

Cosmology with standard sirens at cosmic noon

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Gravitational waves (GWs) from merging black holes and neutron stars directly measure the luminosity distance to the merger, which, when combined with an independent measurement of the source’s redshift, provides a novel probe of cosmology. The proposed next generation of ground-based GW detectors, Einstein Telescope and Cosmic Explorer, will detect tens of thousands of binary neutron stars (BNSs) out to cosmological distances ($z > 2$), beyond the peak of the star formation rate (SFR), or “cosmic noon.” At these distances, it will be challenging to measure the sources’ redshifts by observing electromagnetic (EM) counterparts or statistically marginalizing over a galaxy catalog. In the absence of an EM counterpart or galaxy catalog, Ding *et al.* [*J. Cosmol. Astropart. Phys.* **04** (2019) 033] showed that theoretical priors on the merger redshift distribution can be used to infer parameters in a w CDM (Cold Dark Matter) cosmology. We argue that in the BNS case, the redshift distribution will be measured by independent observations of short gamma ray bursts (GRBs), kilonovae, and known BNS host galaxies. In particular, the peak redshift will provide a clear feature to compare against the peak distance of the GW source distribution and reveal the underlying redshift-distance relation. We show that, in addition to measuring the background cosmology, this method can constrain the effects of dark energy on modified GW propagation. As a simple example, we consider the case in which the BNS rate is *a priori* known to follow the SFR. If the SFR is perfectly known, $\mathcal{O}(10,000)$ events (to be expected within a year of observation with Cosmic Explorer) would yield a subpercent measurement of the combination $H_0^{2.8}\Omega_M$ in a flat Λ CDM model. Meanwhile, fixing H_0 and Ω_M to independently inferred values, this method may enable a 5% measurement of the dark energy equation of state parameter w in a w CDM model. Fixing the background cosmology and instead probing modified GW propagation, the running of the Planck mass parameter c_M may be measured to ± 0.02 . Although realistically, the redshift evolution of the merger rate will be uncertain, prior knowledge of the peak redshift will provide valuable information for standard siren analyses.

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I. INTRODUCTION

As first pointed out by Schutz [1], gravitational waves (GWs) from the coalescence of two compact objects, like black holes and neutron star mergers, provide an absolute distance measurement to the source. In an analogy to “standard candles,” like type Ia supernovae, these GW sources are known as “standard sirens” [2]. If the redshift corresponding to the GW source can also be determined, it is possible to constrain the distance-redshift relation and thereby measure cosmological parameters. However, the redshift cannot be directly extracted from the GWs, because the redshift of the GW signal is degenerate with the mass of the system. Standard siren cosmology therefore relies on external data to infer the redshift of the GW source.

The most straightforward approach to determine the redshift of a GW observation is to identify an associated

electromagnetic (EM) counterpart, like a short gamma ray bursts (GRB) or a kilonova, which in turn allows for a unique host galaxy identification and redshift determination [2–4]. A counterpart standard siren measurement was first carried out following the spectacular multimessenger detection of the binary neutron stars (BNS) merger GW170817 [5,6], yielding a measurement of the Hubble constant $H_0 = 70_{-8}^{+12}$ km s⁻¹ Mpc⁻¹ [7]. Chen *et al.* [8] and Feeney *et al.* [9] showed that 50 detections of BNS mergers by the Advanced LIGO [10] and Virgo [11] GW detector network with associated EM counterparts will enable a $\sim 2\%$ measurement of H_0 , which would provide an important test of the Λ CDM (Cold Dark Matter) cosmological model and may help shed light on the persistent H_0 tension [12–17].

Nevertheless, the majority of GW events do not have identified EM counterparts. In the absence of a counterpart,

it is possible to statistically marginalize over the redshifts of all of the potential host galaxies in the GW localization volume using a galaxy catalog [1,18–20]. This statistical standard siren approach has been applied to several GW events [21–25]. The most promising dark sirens for the statistical method are nearby, well-localized events, where the number of galaxies in the volume is relatively small and available galaxy catalogs are relatively complete [8,23,25,26]. When catalogs are incomplete but GW events are well-localized, it may be possible to compare the spatial clustering of GW sources and galaxies as a function of redshift to infer cosmological parameters [18,27–30]. Finally, in the absence of counterparts or catalogs, several authors have proposed GW-only standard siren analyses. Known properties of the source population, such as features in the source-frame mass distribution [31–36] or knowledge of the neutron star equation of state [37,38], can be used to extract the redshift from the observed GW frequency. Ding *et al.* [39] pointed out that even if the redshifts of individual GW events cannot be identified, it is possible to extract cosmological information from a population of standard sirens if their redshift distribution is theoretically known from population synthesis simulations.

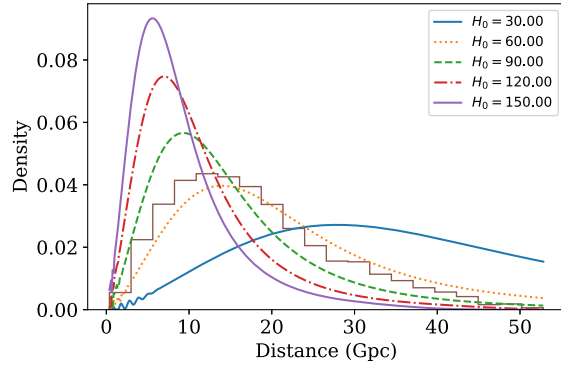
In this study, we build on Ding *et al.* [39] and further explore the potential of standard siren cosmology without counterparts. We argue that external EM observations, not necessarily associated with GW events, provide a measurement of the BNS redshift distribution that can be leveraged in a standard siren measurement. For example, if the BNS merger rate is known to follow the star formation rate (SFR) with short typical time delays [40–42], we will know that there is an abundance of BNS host galaxies near the peak of the SFR at $z \sim 2$ [43,44] without comparing a galaxy catalog against GW events.

This method would be particularly relevant for the next generation of ground-based GW observatories, the proposed detectors Cosmic Explorer [45] and Einstein Telescope [46,47], which are currently under consideration. These third-generation (3G) detectors would dramatically increase the distance out to which BNS mergers can be observed, from a few hundred Mpc with current detectors [48,49] to tens of Gpc [50–52]. The 3G detectors will thus most likely observe these mergers past the peak redshift of the merger rate distribution. Depending on the detector network, the BNS rate, and the mass distribution, they will observe on order of 10^5 BNSs annually [52]. Although some of these GW signals will overlap, the parameters of these sources can nevertheless be measured reliably [53–55]. This large GW dataset will provide a novel probe of the high-redshift universe [56]. For example, assuming the distance-redshift relation is known, the distribution of their luminosity distances will enable precise measurements of the time delay distribution between star formation and compact object merger [33,57,58]. Another natural application of 3G detectors is standard siren cosmology

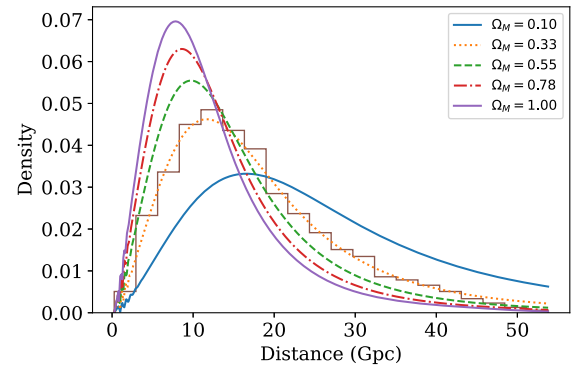
out to high redshifts, which can provide independent constraints on dark energy, alternative cosmological models and modified gravity [33,35,50,52,59–66]. However, at $z > 1$, it will become increasingly difficult to observe EM counterparts, both because of their reduced apparent brightness and the large GW localization areas [67]. The statistical method will also face challenges, because galaxy catalogs will be increasingly incomplete at high redshift. GW-only methods drawing on knowledge of the source-frame population, such as the BNS mass distribution [32,33] or the pair-instability feature in the BBH mass distribution [34,35] may prove useful; the latter technique may even provide an $\mathcal{O}(10\%)$ measurement of the dark energy equation of state with the current GW detector network [34]. However, these methods rely on some understanding of the evolution of the source population with redshift, which remains observationally and theoretically uncertain [68].

These anticipated challenges for standard siren cosmology in the 3G era motivate us to consider supplementary sources of redshift information. Although we cannot always observe the EM counterpart to a GW detection of a BNS, we nevertheless have an external EM sample of short GRBs and kilonovae, the progenitors of which are probably BNS (or neutron-star black hole) mergers [6]. This sample will grow in the coming years with the improved sensitivity of upcoming/proposed observing facilities like the Vera Rubin Observatory, JWST and the Roman Space Telescope for observing kilonovae [69] and BurstCube, SVOM, THESEUS and ULTRASAT for observing GRBs and afterglows, among many other proposed missions. We expect that external EM observations of short GRBs and kilonovae will constrain the number density of these sources as a function of redshift, either by accumulating a large sample of sources with known redshifts, as identified through afterglow or host galaxy spectroscopy, or jointly fitting the GRB observed flux distribution to the underlying luminosity function and redshift distribution [40,41,70–73]. Even observations within a limited redshift range can provide valuable information about the redshift evolution of the merger rate if host galaxies can be identified [74]. The properties of host galaxies (e.g., their masses and specific star formation rates) can be used to measure the time delay distribution [75–77] and therefore, assuming some knowledge about the cosmic SFR, the BNS merger rate as a function of redshift. This measurement may be assisted by theoretical guidance that the BNS progenitor formation rate probably traces the SFR, independent of the (largely uncertain) metallicity evolution [78]. Additional information about the BNS merger rate evolution can be gained by studying the Galactic evolution of r-process elements [42].

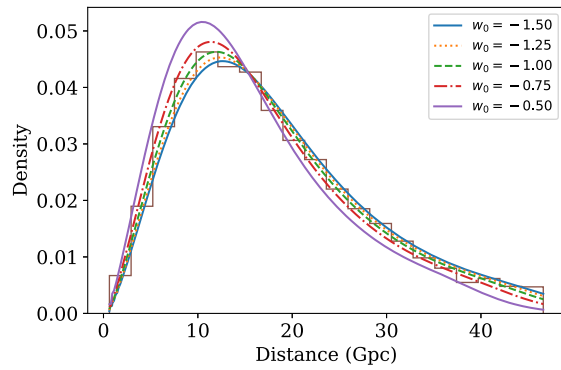
Figure 1 illustrates how, for a fixed merger rate evolution $\mathcal{R}(z)$, the luminosity distance distribution, as measured with GW standard sirens, depends on cosmology. For



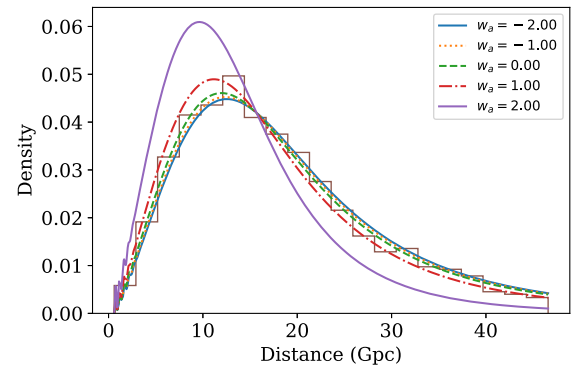
(a) Distribution of source luminosity distances for different values of the Hubble constant H_0 . Smaller values of H_0 result in larger observed distances on average.



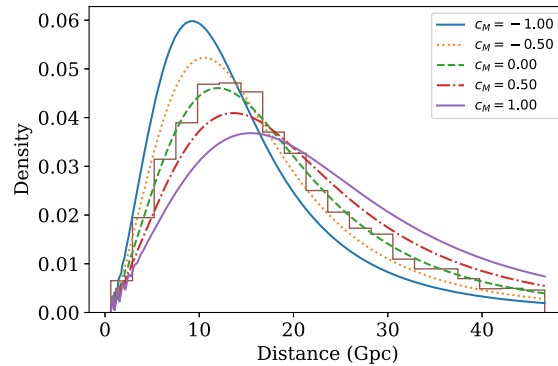
(b) Distribution of source luminosity distances for different values of the matter density Ω_M . Smaller values of Ω_M result in larger observed distances on average.



(c) Distribution of source luminosity distances for different values of the dark energy equation of state parameter w_0 (with w_a fixed to zero). The effect is small compared to the influence of Ω_M and H_0 , but visible. Smaller (more negative) values of w_0 result in larger observed distances on average.



(d) Distribution of source luminosity distances for different values of the dark energy equation of state parameter w_a (with w_0 fixed to -1). The luminosity distance distribution is much more sensitive to variations towards positive w_a than negative w_a . Positive w_a results in smaller observed distances on average.



(e) Distribution of GW luminosity distances for different values of c_M , the parameter controlling the running of the Planck mass in a modified gravity theory, in a flat Λ CDM background cosmology. Positive c_M results in larger observed distances on average.

FIG. 1. Effect of cosmological and modified gravity parameters on the GW luminosity distance distribution, for a fixed merger rate evolution $\mathcal{R}(z)$ that follows the Madau-Dickinson SFR. We assume a flat background $w_0 w_a$ CDM cosmology, with possible modifications to GW propagation parametrized by c_M . Each panel varies one parameter at a time, fixing the other parameters to a “concordance cosmology” of Λ CDM with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$. Histograms show simulated distance measurements with concordance cosmology.

simplicity, we fix the underlying merger rate to the Madau-Dickinson SFR [see Eq. (12)]. Building upon this intuition, in the following we explore how comparing GW observations of luminosity distances to a known redshift distribution can be used to extract cosmological and modified gravity parameters, including the Hubble constant, the matter density and the dark energy equation of state in a flat w_0w_a -CDM cosmology [79,80], and the running of the Planck mass following the framework of Lagos *et al.* [81]. Throughout, we fix the geometry of the Universe to be spatially flat, motivated by tight constraints on the curvature from cosmic microwave background (CMB) and baryon acoustic oscillation (BAO) data [82].

The general idea is that the GW luminosity distance D_L^{GW} , as a function of the redshift z , depends on these constants that we wish to measure. In general relativity (GR), the luminosity distance D_L^{GW} that is extracted from the GW signal is the same luminosity distance D_L^{EM} that governs electromagnetic sources, and we have [83]

$$D_L^{\text{EM}}(z) = (1+z) \frac{c}{H_0} \int_0^z \frac{dz'}{E(z')}, \quad (1)$$

where c is the speed of light, H_0 is the Hubble constant, and assuming a flat universe, $E(z)$ is

$$E(z) = \sqrt{\Omega_M(1+z)^3 + (1-\Omega_M)I(z)}, \quad (2)$$

with Ω_M being the dimensionless matter density today, $(1-\Omega_M)$ the dark energy density (in a flat universe with a negligible radiation density today), and $I(z)$, in the w_0w_a (CPL model) for the dark energy equation of state, given by [79,80,84]

$$I(z) = (1+z)^{3(1+w_0+w_a)} \exp\left(-3w_a \frac{z}{1+z}\right). \quad (3)$$

The above reduces to w CDM for $w_a = 0$ and Λ CDM for $w = -1$, $w_a = 0$. We use ASTROPY [84] for cosmological calculations.

Modified gravity theories [85–87], including models of dynamical dark energy, may alter the amplitude of the GW signal compared to GR in addition to altering the background cosmology away from Λ CDM, so that the measured D_L^{GW} differs from the electromagnetic luminosity distance [81,88–104]. The effect of the GR deviations on GW propagation may be much more significant, and therefore easily measurable with GW events, than the modifications to the background expansion [105]. While the multimessenger detection of GW170817 has put tight constraints on the speed of GW propagation, deviations affecting the GW amplitude remain relatively poorly constrained [106]. In this paper, we consider the example of GW damping caused by an effective running of the Planck mass. Following

Lagos *et al.* [81], we model the time evolution of the Planck mass with an additional parameter c_M on top of the background cosmology, assumed to follow flat Λ CDM. The GW luminosity distance D_L^{GW} is then the product of Eq. (1) (with $w = -1$, $w_a = 0$ for Λ CDM) with the extra factor,

$$\frac{D_L^{\text{GW}}}{D_L^{\text{EM}}} = \exp\left(\frac{c_M}{2(1-\Omega_M)} \ln \frac{1+z}{(\Omega_M(1+z)^3 + 1 - \Omega_M)^{1/3}}\right), \quad (4)$$

where $c_M = 0$ reduces to GR, i.e., $D_L^{\text{GW}} = D_L^{\text{EM}}$.

The remainder of the paper is organized as follows. Section II describes the statistical framework that we apply to simulated GW data. We show the results of the simulations in terms of projected constraints in the cosmological parameters in Sec. III. We conclude in Sec. IV.

II. METHODS

This section describes the analysis and simulation methods. We derive the hierarchical Bayesian likelihood for the joint inference of the cosmological parameters and the redshift distribution parameters in Sec. II A and describe the application of this likelihood to simulated data in Sec. II B.

A. Statistical framework

We assume that the underlying redshift distribution of sources can be described by some parameters λ with some additional possible dependence on the cosmological parameters \mathcal{H} . We write this as $p(z|\lambda, \mathcal{H})$. As a probability density function, $p(z|\lambda, \mathcal{H})$ integrates to unity over $0 < z < z_{\text{max}}$. The population-level parameters are therefore λ and \mathcal{H} . Often the redshift distribution is expressed as a merger rate density $\mathcal{R}(z)$, which refers to the number of mergers per comoving volume and source-frame time, and can be equivalently written as $\frac{dN}{dV_c dt_s}$ where V_c is the comoving volume and t_s is the source-frame time. The redshift distribution $p(z)$ is related to the redshift-dependent merger rate density $\mathcal{R}(z)$ by

$$p(z) \propto \mathcal{R}(z) \frac{dV_c}{dz} \frac{1}{1+z}. \quad (5)$$

We note that the conversion between $\mathcal{R}(z)$ and $p(z)$ depends on the differential comoving volume element $\frac{dV_c}{dz}$, which depends on cosmology. Assuming a flat universe [83],

$$\frac{dV_c}{dz} = \frac{c}{H_0} \frac{D_L^{\text{EM}}(z)^2}{(1+z)^2 E(z)}, \quad (6)$$

with $D_L^{\text{EM}}(z)$ given by Eq. (1) and $E(z)$ given by Eqs. (2)–(3). Depending on the type of observations, the measurement of

$p(z)$ and/or $\mathcal{R}(z)$ may depend on the assumed cosmology. If we have a redshift catalog of sources; i.e., the number of sources per redshift, we have a direct measurement of $p(z|\lambda)$ independent of cosmology. However, if we use observed fluxes to reconstruct the redshift evolution, we may measure $\mathcal{R}(z)$ more directly. The method described below applies to either scenario, but in our simulations, we consider the case where a measurement of $\mathcal{R}(z)$ is available.

We use a hierarchical Bayesian framework [107–109] to write the likelihood of the data d_i from event i , given the population-level parameters, as

$$\begin{aligned} p(d_i|\lambda, \mathcal{H}) &= \int_0^{z_{\max}} p(d_i, z_i|\lambda, \mathcal{H}) dz_i \\ &= \int_0^{z_{\max}} p(d_i|D_L(z_i, \mathcal{H})) p(z_i|\lambda, \mathcal{H}) dz_i, \end{aligned} \quad (7)$$

where $D_L(z_i, \mathcal{H})$ denotes the luminosity distance corresponding to the redshift z_i and the cosmology \mathcal{H} . For simplicity of notation, we use D_L to denote the GW luminosity distance D_L^{GW} throughout, even when we consider modifications to GR [e.g., Eq. (4)]. In the above we have implicitly marginalized over any other parameters of the GW signal, so that the marginal likelihood of d_i depends only on the GW luminosity distance $D_L(z_i, \mathcal{H})$. In reality, the GW data also depends on the detector-frame (redshifted) masses of the source; this is discussed further below.

In the presence of GW selection effects, we must modify the likelihood of Eq. (7) to account for the fact that some mergers do not produce detectable data d_i . If only data passing some threshold d^{thresh} are detected, the likelihood from each event must be normalized by a factor $\beta(\lambda, \mathcal{H})$ [8,109],

$$\beta(\lambda, \mathcal{H}) = \int_{d>d^{\text{thresh}}} \int_0^{z_{\max}} p(d|D_L(z, \mathcal{H})) p(z|\lambda, \mathcal{H}) dz dd. \quad (8)$$

The single-event likelihood, corrected for selection effects, is then

$$p(d_i|\lambda, \mathcal{H}) = \frac{\int_0^{z_{\max}} p(d_i|D_L(z_i, \mathcal{H})) p(z_i|\lambda, \mathcal{H}) dz_i}{\int_{d>d^{\text{thresh}}} \int_0^{z_{\max}} p(d|D_L(z, \mathcal{H})) p(z|\lambda, \mathcal{H}) dz dd}. \quad (9)$$

This differs from the likelihood used in Ding *et al.* [39], which incorporated selection effects by replacing the astrophysical redshift distribution $p(z|\lambda, \mathcal{H})$ with the redshift distribution of detected GW events; see Mandel *et al.* [109] for a derivation of the hierarchical Bayesian likelihood in the presence of selection effects.

The total likelihood of N GW events with data \mathbf{d} is the product of the individual-event likelihoods of Eq. (9),

$$p(\mathbf{d}|\lambda, \mathcal{H}) = \prod_{i=1}^N p(d_i|\lambda, \mathcal{H}). \quad (10)$$

Using Bayes' rule, we get the posterior on the cosmological parameters \mathcal{H} , given some prior $p_0(\mathcal{H})$,

$$p(\mathcal{H}|\mathbf{d}, \lambda) \propto p(\mathbf{d}|\lambda, \mathcal{H}) p_0(\mathcal{H}). \quad (11)$$

In the above, we have made the simplifying assumption that the data (and their detectability) depend on the source's redshift only through the GW luminosity distance. This is a simplification because in reality, the amplitude and frequency of a signal also depends on the source's redshifted masses and spins; in fact, if we have prior knowledge about the source-frame mass distribution, observing the redshifted masses can by itself probe the distance-redshift relationship [32,33]. Nevertheless, because we wish to isolate the information available from the luminosity distance distribution alone, for this proof-of-principle study we approximate that the GW data depends only on the observed luminosity distance. The masses $m_1(1+z)$ and $m_2(1+z)$ can be easily added into the likelihood of Eq. (9) by considering the GW likelihood $p(d|D_L(z, \mathcal{H}), m_1(1+z), m_2(1+z))$ and a population model $p(m_1, m_2, z|\lambda)$.

We have also ignored the additional distance uncertainty due to the effects of weak gravitational lensing, which will contribute an additional 1σ uncertainty of $\sim 0.05z$ to the measured distance depending on the source redshift z . If the distribution of lensing magnifications is known, this contribution can be marginalized over in the GW likelihood without affecting the rest of our formalism [2,50,59,110,111]. The statistical uncertainties we assume for mock data in the following subsection are large enough to encompass this additional contribution. Alternatively, one can simultaneously fit for the magnification distribution or power spectrum as a function of redshift, which may provide useful constraints on large-scale structure [100,112,113]. An additional source of uncertainty will be the calibration uncertainty due in the detector response. This will likely contribute a systematic uncertainty that will limit the accuracy of any standard siren cosmological analyses.

B. Simulations

We apply the likelihood analysis described in the previous Sec. II A to mock data. For simplicity, we assume that the evolution of the merger rate is perfectly known to follow the Madau-Dickinson SFR [43], peaking at $z \sim 2$,

$$\mathcal{R}(z) \propto \frac{(1+z)^{2.7}}{1 + \left(\frac{1+z}{2.9}\right)^{5.6}}, \quad (12)$$

and so the redshift distribution follows,

$$p(z|\lambda, \mathcal{H}) = A \frac{dV_c}{dz} \frac{1}{1+z} \frac{(1+z)^{2.7}}{1 + \left(\frac{1+z}{2.9}\right)^{5.6}}, \quad (13)$$

where A is a normalization constant ensuring that the redshift distribution integrates to unity over the range $0 < z < z_{\max}$. We take $z_{\max} = 8$, which ensures that it is larger than the maximum detected BNS distance for any choice of cosmological parameters in our prior. If the maximum astrophysical merger redshift is within the GW detector horizon, it may serve as another feature that can be leveraged for cosmological analyses. We stress that in reality, we do not expect the redshift distribution to be known perfectly, so that instead of using a δ -function prior on λ as we effectively assume here, future measurements will use a posterior probability distribution on λ inferred from external observations.

For our simulations, we fix a cosmology and draw redshifts z_i from the redshift distribution of Eq. (13). Under the fixed cosmology, this gives us the true luminosity distances D_L^i of our mock sources. We then simulate measurement uncertainty, drawing observed distances $D_L^{\text{obs},i}$ assuming that the GW distance likelihood follows a lognormal distribution with roughly 10% measurement uncertainty (see Appendix B of Safarzadeh *et al.* [58]),

$$\log D_L^{\text{obs},i} \sim \mathcal{N}(\mu = \log D_L^i, \sigma = 0.1), \quad (14)$$

where $\mathcal{N}(\mu, \sigma)$ denotes the normal distribution with mean μ and standard deviation σ . In other words, we write the GW likelihood $p(d_i|D_L(z, \mathcal{H}))$ of Eq. (9) as

$$p(d_i|D_L(z, \mathcal{H})) = p(D_L^{\text{obs},i}|D_L(z, \mathcal{H})) \quad (15)$$

$$\propto \frac{1}{D_L^{\text{obs},i}} \exp\left(-\frac{1}{2} \left(\frac{\log D_L^{\text{obs},i} - \log D_L(z, \mathcal{H})}{0.1}\right)^2\right). \quad (16)$$

This is a conservative assumption compared to parameter estimation simulations and Fisher matrix analyses [52,58]. Next we apply selection effects. We neglect the effects of the sky-dependent GW detector sensitivity and detector-frame mass (see the discussion in the previous subsection) and simply assume that GW sources are detected if and only if their observed distance is within some maximum D_L^{max} . We throw out all simulated $D_L^{\text{obs},i} > D_L^{\text{max}}$ as below the detection threshold. As the observed luminosity distance includes a log-normal error term, the detection probability as a function of the true luminosity distance follows a smooth sigmoid function. The detectability of BNS mergers as a function of distance for 3G observatories has large uncertainties, stemming from the BNS mass distribution and details about the 3G detector network. We bound this uncertainty by exploring two choices for the D_L^{max} parameter, 20 Gpc and 40 Gpc. These roughly correspond to Cosmic Explorer’s 50% “response distance,”

or the distance at which 50% of sources are detectable [49], for binaries with total source-frame masses of $3 M_\odot$ and $4 M_\odot$, respectively (see Fig. 1 of Hall and Evans [51], assuming a *Planck* 2018 cosmology).

Again writing $p(d|D_L(z, \mathcal{H})) = p(D_L^{\text{obs}}|D_L(z, \mathcal{H}))$, Eq. (8) then becomes

$$\beta(\lambda, \mathcal{H}) = \int_0^{D_L^{\text{max}}} \int_0^{z_{\max}} p(D_L^{\text{obs}}|D_L(z, \mathcal{H})) p(z|\lambda) dz dD_L^{\text{obs}}. \quad (17)$$

Under the assumption that $p(D_L^{\text{obs}}|D_L)$ is a log-normal distribution, we can simplify the integral over D_L^{obs} ,

$$\beta(\lambda, \mathcal{H}) = \int_0^{z_{\max}} \frac{1}{2} \left(1 + \text{erf}\left(\frac{\log D_L^{\text{max}} - \log D_L(z, \mathcal{H})}{\sqrt{2}\sigma}\right)\right) \times p(z|\lambda) dz, \quad (18)$$

where $\text{erf}(x)$ is the error function, and we have picked $\sigma = 0.1$.

For all the $D_L^{\text{obs},i}$ that are “detected,” we compute the likelihood of Eq. (9). The final posterior probability on the cosmological parameters \mathcal{H} is proportional to the product of these likelihoods multiplied by the prior on \mathcal{H} , as in Eq. (11).

III. RESULTS

To study the ability of our proposed method to constrain cosmology, we simulate mock luminosity distance measurements according to Eq. (14). We test two different detection thresholds to test the impact of the assumed D_L^{max} . By default we assume that all systems with $D_L^{\text{obs}} < 40$ Gpc are detectable, but for comparison, we also explore constraints with an observed distance limit of 20 Gpc. Given 10, 100, 1000, and 10 000 mock GW observations within the chosen distance limit, we calculate posteriors over cosmological parameters. All parameter inference is done with an input flat Λ CDM cosmology with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$. For extensions to Λ CDM, we use default values of $w_0 = -1$, $w_a = 0$, and $c_M = 0$. We assume the merger rate evolution is known perfectly as a function of redshift according to Eq. (12).

A. H_0 and Ω_M

We begin by assuming a flat Λ CDM universe and calculating 2D posteriors in H_0 and Ω_M given our simulated distance measurements. Figure 2 shows an example posterior from 10 000 GW events, given flat priors in Ω_M and H_0 . The 2D posterior is highly degenerate and unsurprisingly constrains H_0 much more strongly than Ω_M . By empirically fitting the degeneracy, we find that our method is most sensitive to the combination $H_0^{2.8} \Omega_M$, which

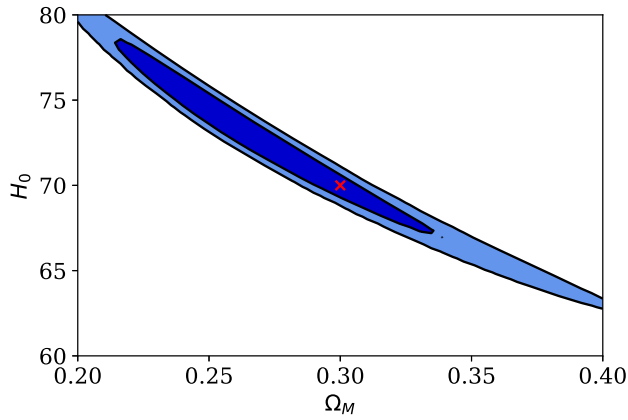


FIG. 2. 2D posterior distribution in H_0 and Ω_M for a flat Λ CDM model, inferred from 10 000 simulated luminosity distance measurements with $D_L^{\max} = 40$ Gpc. We assume flat priors on H_0 and Ω_M . 68% and 95% contours are shown. Input cosmology is marked with an X.

differs from the combination $H_0^2 \Omega_M$ best-measured by the CMB. This method, if used as a joint probe, can help break the degeneracy in H_0 and Ω_M in measurements by current or future CMB experiments.

We estimate the expected constraints in terms of $H_0^{2.8} \Omega_M$ for different sample sizes in Fig. 3. We find that the convergence of the 1σ (68% credibility) constraint in $H_0^{2.81} \Omega_M$ scales with the number of events N as $\frac{18\%}{N^{0.5}}$ for a distance limit of $D_L^{\max} = 40$ Gpc. For a distance limit of 20 Gpc, the expected precision is degraded to $\frac{50\%}{N^{0.5}}$. Much of the cosmology information appears to come from distances greater than 20 Gpc, as expected from Fig. 1. If H_0 is measured at subpercent levels from nearby BNS mergers with counterparts and the merger rate evolution is known, we expect to constrain Ω_M to the 1% level with a couple of hundred of observations (to be expected within a few weeks of observing with 3G detectors).

B. Dark energy parameters

Next we consider extensions to flat Λ CDM and their effect on the background cosmology. We use the $w_0 w_a$ parametrization of the equation of state with free parameters w_0 (the equation of state parameter at $z = 0$) and w_a (the evolution of the equation of state with scale factor $a = \frac{1}{1+z}$). While our method is sensitive to the dark energy equation of state, the resulting constraints on the dark energy parameters are largely degenerate with

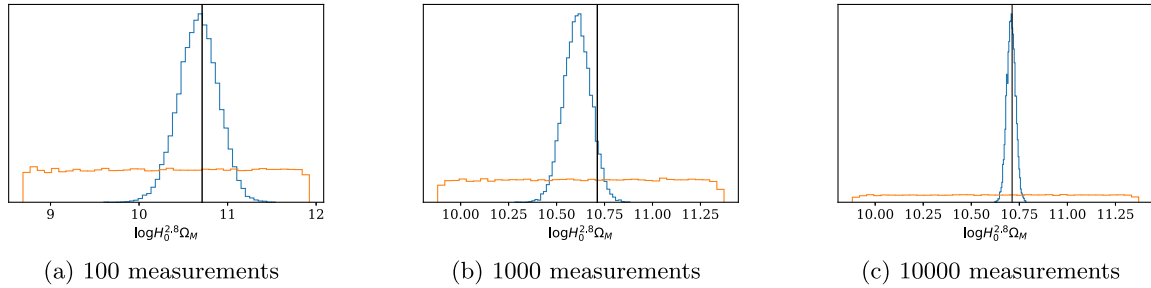


FIG. 3. Posterior (blue) and prior (orange) on $H_0^{2.81} \Omega_M$ inferred from simulated luminosity distances. We reweigh the flat prior on Ω_M and H_0 of Fig. 2 to a flat prior on the combination $H_0^{2.81} \Omega_M$.

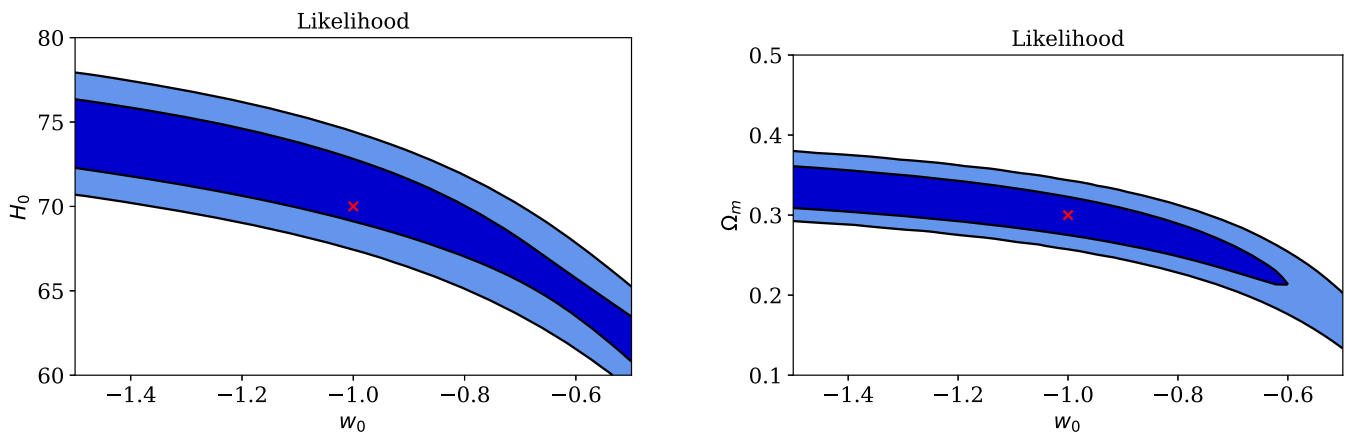


FIG. 4. 2D posterior distribution in H_0 and w_0 (left) and Ω_M and w_0 (right) in a flat w CDM cosmological model ($w_a = 0$), inferred from 1000 GW observations with $D_L^{\max} = 40$ Gpc. In each plot, the parameter not shown is fixed to its true value ($\Omega_M = 0.3$ on the left, $H_0 = 70$ km/s/Mpc on the right), and we take flat priors on the two free parameters. Contours show 68% and 95% credible regions.

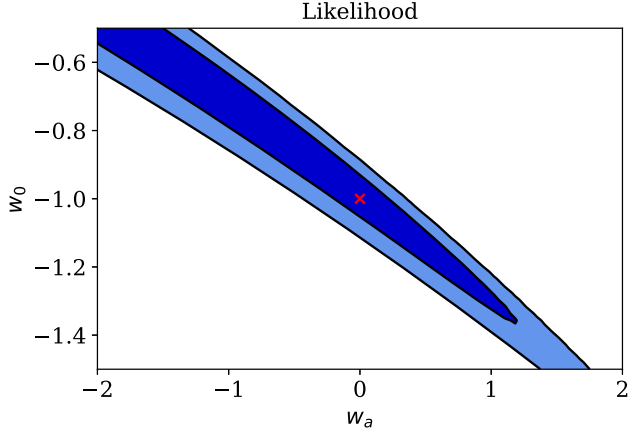


FIG. 5. 2D posterior distribution in w_0 and w_a for a $w_0 w_a$ CDM model, inferred from 10 000 simulated luminosity distance measurements with $D_L^{\max} = 40$ Gpc. H_0 and Ω_M are fixed to their true values, and we adopt flat priors on w_0 and w_a . Contours show 68% and 95% credible regions.

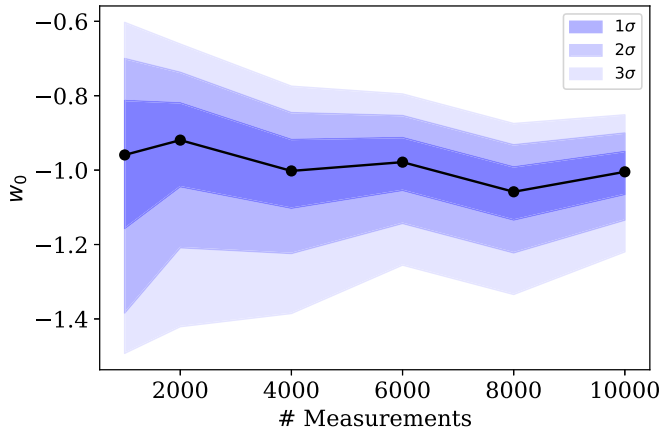
measurements of Ω_M and H_0 , which dominate the constraints, as seen in Fig. 4. Nevertheless, with external cosmological priors on H_0 and Ω_M , we can derive meaningful constraints on w_0 and w_a . Fixing $H_0 = 70$, $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$, we derive joint constraints on w_0 and w_a in Fig. 5. These two parameters are degenerate, such that a larger value of w_0 and a smaller w_a are consistent with the input cosmology. Fixing one parameter and constraining the other, the convergence of the 1σ constraint in w_0 scales as $\frac{500\%}{N^{0.5}}$ assuming a distance limit of 40 Gpc [see Fig. 6(a)], and also scales as \sqrt{N} for w_a in a fixed cosmology [Fig. 6(b)]. The width of the credible intervals in w_a are highly dependent on the maximum prior bound

considered for w_a , where positive w_a is constrained much more strongly. If we work with a w CDM model ($w_a = 0$) and adopt subpercent prior constraints on H_0 and Ω_M , we expect that 10 000 events can constrain the dark energy equation of state parameter w_0 to 5%, comparable to, but completely independent of, the available constraints from the combination of CMB, BAO, supernovae and weak lensing data [114].

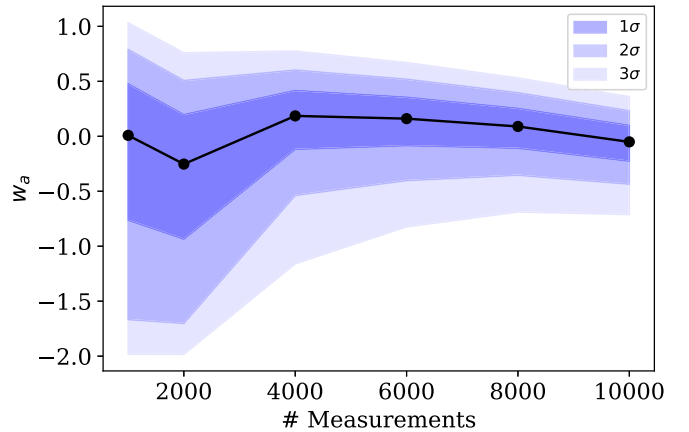
C. Modified gravity

Finally, we study extensions to Λ CDM through their effect on modified GW propagation. We assume the effect of the modified theory on the background expansion is minimal, so that the background expansion can be described by Λ CDM, and forecast constraints on c_M , the running of the Planck mass, based on the parametrization given in Eq. (4). Using standard sirens, it is possible to jointly constrain c_M , H_0 , and Ω_M , although the joint posterior has strong degeneracies (see Fig. 7). Jointly inferring H_0 , Ω_M , and c_M with broad priors, the 1σ width of the marginal posterior on c_M converges roughly as $\frac{60}{N^{0.5}}$.

Fixing all other parameters, including H_0 , the width of the 1σ constraint in c_M scales approximately as $\frac{3.4}{N^{0.5}}$, with N the number of events, as shown in Fig. 8. Current cosmological measurements constrain c_M to $\pm\mathcal{O}(1)$ [115], while BNS observations with counterparts in the advanced LIGO era can constrain c_M to ± 0.5 [81]. We find that if the merger redshift distribution is known and H_0 and Ω_M are perfectly measured, a hundred BNS observations within a distance limit of 40 Gpc can already surpass these projected limits. Without using external measurements on H_0 and Ω_M , it would take $\sim 10,000$ events to surpass these limits. We can interpret these constraints in terms of the



(a) Projected convergence of measurements of w_0 with fixed cosmology and $w_0 = -1$.



(b) Projected convergence of measurements of w_a with fixed cosmology and $w_a = 0$. Plot shows one realization of each measurement count. Constraints on w_a are strongly dependent on the value of w_a , and so statistical fluctuations are reflected in the credible interval width

FIG. 6. Projected convergence of dark energy equation of state measurements.

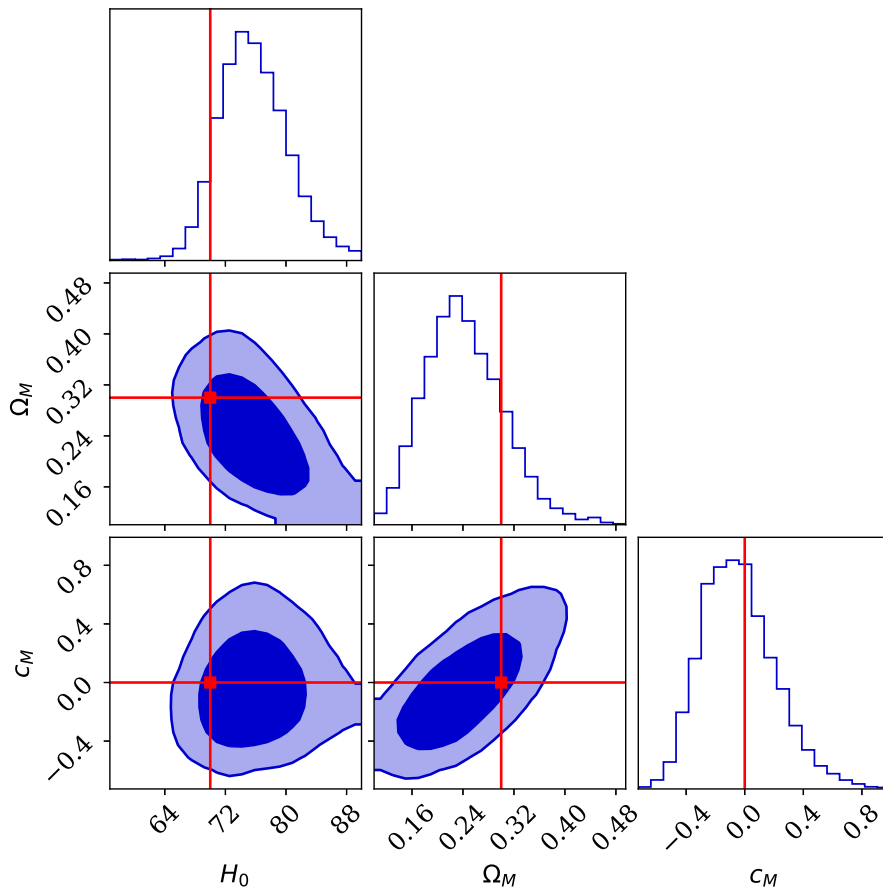
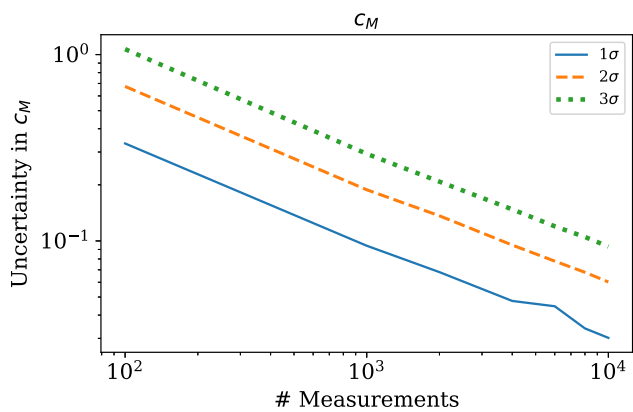


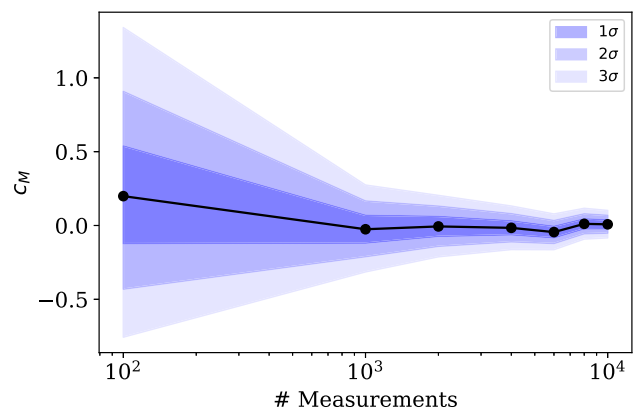
FIG. 7. Joint posterior distribution in H_0 , c_M and Ω_M and c_M in a flat Λ CDM background cosmology, inferred from 10 000 simulated distance measurements. We pick flat priors in all parameters. Contours show 68% and 95% credible regions.

value of the effective Planck mass or Newton’s constant at redshift z compared to today [95,103]. For $c_M = 0$ and fixed H_0 and Ω_M , the 1σ measurement in c_M from 10 000 GW events translates to an effective Planck mass of $2.172 \pm 0.017 \times 10^{-8}$ kg, or an effective Newton’s

constant of $6.70 \pm 0.11 \times 10^{-11} \frac{\text{N}\cdot\text{m}^2}{\text{kg}^2}$ at $z = 2$. Additionally, we can repeat the analysis using the modified GW propagation model proposed by Belgacem *et al.* [96], parametrized in terms of Ξ_0 and n . As an example, we fix $n = 1.91$, as predicted by the RT nonlocal gravity



(a) Uncertainty



(b) Credible intervals

FIG. 8. Projected convergence of measurements of the running of the Planck mass parameter c_M in a flat Λ CDM background cosmology. The parameters of the background cosmology are assumed to be perfectly measured.

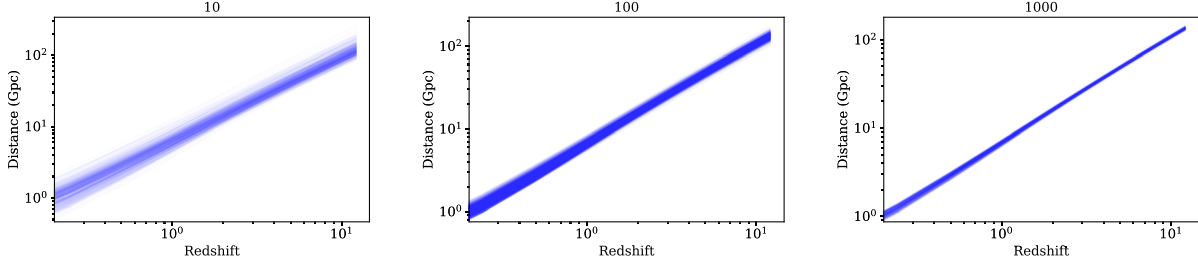


FIG. 9. Distance-redshift curves sampled from the joint posterior in H_0 and Ω_M , calculated with a flat prior in H_0 and Ω_M . We show the inference with 10, 100 and 1000 events within a distance limit of 40 Gpc. The degeneracy in H_0 and Ω_M captures a consistent distance-redshift relation.

model [25,116]. With all other cosmological parameters fixed, a simulated 10 000 events yields a measurement $\Xi_0 = 1.002 \pm 0.009$ ($\Xi_0 = 1$ for GR). These measurements at $z \sim 2$ could complement observations by the Laser Interferometer Space Antenna (LISA), which will probe modified GW propagation out to even higher redshifts ($z \lesssim 10$) by observing GWs from supermassive BBH mergers with possible EM counterparts [104].

D. Discussion

Comparing a catalog of GW luminosity distances against a known redshift distribution is ultimately sensitive to the underlying distance-redshift relation, as also pointed out by Ding *et al.* [39]. For the flat Λ CDM and w CDM models also considered by Ding *et al.* [39], we find similar results for the expected constraints on H_0 , Ω_M and w with 10 000 events (compare their Fig. 2 with our Fig. 2, for example). Regardless of the assumed cosmological model, which provides a parametrization for the distance-redshift relation, we can examine our parameter measurements from the previous subsections in terms of constraints on the luminosity distance-redshift relation or the expansion rate $\frac{H(z)}{1+z}$. Figure 9 shows posterior draws from the distance-redshift relation inferred in a flat Λ CDM model with flat priors on H_0 and Ω_M . Draws of H_0 and Ω_M within our posterior are such that the luminosity distance is the same for a given redshift, and so H_0 has a dominant effect.

Drawing H_0 and Ω_M from the joint posterior, we also look at the expected constraints on the $H(z)$ evolution as a

function of redshift, as in Fig. 10. The spread in $\frac{H(z)}{1+z}$ is smallest at redshifts $z \gtrsim 1$. In a $w_0 w_a$ CDM model, the joint posterior in w_0 and w_a with fixed H_0 and Ω_M yields the lowest spread in $\frac{H(z)}{1+z}$ (at a nonzero redshift) at around $z = 2$. This is consistent with our expectations that most of the cosmological information comes from knowledge of the redshift at which the merger rate peaks.

The forecasts described in this section depend on the true redshift distribution of GW sources, and how well it can be measured. Motivated by recent measurements that favor short delay times for BNS mergers [40–42], we have assumed that the BNS rate density peaks around $z = 2$ like the SFR. A recent analysis of *Fermi* and *Swift* short GRBs finds that their rate density peaks between $z \sim 1.5$ – 2 [41]. While current constraints on the BNS merger rate evolution are broad, as discussed in Sec. I, we expect the measurements to improve significantly over the next decade with upcoming observations of GRBs, kilonovae, and BNS host galaxies. Because we expect to best constrain the cosmological expansion rate near the peak redshift, if it turns out that time delays are long and the peak is at $z < 2$, our projected constraints will differ. Crucially, if the wrong redshift evolution is assumed, the resulting cosmological inference will be biased, as explicitly demonstrated in Ding *et al.* [39]. We therefore expect that the redshift evolution will be inferred jointly with the cosmological parameters, so that its uncertainty can be marginalized over.

Additionally, most of our forecasts have assumed that all BNS mergers within an observed distance of 40 Gpc can be

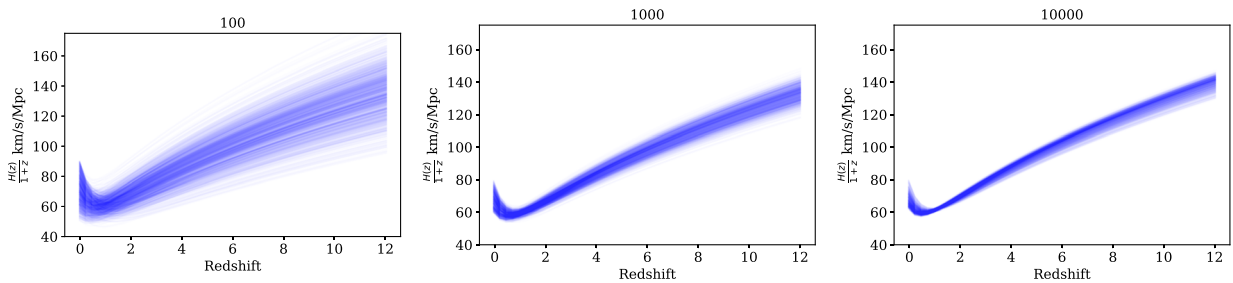


FIG. 10. $H(z)$ evolution sampled from the joint posterior in H_0 and Ω_M in a flat Λ CDM model, calculated with a flat prior in H_0 and Ω_M . We show the inference with 100, 1000 and 10,000 events within 40 Gpc. $\frac{H(z)}{1+z}$ is best constrained around $z \sim 1$ – 2 .

detected, and we have shown that we expect worse constraints, typically by a factor of a few, if the observed distance limit is lowered to 20 Gpc. The sensitivities of the proposed next-generation GW detectors are not yet finalized, and we expect this to affect the projections here, modifying the number of events needed to reach the desired accuracy in the cosmological parameters. Finally, we have considered the case in which the merger rate density $\mathcal{R}(z)$ is directly measured, rather than $p(z)$. Because of the cosmological dependence of the comoving volume element, if $\mathcal{R}(z)$ is perfectly known, there is cosmological information in $p(z)$. This effect is subdominant to the distance-redshift relation probed by the GW luminosity-distance relation, and only affects Ω_M and to a lesser extent w_0 and w_a . We expect our results to differ slightly in the case that $p(z)$ is more directly available.

Standard sirens are an independent probe to address the tension in H_0 measurements between so-called “early-universe” and “late-universe” estimates. While with a flat prior, H_0 and Ω_M are strongly degenerate, a precise measurement of H_0 is possible with our method using an outside prior on Ω_M , such as from measurements of the CMB, galaxy clustering, or weak lensing. Given that the joint posterior in H_0 and Ω_M is captured by $H_0^{2.8}\Omega_M$, when used with experiments sensitive to a different combination of H_0 and Ω_M , our method can help break this degeneracy. Standard sirens are also uniquely poised to probe the nature of dark energy, not only through its effect on the background expansion parametrized by the dark energy equation of state w , but primarily on its effect on GW propagation, parametrized by c_M here. To constrain the dark energy parameters w_a and w_0 , or the running of the Planck mass in modified gravity c_M , outside priors on both H_0 and Ω_M are necessary to reveal the subdominant effects on the GW distance distribution.

IV. CONCLUSION

GW standard sirens can independently test the Λ CDM cosmological model and provide insight into the mysterious dark sector, namely dark matter and dark energy. The next generation of GW detectors, the proposed Einstein Telescope and Cosmic Explorer, would revolutionize standard siren science by observing the GW universe out to tens of Gpc. The challenge for GW cosmology will be to measure the redshifts of these mergers, especially considering the difficulties of identifying EM counterparts and potential host galaxies at $z \gg 1$.

Previous work [39] showed that, in the absence of targeted EM followup campaigns or complete galaxy catalogs, prior knowledge of the *distribution* of merger redshifts can be compared against GW luminosity distances

to infer cosmological parameters. In this work we argue that we can leverage external measurements of the evolution of the BNS merger rate, which, in particular, is expected to peak at some redshift. This provides a redshift feature which can be used in a standard siren analysis to constrain cosmology and modified gravity. As a demonstration of this method, we used a simple toy model in which the evolution of the BNS merger rate as a function of redshift is known perfectly, and studied how the observed GW luminosity distance distribution alone can measure parameters of the w_0w_a CDM model and the running of the Planck mass. This allows us to isolate the available information in a catalog of GW distances, compared to the additional information that enters from the mass distribution.

In reality, we expect this method to be used jointly with fits to the mass distribution and/or available galaxy information. The information from the mass distribution will likely dominate the inference if there is a sharp, redshift-independent mass feature like a NS-BH mass gap at low masses [32] or a pair-instability mass gap at high masses [34]. Because the GW luminosity distance distribution inherently carries information about cosmology, even if it is not used as the primary observable to measure cosmology, it must be taken into account in all standard siren analyses at high redshifts to avoid biasing the cosmological constraints [23,25,117,118].

We have focused on the next generation of detectors in our analysis because they will likely observe GW mergers past cosmic noon, or the peak redshift of the merger rate, providing a clear feature whose feature can be measured in both redshift and distance space. Similar analyses can in principle be carried out on existing GW catalogs; in combination with measurements of the stochastic GW background, current GW observatories will constrain the peak of the BBH merger rate distribution [119]. However, currently the distance distribution is only meaningfully constrained for high-mass BBH mergers, while the corresponding redshift distribution is not well-constrained from EM observations. Existing BBH observations can only constrain large deviations from GR; for example, GW leakage in large extra dimensions [120,121].

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