


Testing the nonclassicality of spacetime: What can we learn from Bell–Bose *et al.*–Marletto–Vedral experiments?

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The Bose *et al.*–Marletto–Vedral (BMV) experiment [S. Bose *et al.*, *Phys. Rev. Lett.* **119**, 240401 (2017); C. Marletto and V. Vedral, *Phys. Rev. Lett.* **119**, 240402 (2017)] aims to prove that spacetime is nonclassical by observing entanglement generated by gravity. However, local hidden variable theories (LHVTs) can simulate the entangled correlations. We propose to extend the entanglement generated by the BMV experiment to distant quantum particles in a Bell experiment. Violating a Bell inequality would rule out LHVTs, providing a stronger proof of the nonclassicality of spacetime than the BMV proposal.

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I. INTRODUCTION

We do not yet know how to unify our two most fundamental theories of physics, quantum theory (QT) and general relativity (GR). Each is supported by overwhelming experimental evidence in its respective domain. However, they are formulated very differently. GR is an essentially classical theory in which the stress-energy tensor, the spacetime geometry and all physical quantities take well-defined values; QT predicts that distinct mass-energy configurations can be in quantum superposition. One popular approach to resolving this tension is to seek a quantum theory of gravity, such as string theory [1] or loop quantum gravity [2]. It has even been claimed that there is no consistent alternative (e.g., [3–6]). However, this has also often been questioned (e.g., [7]). In particular, the argument of Ref. [4] has been refuted (e.g., [8–11]). It seems the question can only be settled empirically.

In their simplest form, semiclassical gravity models [12–15] are defined by taking

$$G_{\mu\nu} = \kappa \langle \hat{T}_{\mu\nu} \rangle, \quad (1)$$

where the left-hand side is the Einstein tensor of classical spacetime, and the right-hand side is the expectation value of the stress-energy tensor of quantum matter propagating in that spacetime, with κ being a proportionality constant. If this is taken to be defined by Everettian quantum theory with purely unitary evolution, Eq. (1) is inconsistent with observation (e.g., [16]); it is also unclear that it defines a

consistent theory. However, it remains possible that (1) holds in some regime (e.g., [17–20]). For example, given a version of quantum theory with explicit localized collapses, a classical gravitational field could couple to the local quantum state, which is defined by the initial conditions, evolution, and collapse events (only) in the causal past of the relevant point [11,20,21]. Another possibility is that quantum superpositions of sufficiently distinct energy-mass configurations are dynamically suppressed (e.g., [22–29]).

All these options are problematic. For example, a recent experiment [30] gives strong evidence against the Diósi–Penrose proposals [24–26,28,29] for gravitationally induced collapse and suggests that a radically new approach may be needed to pursue this idea. However, quantum gravity theories also have well known problems (see e.g., [31] for some discussion).

Recent proposals for table-top tests of quantum gravity, beginning with so-called Bose *et al.*–Marletto–Vedral (BMV) experiments [32,33], aim to give strong evidence as to whether gravity is mediated by quantum information exchange. This would also give indirect evidence that spacetime is quantum [34]. The essential idea of BMV’s proposal is to place two adjacent mesoscopic masses in position superposition states and allow them to fall along paths such that one pair of paths is significantly closer than the others, before recombining the paths interferometrically. With appropriate masses, separations, and fall times, which it is hoped will be experimentally feasible in the foreseeable future, the Schrödinger evolution with a Newtonian potential implies that an entangled final state can be generated from a separable initial state. As the position degrees of freedom are correlated with internal degrees of freedom, this entanglement can be tested by measuring entanglement witnesses.

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Refinements to the experiments have since been proposed [35], as have alternative experiments aiming to witness (non)quantum behavior of gravity by other methods (e.g., [29,36–38]). There is also ongoing debate (e.g., [37,39,40]) over how definitively the quantum nature of gravity would be demonstrated by the generation of entanglement in a BMV experiment.

A separate issue, our focus here, is how definitively passing the type of entanglement witness test proposed by BMV would establish that entanglement had indeed been generated. There is clearly a logical loophole in this inference, since we know that local hidden variable theories can simulate arbitrary entangled correlations in any experiment—such as the BMV proposals—that does not involve measurements in spacelike separated wings. Of course, Bell experiments give very strong—even if not yet loophole-free [41,42]—evidence against local hidden variable theory explanations for correlations between measurements on entangled matter. However, Bell experiments to date have not tested states in which (according to quantum gravity intuitions) entanglement is generated by gravitational interactions. As already noted, there is some motivation for considering models in which classical degrees of freedom associated with gravity couple to, and might indeed contain complete information about, the local quantum state. In principle, this would allow classical simulation of the correlations that a quantum analysis would ascribe to gravitationally generated entanglement.

We are not aware of any convincing model that would reproduce the correct correlations for all the entangled states that could (on a quantum analysis) be generated by varying the experimental parameters. Still, BMV experiments aim to establish a fundamental feature of nature, and in our view it is worth striving to eliminate loopholes in their interpretation, for reasons similar to those motivating the ongoing quest to remove Bell experiment loopholes. Although it is a theoretically familiar idea that the space and time we inhabit somehow emerge from a more fundamental quantum description, it remains an extraordinary claim, which justifies extraordinary care in assessing evidence.

We propose here extending the entanglement generated in a BMV experiment to distant quantum particles in a Bell experiment. For the reasons just given, this would provide a more definitive test of the nonclassicality of gravity and spacetime.

II. A BELL–BOSE *ET AL.*-MARLETTO-VEDRAL EXPERIMENT

We propose an experimental test for the nonclassicality of gravity comprising three general steps. Broadly, these steps are the following (see Fig. 1).

- (1) Perform the BMV experiment [32,33] with a pair of mesoscopic masses with spin. Let us call these systems “particle 1” and “particle 2”, and their spin degrees of freedom “ S_1 ” and “ S_2 ”. Arrange the

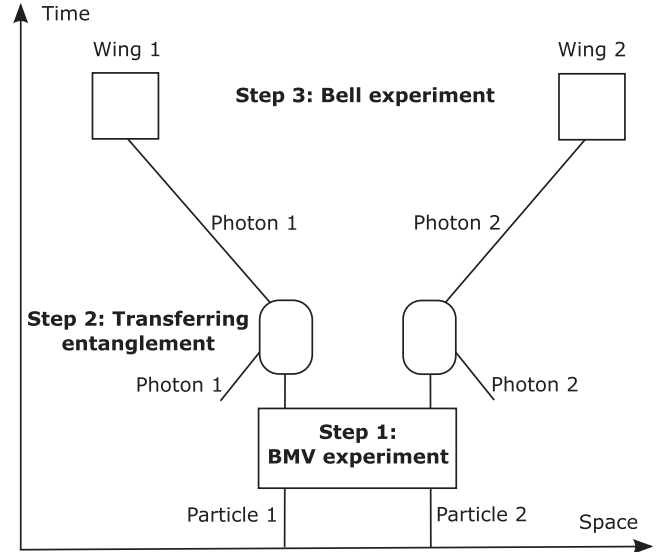


FIG. 1. Our proposed BBMV experiment. We illustrate our proposed BBMV experiment in a spacetime diagram in $1+1$ dimensions in a background Minkowski spacetime. Our proposed experiment comprises three broad steps. In Step 1 (large rectangle), the BMV experiment [32,33] is applied on Particles 1 and 2 (vertical lines). The entanglement generated via the BMV effect is transferred to Photons 1 and 2 (diagonal lines) in Step 2 (rounded rectangles). In Step 3 (small squares), a Bell experiment [43] is applied on Photons 1 and 2, testing (for example) CHSH inequalities [45]. If conditions 1–4 hold and a Bell inequality is violated then we can exclude explanations via local hidden variables, obtaining a stronger proof for the nonclassicality of the gravitational field than the original BMV experiment. The illustration is not at scale. In particular, we expect the distance between the wings in the Bell experiment to be much larger than the distance between the particles in the BMV experiment.

experiment so that the particles’ states are initially unentangled, and the only non-negligible interaction between the particles during the experiment is the gravitational field. The hypothesis that gravity is quantum implies that S_1 and S_2 become entangled due to the gravitational interaction of the particles’ masses [32,33,46]. Assuming this hypothesis is correct, arrange the experiment so that the final state between S_1 and S_2 is approximately the maximally entangled singlet state:

$$\rho_{S_1 S_2}^{(1)} \approx (|\Psi^-\rangle\langle\Psi^-|)_{S_1 S_2}, \quad (2)$$

where

$$|\Psi^-\rangle = \frac{1}{\sqrt{2}}(|0\rangle|1\rangle - |1\rangle|0\rangle). \quad (3)$$

- (2) Transfer the quantum state from the spin degrees of freedom S_i of particle i to some degrees of freedom P_i (e.g., polarization) of a photon, which we call “photon i ”, for $i = 1, 2$. Arrange the experiment so

that the interaction between particle 1 (respectively 2) and photon 1 (2) is local, i.e., without interacting with particle 2 (1) or photon 2 (1). If at the end of step 1, the joint quantum state of the particles' spins S_1 and S_2 satisfies (2), then at the end of this step the joint quantum state of the photons' degrees of freedom P_1 and P_2 is

$$\rho_{P_1 P_2}^{(2)} \approx (|\Psi^-\rangle\langle\Psi^-|)_{P_1 P_2}. \quad (4)$$

- (3) Send photons 1 and 2 to distant wings and implement a Bell [43] experiment on systems P_1 and P_2 at spacelike separation, testing (for example) Clauser-Horne-Shimony-Holt (CHSH) [45] inequalities.

We call this a Bell-Bose *et al.*-Marletto-Vedral (BBMV) experiment. We emphasize the conditions that need to be satisfied experimentally.

Condition 1. At the beginning of step 1, the states of particles 1 and 2 are initially unentangled.

Condition 2. During step 1, the only non-negligible interaction between particles 1 and 2 is gravity.

Condition 3. During step 2, transferring the quantum state from the spin S_i of particle i to the degrees of freedom P_i of photon i does not include any interaction between the joint system of particle i and photon i with the joint system of particle \bar{i} and photon \bar{i} , either directly or indirectly via another system, for $i = 1, 2$.

Condition 4. The measurements in wing i in the Bell experiment of step 3 are applied on the degrees of freedom P_i of photon i , for $i = 1, 2$.

Suppose that the experiment is repeated $N \gg 1$ times and that conditions 1–4 hold in each run of the experiment. In this case, observing correlations approximately consistent with CHSH measurements on a quantum singlet, and so violating the CHSH inequality, in our BBMV experiment would imply (modulo any remaining loopholes) that the correlations obtained cannot be explained by locally causal hidden variable models [43,45]. We see relatively little motivation for considering nonlocally causal hidden variable models in the context of unifying quantum theory and general relativity, and in particular for considering such models as an explanation for correlations observed in the specific BBMV experiment described. The natural inference would thus be that the correlations are indeed generated by measurements on an approximate singlet state, and hence that the BMV subexperiment did indeed generate entanglement.

III. DISCUSSION

BMV experiments involve testing for correlations that would arise from measurements on the entangled quantum states that should be generated, according to nonrelativistic quantum theory using Newtonian gravitational potentials. In principle, such correlations could also be produced by a nonquantum theory of gravity involving local hidden

variables. If our proposed experiment produced results consistent with quantum theory and the BMV analyses, it would exclude the latter explanation, and hence provide stronger evidence for quantum gravity. This is not to diminish the importance of BMV's [32,33] crucial insight. Arguably, given most theorists' Bayesian priors, our proposed experiment would further (beyond a BMV experiment) enhance their credence in quantum gravity only marginally. Still, given our incomplete understanding of nature, empirical proof is preferable to confident priors.

Ideally, all else being equal, the Bell experiment of step 3 should be as loophole free as possible. However, we see more motivation for closing some loopholes than others. There is very strong motivation for closing the locality loophole [43], since BBMV experiments are motivated by the concern that gravitational effects might be mediated by locally causal hidden variables. Ideally, the collapse locality loophole [41,42] should also be closed, since one reason for considering this loophole is the possibility that measurement events occur only when their outcomes leave a record in the gravitational field and hence, in order to verify genuine nonlocally causal correlations, we need to ensure these records are created in spacelike separated regions. Although this is more challenging, recent advances in experimental quantum information science in space (e.g., [47–52]) suggest that the collapse locality loophole could be closed in the foreseeable future [42].

It would also be worthwhile to close the freedom-of-choice loophole [43], to eliminate any possibility that information about choices made at earlier times on one wing propagates via locally causal hidden variables (which again might be associated with the gravitational field) to influence outcomes on the other. We see less motivation for closing the detector efficiency loophole [53]; we find it harder to imagine a plausible theory in which locally causal hidden variables affect the behavior of detectors in BBMV experiments but not in standard Bell experiments.

It is hoped that the BMV experiment could be feasible in the foreseeable future (e.g., [32,33,35]). Our proposed BBMV experiment appears not substantially more challenging to implement than the BMV experiment and a Bell experiment. State of the art techniques for spin-photon coupling (e.g., [54,55]) suggest that step 2 in our experiment could be possible in practice.

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