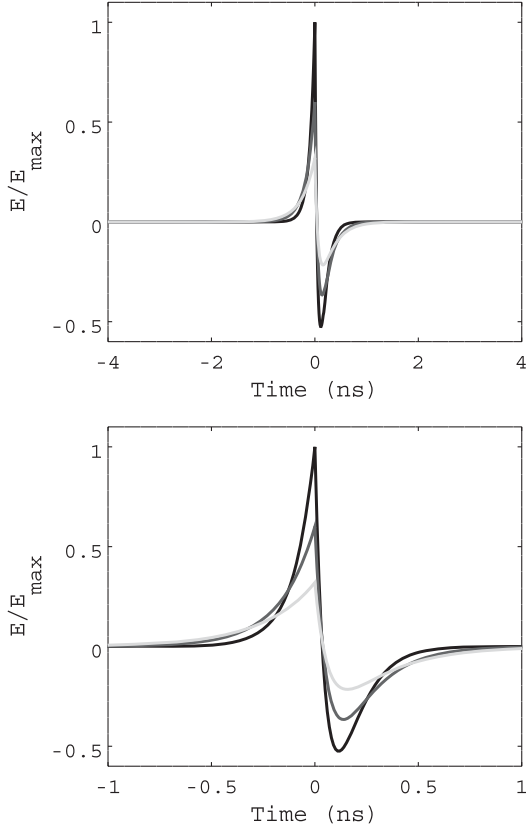


FIG. 3. Top: Eq. (15) from $[-4, 4]$ ns, with (black) $\omega_C = 2\pi(1.25)$ GHz, $\omega_0 = 2\pi(1.56)$ GHz, $\epsilon = 1.25$, (gray) $\omega_C = 2\pi(1.25)$ GHz, $\omega_0 = 2\pi(0.94)$ GHz, $\epsilon = 0.75$, (light gray) $\omega_C = 2\pi(1.25)$ GHz, $\omega_0 = 2\pi(0.625)$ GHz, and $\epsilon = 0.5$. The amplitudes of all curves are normalized to the peak of the $\epsilon = 1.25$ (black) data. Bottom: same as top panel, plotted between $[-1, 1]$ ns.

$$\sigma_t = \bar{t}_{r,+} - \bar{t}_{r,-} = \frac{\epsilon + 2}{\omega_0} = \frac{1}{\omega_C} + \frac{2}{\omega_0}. \quad (16)$$

The result has the correct units and the limiting cases are sensible. If $\epsilon \rightarrow 0$ ($\omega_C \gg \omega_0$), then $\sigma_t \rightarrow 2/\omega_0$, which is expected from observing Eq. (15) if the ω_C exponential disappears. If $\epsilon = 1$ ($\omega_C = \omega_0$), then $\sigma_t = 3/\omega_0$. That is, the pulse is wider if there is more than one relevant cutoff frequency.

The expression for σ_ν is given by Eq. 36 of JCH + AC [34]:

$$\sigma_\nu = \frac{c}{2\pi a \Delta \cos \theta} (1 + \eta^2)^{1/2}. \quad (17)$$

Expanding to first order in $\Delta \cos(\theta) = \cos(\theta) - \cos(\theta_C)$,

$$\sigma_\nu \approx \frac{c}{2\pi a \sin(\theta_C) \Delta \theta} (1 + \eta^2)^{1/2}. \quad (18)$$

From Table I: $\omega_C^{-1} = na^2 \sin^2(\theta_C)/(rc)$, and $\omega_0^{-1} = nl \sin(\theta_C)/c$, with $l = \sqrt{3/2}/(\sqrt{2\pi\rho_0})$. (Recall that ρ_0

is a parameter discussed in Ref. [34]). Multiplying σ_t and σ_ν with the far-field limit ($\eta < 1$) gives the inequality

$$\sigma_\nu \sigma_t \geq \frac{n}{2\pi} \left(\left(\frac{a}{r} \right) \frac{\sin(\theta_C)}{\Delta \theta} + 2 \left(\frac{l}{a} \right) \frac{1}{\Delta \theta} \right). \quad (19)$$

Therefore, in order to satisfy $\sigma_\nu \sigma_t > 1/(2\pi)$,

$$n \left(\frac{a}{r} \right) \sin(\theta_C) + 2n \left(\frac{l}{a} \right) > \Delta \theta. \quad (20)$$

Although $a/r \ll 1$ and $l/a \ll 1$, as long as these expressions do not approach zero as fast as $\Delta \theta \rightarrow 0$ in Eq. (20), the uncertainty principle holds. Yet these are exactly the conditions of the problem: a displacement r in the far field (but not infinitely far away) and a longitudinal length a much larger (but not infinitely larger) than the lateral ICD width l . Thus, $\sigma_\nu \sigma_t > 1/(2\pi)$ holds.

V. OFF-CONE FIELD EQUATIONS

Turning to the case for which $\theta \neq \theta_C$, the $\hat{\theta}$ component of the electromagnetic field will now be built in the time domain. The RB field equations for the $\hat{\theta}$ and \hat{r} components were summarized by both RB and JCH + AC [33,34], and are included here in Appendix A. Recall the general form of the electromagnetic field, given in Eq. (4):

$$r\vec{E}(\omega, \theta) = E_0 \left(\frac{\omega}{2\pi} \right) \psi \vec{\mathcal{E}}(\omega, \theta) \vec{F}(\omega, \theta). \quad (21)$$

The first task is to simplify $\vec{\mathcal{E}}(\omega, \theta)$ before taking the inverse Fourier transform. The simplification involves expanding $\vec{\mathcal{E}}(\omega, \theta)$ in a Taylor series such that $u = 1 - i\eta \approx 1$, restricting $\eta < 1$ (far-field). Once $\vec{\mathcal{E}}(\omega, \theta)$ is simplified, the inverse Fourier transform of Eq. (21) may be evaluated to produce the result. Table II contains useful variable definitions, Table III contains useful function definitions, and Table IV contains special cases of the functions in Table III.

The original form of $\vec{\mathcal{E}}(\eta, \theta)$ is shown in Appendix A. Changing variables to u and x (Table II) and using the

TABLE II. Useful variables for the derivation of the off-cone Askaryan electromagnetic field.

Variable	Definition
u	$1 - i\eta$
x	$\cos(\theta)$
x_C	$\cos(\theta_C)$
q	$(xx_C - x_C^2)/(1 - x^2)$
y	$(\frac{1}{2})(ka)^2(\cos \theta - \cos \theta_C)^2$
p	$\frac{1}{2}(\frac{a}{r})^2(\cos \theta - \cos \theta_C)^2$

TABLE III. Useful functions for the derivation of the off-cone Askaryan electromagnetic field. The last row contains the vector structure of the $\hat{\theta}$ component of the field.

Function	Definition
$f(u, x)$	$(u + 3 \frac{(1-u)^2 x^2 - x x_c}{1-x^2})^{-1/2}$
$g(u, x)$	$\exp(-\frac{1}{2}(ka)^2(x - x_c)^2 u^{-1})$
$h(u, x)$	$(\frac{1-u}{u})q$
$\vec{\mathcal{E}}(u, x) \cdot \hat{\theta}$	$f(u, x)g(u, x)(1 - h(u, x))$

TABLE IV. Special cases of the functions defined in Table III, when $u = 1$.

Function ($u = 1$)	Result
$f(x, 1)$	1
$\dot{f} _{u=1}$	$-\frac{1}{2}$
$g(x, 1)$	$\exp(-y)$
$\dot{g} _{u=1}$	$y \exp(-y)$
$h(x, 1)$	0
$\dot{h} _{u=1}$	$-q$

function definitions and values in Tables III–IV, $\vec{\mathcal{E}}(u, x) \cdot \hat{\theta} = \mathcal{E}(u, x)$ becomes

$$\mathcal{E}(u, x) = f(u, x)g(u, x)(1 - h(u, x)). \quad (22)$$

Expanding $\mathcal{E}(u, x)$ near $u = 1$ gives

$$\mathcal{E}(u, x) = \mathcal{E}(x, 1) + (u - 1)\dot{\mathcal{E}}(x, 1) + \mathcal{O}(u - 1)^2. \quad (23)$$

The details of the expansion are shown in Appendix B. The result is.

$$\mathcal{E}(x, u) = e^{-y} \left(1 - \frac{1}{2} j\eta(2y + 2q - 1) \right). \quad (24)$$

The inverse Fourier transform of the $\hat{\theta}$ component gives the time-domain results, after including the expanded $\mathcal{E}(u, x)$:

$$rE(t, \theta) = \mathcal{F}^{-1} \left\{ E_0 \left(\frac{\omega}{2\pi} \right) \tilde{F}\psi\mathcal{E} \right\}. \quad (25)$$

Intriguingly, the result is proportional to the *line-broadening function*, H (DLMF 7.19, [39]) common to spectroscopy applications. There are three terms in Eq. (24). Two terms ultimately vanish, being integrals over odd integrands (see Appendix B). The integral that remains contains H , with $\omega_1 = t_r/(2p)$:

$$I_0 = 2\pi i \left(\frac{\omega_C}{\omega_0} \right) e^{-\frac{t_r^2}{4p}} H(\sqrt{p}\omega_0, i\omega_1\sqrt{p}). \quad (26)$$

The line-broadening function is similar to a convolution between a Gaussian function and a Lorentzian function, and cannot be expressed analytically, though there are examples of polynomial expansions [40]. Note that, for situations relevant to the current problem, $\omega > \omega_1$. Requiring that $\omega > \omega_1$ amounts to a restriction between $\Delta\theta$ and $|t_r|$:

$$|t_r| < |2p\omega|. \quad (27)$$

It is shown in the next section that $\sqrt{2p}$ is the pulse width σ_t , so $|2p\omega|$ has units of time. Using the results of Sec. VA below, the restriction on the retarded time may be written as $|t_r|/\sigma_t < \omega\sigma_t = 2\pi(\sigma_t/T)$. That is, the accuracy of the waveform should be trusted within a number of pulse widths that is less than 2π times the ratio of the pulse width to the period of the lowest frequency. This is not a strong requirement, since the field quickly approaches zero after several pulse widths. Hereafter, this step will be called the *symmetric approximation*, because the result for $r\vec{E}(t_r, \theta)$ in Eq. (28) has equal positive and negative amplitudes. Evaluating the line-broadening function numerically would account for amplitude asymmetry. The restriction on $\Delta\theta$ is formalized in Sec. VB. Solving I_0 using the symmetric approximation clears the way for the final result (see Appendix B):

$$rE(t, \theta) = -\frac{E_0\omega_0 \sin(\theta)}{8\pi p} t_r e^{-\frac{t_r^2}{4p} + p\omega_0^2} \operatorname{erfc}(\sqrt{p}\omega_0) \quad (28)$$

Equation (28) represents the time-domain solution for the off-cone $\hat{\theta}$ component of the Askaryan electric field. Equation (28) is graphed in Figs. 4 and 5. In Fig. 4 (top), $E(t, \theta)$ is shown normalized to the maximum value for the angular range displayed, $[\theta_C + 1.5^\circ, \theta_C + 5.5^\circ]$, from $t = [-5, 5]$ ns. Pulses with viewing angles closer to θ_C have larger relative amplitudes and shorter pulse widths. Figure 4 (bottom) contains the same results, but for $t = [-1.5, 1.5]$ ns. The pulses are symmetric and all zero crossings are at $t_r = 0$ ns as a result of the symmetric approximation. Figure 5 contains contours of the same results as in Fig. 4.

As in the on-cone result, the overall field amplitude scales with energy ($E_0 \sim n_{\max} a$). However, the amplitude also scales with ω_0/p . The argument of the complementary error function, $\sqrt{p}\omega_0$, is unitless. This factor is strictly positive, so the range of the complementary error function is (0,1). The factor $\sqrt{p}\omega_0$ cannot be zero without setting $\theta = \theta_C$, or setting $\omega_{CF} = 0$. Both cases are not allowed. Equation (28) represents the *off-cone* ($\theta \neq \theta_C$) solution, so $p \neq 0$. Setting $\omega_{CF} = 0$ is not physical, for

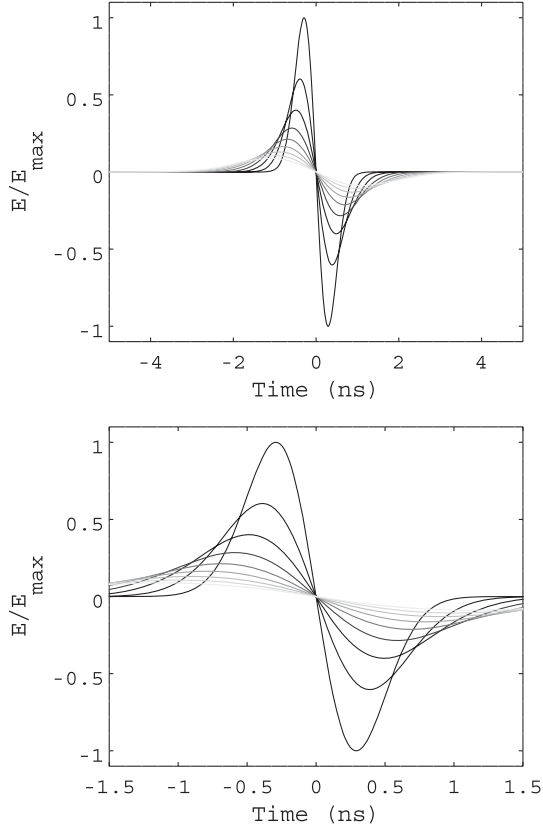


FIG. 4. $E(t, \theta)$ versus t_r [Eq. (28)], normalized. The viewing angle θ is varied from $\theta_C + 1.5^\circ$ to $\theta_C + 5.5^\circ$ in steps of 0.5° . Top: $\omega_0/(2\pi) = 1.0$ GHz. Bottom: same as top, but zoomed in on the central region.

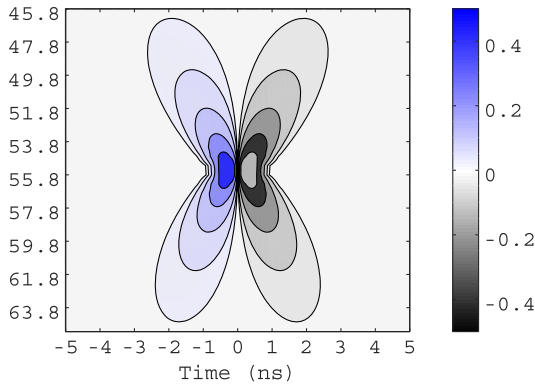


FIG. 5. Contours of $E(t, \theta)$ versus θ versus t_r [Eq. (28)], normalized. The normalization is the same as in Fig. 4. Although the contour lines extend into the region near θ_C , Eq. (5) is only being evaluated at $\Delta\theta > 1.5^\circ$ (see text for details).

this implies infinite lateral width (l) and cascade particles have finite transverse momentum. Another possibility is that $p = 0$ if $a = 0$, but this implies $E_0 = 0$. Therefore, $0 < \text{erfc}(\sqrt{p}\omega_0) < 1$.

A. Verification of the uncertainty principle

As in Sec. IV A, the uncertainty principle should be checked. Equation (28) is an antisymmetric Gaussian function with pulse width $\sigma_t = \sqrt{2p}$. Let $\Delta \cos \theta = (\cos \theta - \cos \theta_C)$. Using Table II, the expression $\sqrt{2p}$ evaluates to

$$\sigma_t = \sqrt{2p} = \left(\frac{a}{c}\right) (\Delta \cos \theta). \quad (29)$$

Recall that σ_ν is given by

$$\sigma_\nu = \frac{c}{2\pi a \Delta \cos \theta} (1 + \eta^2)^{1/2}. \quad (30)$$

The uncertainty product is

$$\sigma_t \sigma_\nu = \frac{1}{2\pi} (1 + \eta^2)^{1/2}. \quad (31)$$

In the far field, $\eta < 1$, so $\sigma_t \sigma_\nu \geq 1/(2\pi)$ holds.

B. Usage of the on-cone versus off-cone fields

The form of Eq. (28), and the restriction between $\Delta\theta$ and $|t_r|$ from the symmetric approximation suggests the limit $\Delta\theta \rightarrow 0$ must be examined carefully. Since $p \propto (\cos \theta - \cos \theta_C)^2$, probing the model near $\theta = \theta_C$ is equivalent to taking the limit that $p \rightarrow 0$. Intriguingly, the p^{-1} dependence in the field does not lead to a divergence. As the field grows in amplitude from p^{-1} as $p \rightarrow 0$, the field width, $\sqrt{2p}$, approaches zero.

Equations (16) and (29) contain the pulse widths of the on-cone and off-cone fields, respectively. Power in the off-cone case is limited by the pulse width $\sqrt{2p}$, and the observed power increases as $\Delta\theta$ and $\sqrt{2p}$ both decrease. Thus, a reasonable constraint on when $\Delta\theta_{\min}$ is large enough to use Eq. (28) is given by setting the off-cone pulse width equal to the on-cone pulse width:

$$\frac{1}{\omega_C} + \frac{2}{\omega_0} = \sqrt{2p}. \quad (32)$$

Expanding the expression for p near $\theta = \theta_C$, and evaluating the square root leads to

$$\frac{1}{\omega_C} + \frac{2}{\omega_0} = \frac{a}{c} \sin \theta_C \Delta\theta_{\min}. \quad (33)$$

Using $\epsilon = \omega_0/\omega_C$, and letting $k_0 = \omega_0/c$, the formula may be rearranged:

$$\epsilon + 2 = a k_0 \sin \theta_0 \Delta\theta_{\min}. \quad (34)$$

Squaring both sides, and then dividing both sides by r yields

$$\frac{(\epsilon + 2)^2}{r} = k_0 \left(\frac{k_0 (a \sin \theta_C)^2}{r} \right) \Delta \theta_{\min}^2. \quad (35)$$

The quantity in parentheses on the right-hand side is η , with $\omega = \omega_0$. Setting $\omega = \omega_0$ means $\eta = \epsilon$. Solving for $\Delta \theta_{\min}$ gives

$$\Delta \theta_{\min} = \frac{\epsilon + 2}{\sqrt{\epsilon k_0 r}}. \quad (36)$$

Assuming $\epsilon \approx 1$, $f_0 \approx 1$ GHz, $n = 1.78$ for solid ice, and $c = 0.3$ m ns⁻¹ (see Sec. VI A), $k_0 \approx 35$ m⁻¹. Taking $r = 1000$ m, $\Delta \theta_{\min} \approx 1^\circ$. Simple rules of thumb for the application of Eq. (28) field are

$$\Delta \theta_{\min} \geq 1^\circ, \quad (37)$$

$$\Delta \theta_{\min} \propto \frac{1}{\sqrt{kr}}. \quad (38)$$

VI. COMPARISON TO SEMIANALYTIC PARAMETRIZATIONS

The fully analytic model will now be compared to the ARVZ semianalytic parametrization used in NuRadioMC to predict signals in IceCube-Gen2 Radio [23]. Specifically, the comparison is between Eqs. (15) and (28) and the NuRadioMC implementation of the semianalytic parametrization given in Ref. [32]. To provide concrete comparisons, a small set of waveforms was generated with NuRadioMC, for both electromagnetic and hadronic cascades, on- and off-cone. The electromagnetic cascades have $E_C = 10^{16}$ eV, while the hadronic cascades have $E_C = 10^{17}$ eV. These choices minimize the impact of the LPM effect, though the LPM effect was activated in the NuRadioMC code.

The comparison involves three stages. First, waveforms and a values are generated for each cascade type, energy, and angle: $\theta = \theta_C + 3.0^\circ$, and $\theta = \theta_C$. Second, Eqs. (15) and (28) are tuned to match the waveforms. In each fit, the Pearson correlation coefficient (ρ) is maximized, and the sum squared of amplitude differences $[(\Delta E)^2]$ is minimized. Finally, best-fit parameters are tabulated.

Two remarks are important regarding the fit criteria. First, the Pearson correlation coefficient is not sensitive to changes in amplitude because it is normalized:

$$\rho = \frac{\text{COV}(f_{\text{data}}, f_{\text{model}})}{\sigma_{\text{data}} \sigma_{\text{model}}}. \quad (39)$$

If E_i represent the samples of the models, then

$$(\Delta E)^2 = \sum_{i=1}^N (E_{i,\text{data}} - E_{i,\text{model}})^2. \quad (40)$$

A. Waveform comparison: $\theta = \theta_C$

1. Electromagnetic case

Six different electromagnetic cascades and the corresponding Askaryan fields were generated using the ARZ2019 model from NuRadioMC [23,32] for comparison to Eq. (15). The cascades have $E_C = 10$ PeV, and $r = 1000$ meters. The LPM effect is activated in NuRadioMC for all comparisons in this work. The units of $\vec{E}(t, \theta_C)$ are mV/m versus nanoseconds, so the units of $r\vec{E}$ are volts. The sampling rate of the digitized semianalytic parametrizations was 100 GHz, with $N = 2048$ samples. Let $f_C = \omega_C/(2\pi)$ and $f_0 = \omega_0/(2\pi)$. The frequencies f_C and f_0 were varied from [0.6–6.0] GHz. The parameter E_0 was varied from [0.05 – 5.0] V GHz⁻². In a simple two-level FOR loop, the Pearson correlation coefficient ρ was maximized by varying f_0 and f_C . Next, the sum of the squared amplitude differences $(\Delta E)^2$ was minimized by varying E_0 , while holding f_0 and f_C fixed. Several other schemes were studied, including a three-level FOR loop, but the two-stage process produced the best results. The results are shown in Fig. 6.

Maximizing ρ corresponds to minimizing $(\Delta E)^2$. In Fig. 7, $(\Delta E)^2$ is graphed versus ρ for one event. Best-fit ρ values are ≈ 0.97 for this set, corresponding to best-fit $(\Delta E)^2$ values of $\approx 7\%$. Contours of $\rho > 0.95$ for f_0 versus f_C are shown in Fig. 6 (left column). The crosses represent the best-fit locations. The dashed gray line at $y = x$ corresponds to $f_0/f_C = \epsilon = 1$. Though Eq. (15) contains an expansion to first order in ϵ , making it resemble the derivative of the vector potential from the ARVZ semianalytic parametrization [32], the expansion is optional. There is a restriction that $\epsilon \neq 2$ [see Eq. (A14) of Appendix A]. Thus, the best-fit ϵ values avoid the solid black lines ($\epsilon = 2$) in Fig. 6, but are large enough to account for pulse asymmetry. The best-fit waveforms are shown in Fig. 6 (right column). The gray curves correspond to the semianalytic parametrization, and the black curves represent Eq. (15).

Table V contains the best-fit results for the Eq. (15) parameters, along with best-fit ρ values and $(\Delta E)^2$ values. The horizontal and vertical distances from the crosses to the $\rho > 0.95$ contour are used as error estimates for f_0 and f_C in Table V. The a errors typically encompass the a values from NuRadioMC. The full region in $[f_0, f_C]$ space for which UHE- ν signals are expected the radio component of the IceCube Gen2 detector will be the topic of future studies, along with the apparent difference in ϵ values depending on the electromagnetic or hadronic classification of the cascade (see Fig. 8).

2. Hadronic case

Using the same procedure as in the electromagnetic case, NuRadioMC was used to generate six hadronic cascades at 100 PeV for comparison to Eq. (15). The energy was

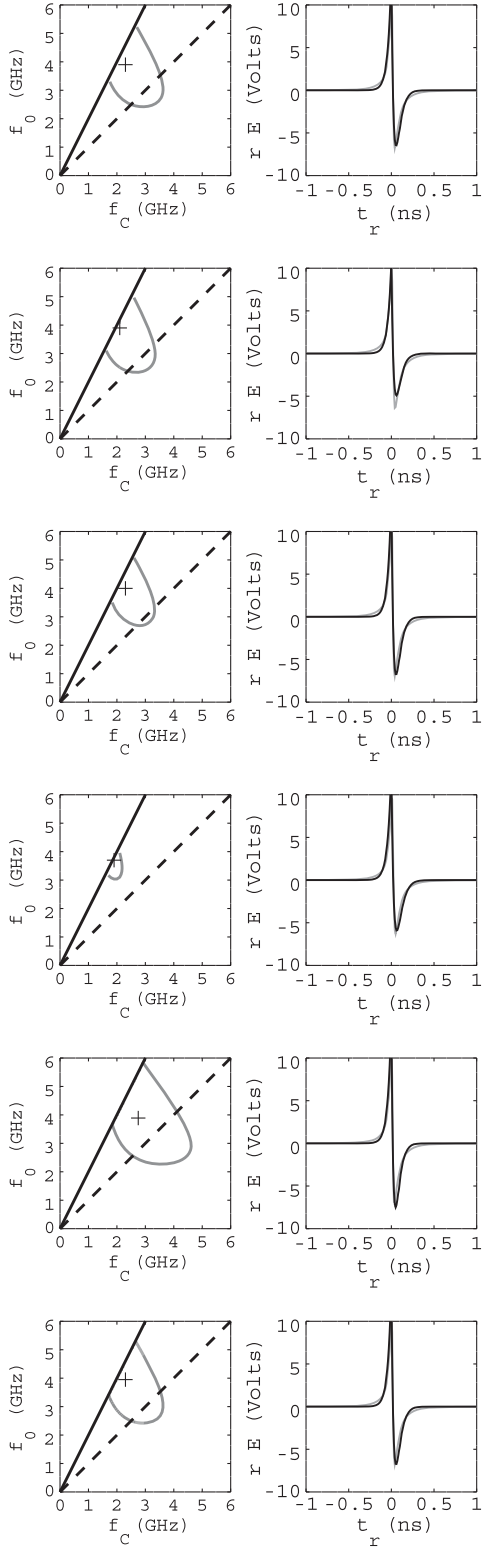


FIG. 6. Fit results: electromagnetic case, $\theta = \theta_C$, $E_C = 10$ PeV. The rows correspond to NuRadioMC waveforms 1–6, with 10 PeV electromagnetic cascades. Left column: best fits for f_0 and f_C . Dashed line: $\epsilon = 1$. Solid line: $\epsilon = 2$. Gray contour: $\rho > 0.95$. Black cross: best fit. Right column: best-fit waveforms. Gray: semianalytic parametrizations from Ref. [23]. Black: Eq. (15).

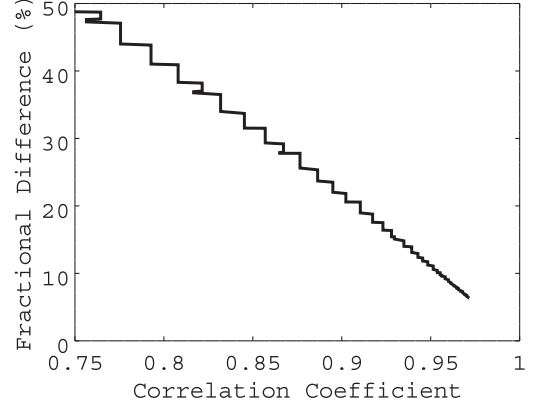


FIG. 7. Fractional difference in the sum of amplitude differences squared $[(\Delta E)^2]$ versus correlation coefficient (ρ) for waveform 1 at $E_C = 10$ PeV, for the electromagnetic case.

increased to show that the model describes a range of energies, so the waveform amplitudes are larger by a factor of 10 relative to the 10 PeV case. The LPM effect is activated in NuRadioMC for all comparisons in this work. The main results are shown in Fig. 8, and the correlation contours represent $\rho = 0.985$.

The results shown in Fig. 8 demonstrate that modeling hadronic cascades at $\theta = \theta_C$ is similar to the electromagnetic case, with one interesting difference. The contours enclose best-fit ϵ values below the dashed line, whereas the fits to the electromagnetic cases were above the dashed line. This could indicate a potential discriminator for cascade classification. Another difference between the electromagnetic and hadronic cases is that the gray contours in Fig. 8 correspond to $\rho = 0.985$, as opposed to $\rho = 0.95$ in the electromagnetic case.

TABLE V. Fit results: electromagnetic case, $\theta = \theta_C$, $E_C = 10$ PeV. The six rows (from top to bottom) correspond to NuRadioMC waveforms 1–6, WITH 10 PeV electromagnetic cascades. From left to right: the form-factor cutoff frequency, coherence cutoff frequency, energy-scaling normalization, longitudinal length parameter, the best-fit correlation coefficient, and the relative power difference between the NuRadioMC semianalytic parametrization and the fully analytic model. The parameter means and errors in the mean are quoted in the bottom two rows.

No.	f_0 (GHz)	f_C (GHz)	E_0 (V GHz ⁻²)	a_{wave} (m), a_{MC} (m)	ρ	$(\Delta E)^2$ (%)
1	$3.9^{+0.2}_{-1.9}$	$2.3^{+1.3}_{-0.3}$	0.3	$4.1^{+1.2}_{-0.3}$, 4.85	0.97	6.5
2	$3.9^{+0.3}_{-1.5}$	$2.1^{0.9}_{-0.1}$	0.5	$4.3^{+1.8}_{-0.2}$, 6.35	0.97	10.9
3	$4.0^{+1.2}_{-1.0}$	$2.3^{+0.8}_{-0.4}$	0.35	$4.1^{+0.7}_{-0.4}$, 4.48	0.96	7.5
4	$3.7^{+0.1}_{-0.5}$	$1.9^{+0.5}_{-0.1}$	1.85	$4.5^{+1.1}_{-0.3}$, 5.6	0.955	8.9
5	$3.9^{+1.4}_{-0.9}$	$2.7^{+1.4}_{-0.8}$	0.18	$4.0^{+2.0}_{-1.2}$, 4.48	0.97	5.7
6	$3.9^{+1.3}_{-1.9}$	$2.3^{+1.3}_{-0.3}$	0.31	$4.1^{+2.0}_{-0.5}$, 4.85	0.97	6.4
Average	3.88	2.3	0.6	4.18	0.966	7.7
Error	3.08	0.1	0.3	0.07	0.003	0.8

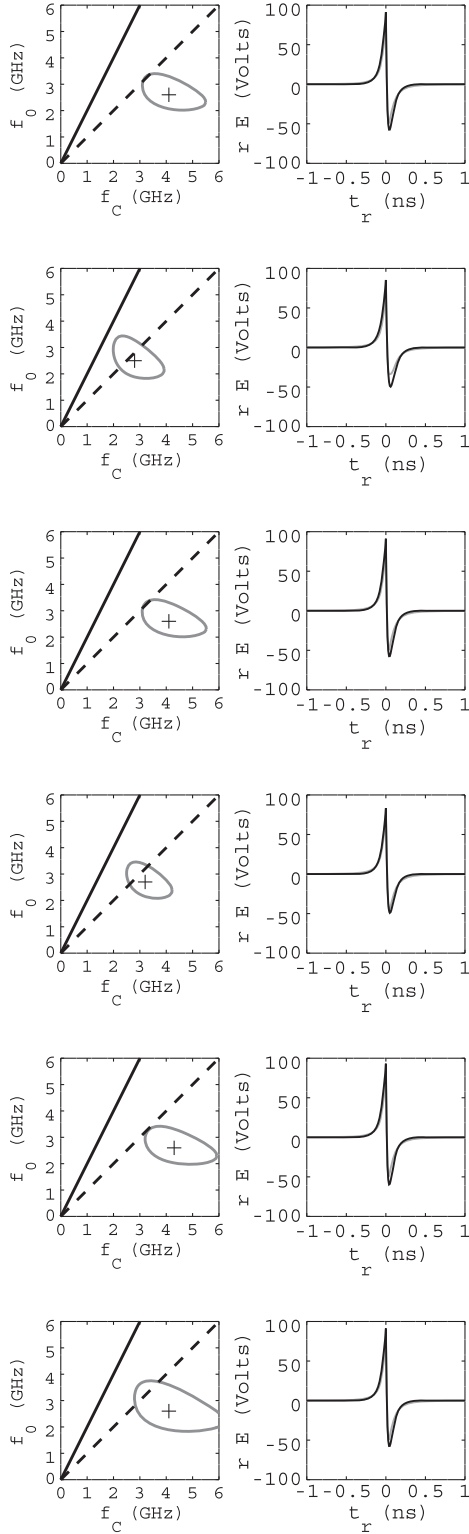


FIG. 8. Fit results: hadronic case, $\theta = \theta_C$, $E_C = 100$ PeV. The six rows (from top to bottom) correspond to NuRadioMC waveforms 1–6, with 100 PeV hadronic cascades. Left column: best fits for f_0 and f_C . Dashed line: $\epsilon = 1$. Solid line: $\epsilon = 2$. Gray contour: $\rho > 0.985$. Black cross: best fit. Right column: best-fit waveforms. Gray: semianalytic parametrizations from Ref. [23]. Black: Eq. (15).

Table VI contains the best-fit parameters corresponding to Fig. 8. The typical power difference $(\Delta E)^2$ has decreased with respect to the electromagnetic case. The ρ values all exceed 0.985, and the $(\Delta E)^2$ results are typically below 2 percent. Intriguingly, $\epsilon < 1$ means higher f_C values, which in turn yields systematically low a values relative to those generated in NuRadioMC, despite the increased energy. Reconstructed a values are still within a factor of 2 of the MC-true values. Despite the systematic offset, the best-fit a and the NuRadioMC a values are tightly correlated (see Fig. 11 below).

B. Waveform comparison: $\theta \neq \theta_C$

1. Electromagnetic case

The general comparison procedure of Sec. VIA was repeated with the same semianalytic parametrization from NuRadioMC, but with 12 new events each viewed at $\theta = \theta_C + 3.0^\circ$ (six electromagnetic cascades, six hadronic). One difference is that ω_0 only changes the waveform amplitude, along with E_0 . The pulse width $\sigma_t = \sqrt{2\rho}$ connects the longitudinal length a and the viewing angle with respect to the Cherenkov angle.

The fit procedure was performed in two stages. First, θ values and a values were scanned from $[\theta_C + 1.5^\circ, \theta_C + 10.0^\circ]$ and $[0.1, 10]$ meters, respectively, to maximize ρ . Once the best-fit values for a and θ were determined, $(\Delta E)^2$ was minimized by varying $f_0 = \omega_0/(2\pi)$ and E_0 from $[0.3, 3.0]$ GHz and $[0.1, 2.0]$ V GHz $^{-2}$, respectively. The (θ, a) scan and the (f_0, E_0) scan were each separate two-level for loops. The results are shown in Fig. 9.

TABLE VI. Fit results: hadronic case, $\theta = \theta_C$, $E_C = 100$ PeV. The six rows (from top to bottom) correspond to NuRadioMC waveforms 1–6, with 100 PeV hadronic cascades. From left to right: the form-factor cutoff frequency, coherence cutoff frequency, energy-scaling normalization, longitudinal length parameter, the best-fit correlation coefficient, and the relative power difference between the NuRadioMC semianalytic parametrization and the fully analytic model. The parameter means and errors in the mean are quoted in the bottom two rows.

No.	f_0 (GHz)	f_C (GHz)	E_0 (V GHz $^{-2}$)	a_{wave} (m), a_{MC} (m)	ρ	$(\Delta E)^2$ (%)
1	$2.6^{+0.6}_{-0.6}$	$4.1^{+1.1}_{-1.0}$	1.0	$3.1^{+0.8}_{-0.8}$, 5.23	0.99	1.86
2	$2.5^{+0.7}_{-0.6}$	$2.8^{+0.9}_{-0.8}$	1.25	$3.75^{+1.2}_{-1.1}$, 6.35	0.99	1.83
3	$2.6^{+0.7}_{-0.6}$	$4.1^{+1.2}_{-0.9}$	1.0	$3.1^{+0.9}_{-0.7}$, 5.23	0.99	1.83
4	$2.7^{+0.6}_{-0.5}$	$3.2^{+0.8}_{-0.6}$	1.0	$3.5^{+0.9}_{-0.7}$, 6.35	0.99	2.5
5	$2.6^{+0.7}_{-0.6}$	$4.3^{+1.4}_{-1.1}$	1.0	$3.0^{+1.0}_{-0.75}$, 4.85	0.99	1.755
6	$2.6^{+1.4}_{-0.7}$	$4.1^{+1.9}_{-1.2}$	1.0	$3.1^{+1.4}_{-0.9}$, 5.23	0.99	1.86
Average	2.60	3.75	1.04	3.3	0.99	1.9
Error	0.03	0.25	0.04	0.1	0.0	0.1

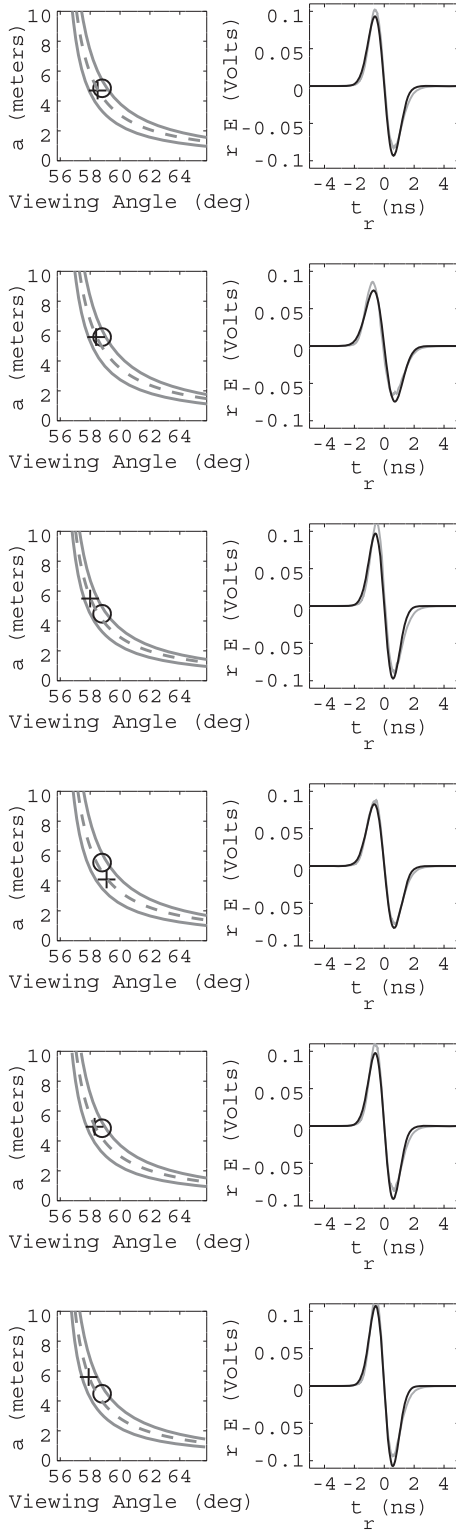


FIG. 9. Fit results: electromagnetic case, $\theta \neq \theta_C$, $E_C = 10$ PeV. The six rows (from top to bottom) correspond to NuRadioMC waveforms 1–6, with 10 PeV electromagnetic cascades. Left column: best-fit θ and a values. Crosses: best fits. Circles: MC true values. Gray contour: $\rho > 0.95$. Dashed line: a versus θ from Eq. (29). Right column: best-fit waveforms. Gray: semianalytic parametrizations from Ref. [23]. Black: Eq. (28).

In Fig. 9 (left column), the best-fit a values and θ values are marked with a cross. The circles represent the MC true values. Circles and crosses lie on the dashed lines, because an uncertainty principle connects a values to θ values (see Sec. VA). Specifically, Eq. (29) may be used to show, to first order in $\Delta\theta = \theta - \theta_C$,

$$a\Delta\theta = \frac{c\sqrt{2p}}{\sin\theta_C} = \text{const.} \quad (41)$$

The pulse width $\sigma_t = \sqrt{2p}$ is a constant derived from the waveform, implying that the product of a and $\Delta\theta$ is constant. The parameters a and $\Delta\theta$ are therefore inversely proportional: $a \propto \Delta\theta^{-1}$. The shape of the $\rho > 0.95$ contour follows this inverse proportionality. The dashed lines represent Eq. (41). These results suggest that a measurement of the Askaryan pulse width would constrain the cascade shape and geometry. The best-fit waveforms are shown in Fig. 9 (right column). Typical correlation coefficients exceed $\rho = 0.98$. Table VII contains the fit results. The fit results include estimates of the lateral width parameter, l , derived from f_0 (see Sec. III A). Despite making the symmetric approximation to arrive at Eq. (28), the fits include fractional power differences of $\approx 3\%$.

TABLE VII. Fit results: electromagnetic case, $\theta \neq \theta_C$, $E_C = 10$ PeV. The six rows (from top to bottom) correspond to NuRadioMC waveforms 1–6, with 10 PeV electromagnetic cascades. From left to right: the viewing angle, longitudinal length parameter, the form-factor cutoff frequency, the energy-scaling normalization, the lateral width of the cascade, the best-fit correlation coefficient, and the relative power difference between the NuRadioMC semianalytic parametrization and the fully analytic model. The parameter means and errors in the mean are quoted in the bottom two rows.

No.	θ_{wave} (deg), θ_{MC} (deg)	a_{wave} (m), a_{MC} (m)	f_0 (GHz)	E_0 (V GHz ⁻²)	l (cm)	ρ	$(\Delta E)^2$ (%)
1	$58.5^{+0.7}_{-0.6}$, 58.8	$4.7^{+1.3}_{-1.0}$, 4.85	0.75	1.2	$3.4^{+0.9}_{-0.7}$	0.99	1.93
2	$58.4^{+0.6}_{-0.5}$, 58.8	$5.6^{+1.4}_{-1.1}$, 5.60	1.0	1.2	$2.6^{+0.4}_{-0.3}$	0.99	2.61
3	$58.0^{+0.5}_{-0.4}$, 58.8	$5.5^{+1.3}_{-1.0}$, 4.48	1.0	1.1	$2.6^{+0.3}_{-0.2}$	0.98	4.47
4	$59.1^{+0.9}_{-0.7}$, 58.8	$4.1^{+1.2}_{-0.9}$, 5.23	0.75	1.2	$3.4^{+0.5}_{-0.5}$	0.995	0.80
5	$58.3^{+0.7}_{-0.5}$, 58.8	$4.95^{+1.4}_{-1.1}$, 4.85	0.75	1.2	$3.4^{+0.4}_{-0.3}$	0.99	1.8
6	$57.9^{+0.6}_{-0.4}$, 58.8	$5.6^{+1.5}_{-1.2}$, 4.48	0.75	1.2	$3.5^{+0.5}_{-0.4}$	0.99	1.83
Average	58.4	5.1	0.83	1.18	3.2	0.989	2.2
Error	0.2	0.2	0.05	0.02	0.2	0.002	0.5

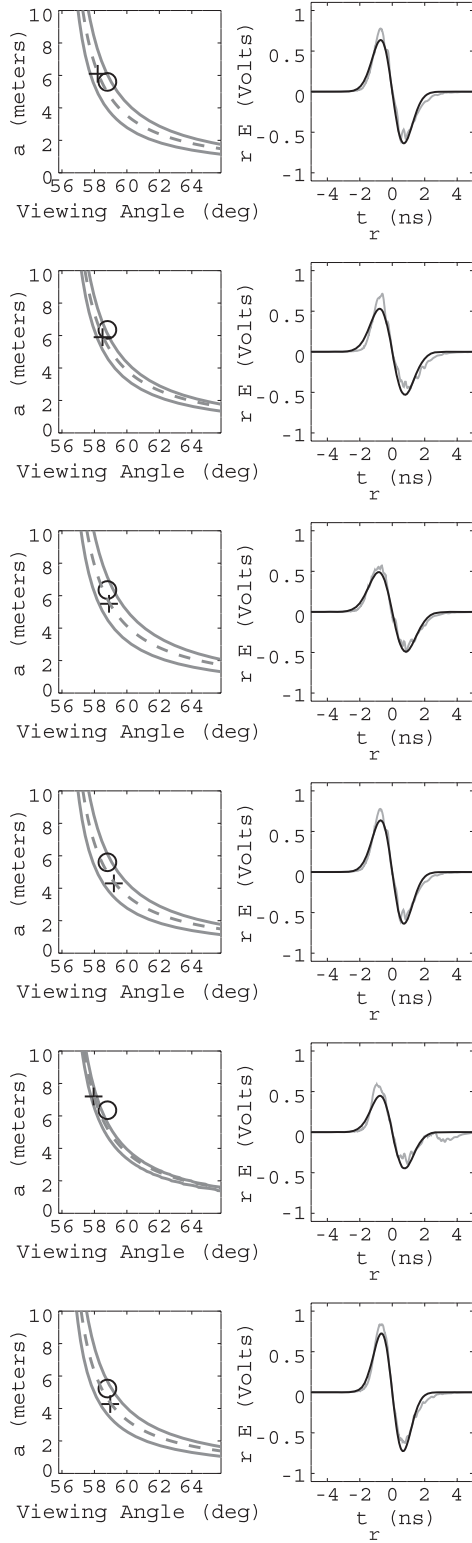


FIG. 10. Fit results: hadronic case, $\theta \neq \theta_C$, $E_C = 100$ PeV. The six rows (from top to bottom) correspond to NuRadioMC waveforms 1–6, with 100 PeV hadronic cascades. Left column: best-fit θ and a values. Crosses: best fits. Circles: MC true values. Gray contour: $\rho > 0.95$. Dashed line: a versus θ from Eq. (29) (uncertainty principle). Right column: best-fit waveforms. Gray: semianalytic parametrizations from Ref. [23]. Black: Eq. (28).

2. Hadronic case

The fit procedure for the hadronic cascades was the same as the electromagnetic case, except that the range for E_0 was expanded to $[1.0, 20.0]$ V GHz $^{-2}$. As in the on-cone procedure, the hadronic cascade energy was $E_C = 100$ PeV. The results are shown in Fig. 10.

As with the electromagnetic case, ρ is maximized and $(\Delta E)^2$ is minimized. Table VIII contains the best-fit parameters, along with ρ and $(\Delta E)^2$. Solutions with $\rho \approx 0.98$ and $(\Delta E)^2 \approx 5\%$ were found. Similar to the results shown in Table VII, the results in Table VIII are in agreement with the MC values from NuRadioMC. The E_0 values match expectations for 100 PeV cascades, because they are a factor of 10 higher than those of the 10 PeV electromagnetic case. The results for a , l , and f_0 , however, are not statistically different between Tables VII and VIII. Future studies will require computing the probability distributions of these parameters from large numbers of UHE- ν cascades.

As a first exercise for statistical energy reconstruction from waveform parameters, assume that $\theta = \theta_C + 3.0^\circ$ is already measured. For example, θ could be determined by measuring the cutoff frequency in the Fourier domain below 1 GHz (see Fig. 5 of Ref. [34], for example). Scanning Eq. (28) over all NuRadioMC waveforms at fixed $\theta = \theta_C + 3.0^\circ$ yields Fig. 11, in which the fitted a value from each waveform

TABLE VIII. Fit results: hadronic case, $\theta \neq \theta_C$, $E_C = 100$ PeV. The six rows (from top to bottom) correspond to NuRadioMC waveforms 1–6, with 10 PeV hadronic cascades. From left to right: the viewing angle, longitudinal length parameter, form-factor cutoff frequency, the energy-scaling normalization, the lateral width of the cascade, the best-fit correlation coefficient, and the relative power difference between the NuRadioMC semianalytic parametrization and the fully analytic model. The parameter means and errors in the mean are quoted in the bottom two rows.

No.	θ_{wave} (deg), θ_{MC} (deg)	a_{wave} (m), a_{MC} (m)	f_0 (GHz)	E_0 (V GHz $^{-2}$)	l (cm)	ρ	$(\Delta E)^2$ (%)
1	$58.2^{+0.6}_{-0.4}$, 58.8	$6.1^{+1.5}_{-1.2}$, 5.6	0.8	10.6	$3.2^{+0.5}_{-0.5}$	0.98	3.55
2	$58.5^{+0.4}_{-0.3}$, 58.8	$5.9^{+0.9}_{-0.8}$, 6.35	0.85	10.3	$3.0^{+0.3}_{-0.2}$	0.96	7.1
3	$58.9^{+0.8}_{-0.6}$, 58.8	$5.5^{+1.4}_{-1.1}$, 6.35	0.9	10.8	$2.8^{+0.5}_{-0.5}$	0.98	2.64
4	$59.2^{+0.8}_{-0.7}$, 58.8	$4.3^{+1.1}_{-0.8}$, 5.6	0.85	10.5	$3.0^{+0.5}_{-0.5}$	0.98	3.10
5	$58.0^{+0.2}_{-0.2}$, 58.8	$7.2^{+0.6}_{-0.6}$, 6.35	0.9	8.2	$2.9^{+0.3}_{-0.3}$	0.955	8.76
6	$59.0^{+0.8}_{-0.6}$, 58.8	$4.3^{+1.1}_{-0.9}$, 5.23	0.85	10.4	$3.0^{+0.5}_{-0.5}$	0.985	3.00
Average	58.6	5.5	0.86	10.1	3.2	0.973	5
Error	0.2	0.5	0.015	0.4	0.2	0.005	1

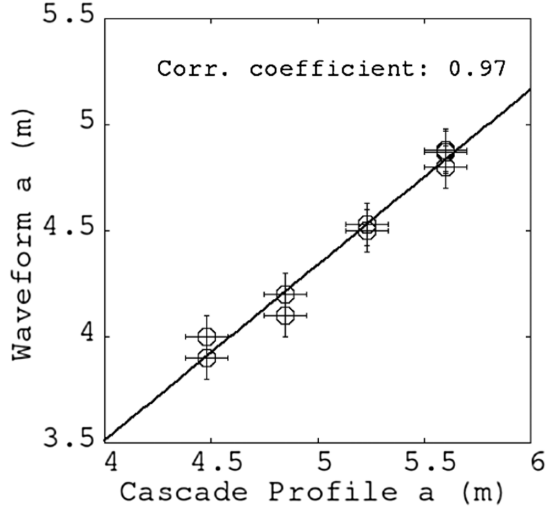


FIG. 11. The longitudinal length parameter a derived from the Eq. (28) best fit versus the a value derived from the cascade profile in NuRadioMC. A linear fit and correlation coefficient are shown (slope: 0.83 ± 0.05 ; intercept: 0.2 ± 0.2 (m); correlation coefficient = 0.97).

is graphed versus the MC-true a value. The a errors in all cases are taken to be ± 10 cm (\pm two Δa step sizes). A least-squares linear fit was applied to the data. The linear function fits the data, and the correlation coefficient is 0.97. The results in Fig. 11 imply an energy reconstruction technique using the formulas found in Sec. III B. Consider the relationship between a and $\ln(E_C/E_{\text{crit}})$: $a = c_1 \sqrt{\ln(E_C/E_{\text{crit}})}$. The fractional error in $\ln(E_C/E_{\text{crit}})$ is proportional to the fractional error in a :

$$\frac{\sigma_{\ln(E_C/E_{\text{crit}})}}{\ln(E_C/E_{\text{crit}})} = 2c_1 \left(\frac{\sigma_a}{a} \right). \quad (42)$$

If a reliable fit for the a parameter is obtained from observed Askaryan waveforms, Eq. (42) shows that the logarithm of the energy can be constrained.

VII. CONCLUSION

We have presented a fully analytic Askaryan model in the time domain, and we have shown that it matches results generated with semianalytic parametrizations used in NuRadioMC. Pearson correlation coefficients between the fully analytic and semianalytic parametrizations were found to be greater than 0.95, and typical fractional differences in total power were found to be $\approx 5\%$. New results and potential applications are summarized in the following sections.

A. Summary of new results

The main results are summarized in Table IX. This work represents the first time the two distinct pole frequencies f_0 and f_C have been used to characterize the time-domain

TABLE IX. Summary of results in this work.

Result	Location
$r\vec{E}(t_r, \theta_C)$, on-cone field ($\hat{\theta}$)	Eq. (15), Sec. IV
$\sigma_t \sigma_\nu \geq 1/(2\pi)$, on-cone	Eq. (20), Sec. IVA
$r\vec{E}(t_r, \theta)$, off-cone field ($\hat{\theta}$)	Eq. (28), Sec. V
$\sigma_t \sigma_\nu \geq 1/(2\pi)$, off-cone	Eq. (31), Sec. VA
On-cone EM comparison to [32]	Fig. 6, Table V
On-cone HAD comparison to [32]	Fig. 8, Table VI
Off-cone EM comparison to [32]	Fig. 9, Table VII
Off-cone HAD comparison to [32]	Fig. 10, Table VIII

field equations of the Askaryan effect for both $\theta = \theta_C$ and $\theta \neq \theta_C$. The uncertainty principle was verified on-cone ($\theta = \theta_C$), serving as a check on the model. By fitting on-cone cascade parameters, we have shown that an analytic model matches semianalytic predictions. The ϵ parameter reveals a potential cascade classification scheme. Next, the off-cone ($\theta \neq \theta_C$) field equations were derived, and again the uncertainty principle was verified. Off-cone cascade parameters were fit, and the results are in excellent agreement with semianalytic results. Fitting a values has revealed a potential energy reconstruction.

To obtain the fields on- and off-cone, $\eta < 1$ was assumed. The restriction $\eta < 1$ means that Eqs. (15) and (28) must be applied to the far field. Given that a and θ_C are fixed by cascade physics and ice density, and that the relevant Askaryan bandwidth for ice is [0.1–1] GHz, the parameter most easily varied within η is the observer distance r . Taking $\nu = 0.5$ GHz, $n = 1.78$, $c = 0.3$ m GHz, $\theta = \theta_C$, and $a = 5$ m, requiring that $\eta = 1$ gives $r \geq 0.4$ km. Scaling to $\nu = 0.25$ GHz gives $r \geq 0.2$ km. According to NuRadioMC [23] (Fig. 13), the r corresponding to UHE- ν at 10^{18} eV ranges from 0.7–3.2 km, and 0.2 km is rare.

The “acceleration argument” invoked by RB in Ref. [33] states that if $r(t)$ points to the ICD, $r(t)$ must be constant enough to ensure that $\Delta r < \lambda$. Using the law of cosines, with two sides being r and $r + \Delta r$, and a third being a , the criteria that $(a/r)^2 \ll 1$ leads to $|\Delta r| \approx a/n$ which is $O(2)$ m. When in doubt about usage and event geometry, determining if $(a/r)^2 \ll 1$ is a good check. If the UHE- ν event is a charged-current interaction with an electromagnetic cascade far above the LPM energy for ice, a grows faster than $\sqrt{\ln(E_C/E_{\text{crit}})}$ [20].

B. Utility of the analytic model

There are at least four advantages of fully analytic Askaryan models. First, when analytic models are matched to observed data, cascade properties may be derived directly from the waveforms. Second, in large-scale simulations, evaluating a fully analytic model technically provides a speed advantage over other approaches. Third, fully analytic models, combined with RF channel

response, can be embedded in firmware to form a *matched filter* that enhances UHE- ν detection probability. Fourth, parameters in analytic models may be *scaled* to produce results that apply to media of different density than ice. This application is useful for understanding potential signals in the Antarctic firm, or the upper layer of snow and ice that is of lower density than the solid ice beneath it.

The ability to fit cascade properties from waveforms will be a useful tool for the radio component of IceCube-Gen2. Examples of current reconstruction techniques include the forward-folding method [25] and information field theory [26]. In particular, the longitudinal length parameter a leads to a reconstruction of $\ln(E_C)$, given knowledge of $\Delta\theta$ [Fig. 11 and Eq. (42)]. Further, all designs for detector stations in IceCube-Gen2 radio include many distinct RF channels and one phased array of channels. Matching our analytic model to each channel waveform will provide a separate measurement of parameters like a and θ (see gray contours of Figs. 4 and 5). The ensuing global fit should constrain the event energy and geometry.

The most intriguing usage for a fully analytic Askaryan model would be to embed the model as a *matched filter* in detector firmware. Because cascade properties are unknown *a priori*, an array of matched filters could be implemented to form a *matched filter bank*. One example of this approach was the TARA experiment [41], which was designed to detect low-SNR cosmic ray radar echoes. This is similar to the challenge faced by the radio component of the IceCube Gen2 detector: pushing the limit of low-SNR RF pulse detection in a remote setting. For example, a matched filter bank could be formed with an array of off-cone field formulas with fixed a value and varying θ values, which would then be convolved with the RF channel impulse response (see Sec. 6 of Ref. [27]).

Finally, a fully analytic model enhances the ability of IceCube-Gen2 radio to identify signals that originate in the firm. At the South Pole, the RF index of refraction begins around 1.35 and does not reach the solid ice value of 1.78 until 150–200 meters [28]. There are at least two signals that could originate in the firm: UHE- ν events that create Askaryan radiation, and UHE cosmic-ray cascades partially inside or fully inside the firm. The altitude of the South Pole makes the latter possible. The Askaryan radiation of the firm UHE- ν events could be modeled via appropriate density scaling of the cascade parameters.

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APPENDIX A: DETAILS OF THE ON-CONE FIELD EQUATION DERIVATION

The original equations for the $\hat{\theta}$ component of $\vec{\mathcal{E}}$ are

$$\mathcal{W}(\eta, \theta) = \frac{\exp\left(-\frac{1}{2}(ka)^2 \frac{(\cos\theta - \cos\theta_C)^2}{1-i\eta}\right)}{\left(1 - i\eta\left(1 - 3i\eta \frac{\cos\theta \cos\theta_C - \cos\theta_C}{\sin^2\theta} \frac{1}{1-i\eta}\right)\right)^{1/2}}, \quad (\text{A1})$$

$$\vec{\mathcal{E}}(\eta, \theta) \cdot \hat{\theta} = \mathcal{W}(\eta, \theta) \left(1 - i\eta \frac{\cos\theta_C \cos\theta - \cos\theta_C}{\sin^2\theta} \frac{1}{1-i\eta}\right). \quad (\text{A2})$$

Letting $\theta = \theta_C$ yields

$$\vec{\mathcal{E}}(\eta, \theta) \cdot \hat{\theta} = \frac{1}{\sqrt{1-i\eta}}. \quad (\text{A3})$$

The complete field from the original RB model [33], including the form factor \tilde{F} , $\psi = -i \exp(ikr) \sin\theta$, and $\vec{\mathcal{E}}$ is

$$r\vec{E}(\omega, \theta) = E_0 \left(\frac{\omega}{2\pi}\right) \psi \vec{\mathcal{E}}(\eta, \theta) \tilde{F}. \quad (\text{A4})$$

Take Eq. 6 for the form factor \tilde{F} . Let E_0 be proportional to cascade energy E_C and let $\theta = \theta_C$:

$$r\vec{E}(\omega, \theta_C) = \frac{(-i\omega)E_0 \sin(\theta_C) e^{i\omega R/c}}{(1 - i\omega/\omega_C)^{1/2} (1 + (\omega/\omega_{CF})^2)^{3/2}}. \quad (\text{A5})$$

Suppose $\omega < \omega_C$, and $\omega < \omega_{CF}$, such that the following approximations of the factors in the denominator are valid:

$$(1 - i\omega/\omega_C)^{1/2} \approx 1 - \frac{i}{2} \frac{\omega}{\omega_C}, \quad (\text{A6})$$

$$(1 + (\omega/\omega_{CF})^2)^{3/2} \approx 1 + \frac{3}{2} \left(\frac{\omega}{\omega_{CF}}\right)^2. \quad (\text{A7})$$

Using the approximations introduces simple poles into the complex formula for the frequency-dependent electric field. Inserting the approximations in the denominator of Eq. (A5), we have

$$r\vec{E}(\omega, \theta_C) = \frac{(-i\omega)E_0 \sin(\theta_C) e^{i\omega R/c}}{\left(1 - \frac{i}{2}\omega/\omega_C\right) \left(1 + \frac{3}{2}\left(\omega/\omega_{CF}\right)^2\right)}. \quad (\text{A8})$$

The denominator can be rearranged by factoring the ω coefficients, and defining $\omega_0 = \sqrt{\frac{2}{3}}\omega_{CF}$:

$$r\tilde{E}(\omega, \theta_C) = \frac{2i\omega_C\omega_0^2(-i\omega)E_0 \sin(\theta_C)e^{i\omega r/c}}{(2i\omega_C + \omega)(\omega + i\omega_0)(\omega - i\omega_0)}. \quad (\text{A9})$$

Let $\hat{E}_0 = E_0 \sin(\theta_C)$, and let the retarded time be $t_r = t - r/c$. Taking the *inverse* Fourier transform, using the same sign convention as RB [33] [$f(t) = (2\pi)^{-1} \int_{-\infty}^{\infty} \tilde{F}(\omega)e^{-i\omega t} d\omega$], converts the field to the time domain:

$$\begin{aligned} rE(t, \theta_C) &= \frac{\hat{E}_0 i\omega_C \omega_0^2}{\pi} \frac{d}{dt_r} \int_{-\infty}^{\infty} \frac{e^{-i\omega t_r}}{(2i\omega_C + \omega)(\omega + i\omega_0)(\omega - i\omega_0)} d\omega. \end{aligned} \quad (\text{A10})$$

- (1) If $t_r > 0$: Consider the contour comprised of the real axis and the clockwise-oriented negative infinite semicircle. On the contour, the exponential phase factor in Eq. (A10) goes as

$$\exp(-i\omega t_r) = \exp(-i(R\cos\phi + iR\sin\phi)t_r). \quad (\text{A11})$$

For the semicircle, $\phi \in [\pi, 2\pi]$, so $\sin\phi < 0$ and $t_r > 0$. Exponential decay occurs and the integrand vanishes on the semicircle for $|\omega| = R \rightarrow \infty$.

- (2) If $t_r < 0$: Consider the contour comprised of the real axis and the counterclockwise-oriented positive infinite semicircle. On the contour, the exponential phase factor in Eq. (A10) goes again as

$$\exp(-i\omega t_r) = \exp(-i(R\cos\phi + iR\sin\phi)t_r). \quad (\text{A12})$$

For the semicircle, $\phi \in [0, \pi]$, so $\sin\phi > 0$ and $t_r < 0$. Exponential decay occurs and the integrand vanishes on the semicircle for $|\omega| = R \rightarrow \infty$.

Using cases 1 and 2, Eq. (A10) can be solved using the Cauchy integral formula. Beginning with $t_r > 0$, two poles are enclosed in the semicircle: one that originated from the coherence cutoff frequency, and the other that originated from the form factor. The Cauchy integral formula yields

$$\begin{aligned} rE(t, \theta_C) &= 2\hat{E}_0\omega_C\omega_0^2 \frac{d}{dt_r} \left(\frac{e^{-2\omega_C t_r}}{i^2(-2\omega_C + \omega_0)(-2\omega_C - \omega_0)} \right. \\ &\quad \left. + \frac{e^{-\omega_0 t_r}}{i^2(-\omega_0 + 2\omega_C)(-2\omega_0)} \right). \end{aligned} \quad (\text{A13})$$

Define the ratio of the cutoff frequencies: $\epsilon = \omega_0/\omega_C$. After evaluating the time derivatives, Eq. (A13) becomes

$$rE(t, \theta_C) = \hat{E}_0\omega_0^2 \left(\frac{e^{-2\omega_C t_r}}{(1 - \frac{\epsilon}{2})(1 + \frac{\epsilon}{2})} - \frac{e^{-\omega_0 t_r}}{(2)(1 - \frac{\epsilon}{2})} \right). \quad (\text{A14})$$

Expanding to linear order in ϵ , assuming $\epsilon < 1$, and recalling that $\omega_0^2 = \frac{2}{3}\omega_{\text{CF}}^2$,

$$rE(t, \theta_C) \approx \frac{1}{3}\hat{E}_0\omega_{\text{CF}}^2 \left(2e^{-2\omega_C t_r} - \left(1 + \frac{\epsilon}{2}\right)e^{-\omega_0 t_r} \right). \quad (\text{A15})$$

Turning to the case of $t_r < 0$, consider integrating Eq. (A10) along the contour comprised of the real axis and the counterclockwise-oriented positive infinite semicircle. The contour encloses one pole, and the exponent ensures convergence:

$$rE(t, \theta_C) = (2\pi i)\hat{E}_0(\pi)^{-1}i\omega_C\omega_0^2 \frac{d}{dt_r} \left(\frac{e^{\omega_0 t_r}}{(2i\omega_C + i\omega_0)(2i\omega_0)} \right). \quad (\text{A16})$$

After evaluating the derivative, the expression simplifies with $\epsilon = \omega_0/\omega_C$:

$$rE(t, \theta_C) = \frac{1}{2}\hat{E}_0\omega_0^2 \left(\frac{e^{\omega_0 t_r}}{1 + \frac{1}{2}\epsilon} \right). \quad (\text{A17})$$

Finally, using the same first-order approximation in ϵ as the $t_r > 0$ case,

$$rE(t, \theta_C) \approx \frac{1}{3}\hat{E}_0\omega_{\text{CF}}^2 \left(1 - \frac{1}{2}\epsilon \right) e^{\omega_0 t_r}. \quad (\text{A18})$$

Collecting the $t_r > 0$ and $t_r < 0$ results together,

$$rE(t, \theta_C) = \frac{1}{3}\hat{E}_0\omega_{\text{CF}}^2 \begin{cases} (1 - \frac{1}{2}\epsilon)e^{\omega_0 t_r}, & t_r < 0, \\ (2e^{-2\omega_C t_r} - (1 + \frac{1}{2}\epsilon)e^{-\omega_0 t_r}), & t_r > 0. \end{cases} \quad (\text{A19})$$

APPENDIX B: DETAILS OF THE OFF-CONE FIELD EQUATION DERIVATION

Using Tables II–IV, Eq. (A2) reduces to

$$\mathcal{E}(u, x) = f(u, x)g(u, x)(1 - h(u, x)). \quad (\text{B1})$$

Expanding to first order with respect to u near ($u = 1$) gives

$$\mathcal{E}(u, x) = \mathcal{E}(x, 1) + (u - 1)\dot{\mathcal{E}}(x, 1) + \mathcal{O}(u - 1)^2. \quad (\text{B2})$$

The first term is $fg(1 - h)$ evaluated at $u = 1$: $\exp(-y)$ (Table IV). The second term requires the first derivative of $\mathcal{E}(u, x)$ with respect to u , evaluated at $u = 1$:

$$\dot{\mathcal{E}}(u, x) = f\dot{g} + \dot{f}g - (f\dot{g}h + f\dot{g}h + \dot{f}gh) \quad (\text{B3})$$

$$\dot{\mathcal{E}}(1, x) = (f\dot{g} + \dot{f}g - (f\dot{g}h + f\dot{g}h + \dot{f}gh))|_{u=1} \quad (\text{B4})$$

The first derivatives of f , g , and h , evaluated at $u = 1$, are given in Table IV. Because $h(x, 1) = 0$, terms proportional to h will vanish. The result is

$$\dot{\mathcal{E}}(1, x) = \frac{1}{2} e^{-y} (2y + 2q - 1). \quad (\text{B5})$$

Inserting Eq. (B5) into Eq. (B2),

$$\mathcal{E}(u, x) = e^{-y} \left(1 + \frac{1}{2} (u - 1) (2y + 2q - 1) \right). \quad (\text{B6})$$

Using the definition of u (Table II), the result may be written as

$$\mathcal{E}(u, x) = e^{-y} \left(1 - \frac{1}{2} j\eta (2y + 2q - 1) \right). \quad (\text{B7})$$

Proceeding with the inverse Fourier transform of the $\hat{\theta}$ component,

$$rE(t, \theta) = \mathcal{F}^{-1} \left\{ E_0 \left(\frac{\omega}{2\pi} \right) \tilde{F} \psi \mathcal{E} \right\}. \quad (\text{B8})$$

Let $\eta = \omega/\omega_C$, $y = p\omega^2$ (Table II). Inserting the Taylor series for \mathcal{E} , the form factor \tilde{F} , and $\psi = -i \exp(ikr) \sin \theta$ (Sec. II), and following the same steps as the on-cone case produces

$$\begin{aligned} 2\pi r E(t, \theta) &= \frac{E_0 \omega_0^2 \sin(\theta)}{4\pi i \omega_C} \\ &\times \frac{d}{dt_r} \int_{-\infty}^{\infty} \frac{e^{-i\omega t_r - p\omega^2} (2i\omega_C + 2p\omega^3 + (2q - 1)\omega)}{\omega^2 + \omega_0^2} d\omega. \end{aligned} \quad (\text{B9})$$

Unlike the on-cone case, Eq. (B9) cannot be integrated with infinite semicircle contours, because the exponential term diverges along the imaginary axis far from the origin.

Let I_0 represent the constant term with respect to ω in the numerator:

$$I_0 = \int_{-\infty}^{\infty} \frac{e^{-i\omega t_r - p\omega^2} (2i\omega_C)}{\omega^2 + \omega_0^2} d\omega. \quad (\text{B10})$$

Further, let I_1 and I_3 represent the linear and cubic terms, respectively. Completing the square in the exponent of I_0 , with $\omega_1 = t_r/(2p)$, yields

$$I_0 = 2i\omega_C e^{-\frac{t_r^2}{4p}} \int_{-\infty}^{\infty} \frac{e^{-p(\omega+i\omega_1)^2}}{\omega^2 + \omega_0^2} d\omega. \quad (\text{B11})$$

Equation (B11) may be recast as the *line-broadening function*, H (DLMF 7.19, [39]) common to spectroscopy applications:

$$I_0 = 2\pi i \left(\frac{\omega_C}{\omega_0} \right) e^{-\frac{t_r^2}{4p}} H(\sqrt{p}\omega_0, i\sqrt{p}\omega_1). \quad (\text{B12})$$

Assume that $\omega > \omega_1$. This approximating step will be called the *symmetric approximation*:

$$I_0 \approx 2i\omega_C e^{-\frac{t_r^2}{2p}} \int_{-\infty}^{\infty} \frac{e^{-p\omega^2}}{\omega^2 + \omega_0^2} d\omega. \quad (\text{B13})$$

The result for I_0 involves the complementary error function (DLMF 7.7.1, [39]):

$$I_0 = 2i\omega_C e^{-\frac{t_r^2}{2p}} \pi \omega_0^{-1} e^{p\omega_0^2} \operatorname{erfc}(\sqrt{p}\omega_0). \quad (\text{B14})$$

The integrals I_1 and I_3 are zero by symmetry, with odd integrands over $(-\infty, \infty)$. Inserting the result for I_0 into Eq. (B9) and evaluating the derivative finishes the problem (see Sec. V).

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