Errata

Weinberg's Perturbation Theory at Infinite Momentum, IRWIN MANNING [Phys. Rev. **184**, 1957 (1969)]. Mention should be made of the work by S. Chang and S. Ma [Phys. Rev. **180**, 1506 (1969)] which appeared while the above Comment was in press. These authors also consider the effect of spin on Weinberg's rules.

Sugawara Model, Stress Tensor, and Spectral Sum Rules, K. T. MAHANTHAPPA AND D. R. PALMER [Phys. Rev. **185**, 1970 (1969)]. (1) In Eqs. (4.7) and (4.11b), instead of "$/(1-\frac{3}{4}N)$" read "$(1-\frac{3}{4}N)^{-1}$." (2) Divide the first term on the right-hand side of Eq. (4.8) by $(1-\frac{3}{4}N)$. (3) In the first line after Eq. (5.1), instead of "(4.7)" read "(4.4)."

The treatment given in this paper is valid for the case of $f(\xi)\sim\xi^n$ with $n=0$. The case of $n>0$ is beset with difficulties the resolution of which is not clear in the framework of the Sugawara current commutators which are the same as those of gauge-field algebra. The sum rules analogous to Eqs. (4.11) can be easily written down for the case of $n=0$ using Eqs. (4.6), (4.9), and (4.10). The conclusion concerning the cross section for $e^+e^-\to$ hadrons remains unaltered. We thank Dr. K. Barnes of Queen Mary College (London) for correspondence in this connection.

Presymmetry. II, H. EKSTEIN [Phys. Rev. **184**, 1315 (1969)]. (1) On p. 1331, the assumption (5) (lines 15 and 16 from top) should read: Each subalgebra $\mathcal{O}_{tce_l}{}^{(n)}$ is causally independent of every algebra $(\bigcup_{m\neq l}\mathcal{O}_{tce_m}{}^{(n)})''$ that is generated by subalgebras $\mathcal{O}_{tce_m}$ associated to other particles $(m\neq l)$.

(2) On p. 1331, Theorem 3 should read: *Theorem 3:* Let $\mathfrak{A}_1\cdots\mathfrak{A}_n\subset\mathfrak{B}$ be causally independent factor subalgebras, algebraically isomorphic to $\mathfrak{B}$, in the algebra $\mathfrak{B}$ of all bounded operators on $\mathcal{H}$, such that the union of these subalgebras generates $\mathfrak{B}$, and let each subalgebra $\mathfrak{A}_i$ be causally independent of every algebra $(\bigcup_{j\neq i}\mathfrak{A}_j)''$ generated by other subalgebras $\mathfrak{A}_{j\neq i}$. Then there exists a tensor decomposition of $\mathcal{H}$, namely,

$$\mathcal{H}=\mathcal{H}_1\otimes\mathcal{H}_2\mathcal{H}\cdots\otimes\mathcal{H}_n\,,$$

and a corresponding tensor product for $\mathfrak{B}$, namely,

$$\mathfrak{B}=\mathfrak{a}_1\otimes\mathfrak{a}_2\otimes\cdots\otimes\mathfrak{a}_n\,,$$

where $\mathfrak{a}_n$ is the irreducible faithful representation of $\mathfrak{A}_n$.

(3) On p. 1335, the Theorem should read: *Theorem:* Let $\mathfrak{A}_i$ $(i=1,\ \ldots,\ n)$ be discrete von Neumann factor subalgebras of the algebra $\mathfrak{B}(\mathcal{H})$ of all bounded linear operators on a Hilbert space $\mathcal{H}$, such that the union $\bigcup_{i=1}{}^n\mathfrak{A}_i$ generates $\mathfrak{B}(\mathcal{H})$, and let each subalgebra $\mathfrak{A}_i$ be causally independent of each algebra $(\bigcup_{j\neq i}\mathfrak{A}_j)''$ generated by the union of other subalgebras $\mathfrak{A}_{j\neq i}$. Let $\mathfrak{a}_i$ be the irreducible representation [isomorphic to $\mathfrak{B}(\mathcal{H})$] of $\mathfrak{A}_i$. Then $\mathfrak{B}$ is the tensor product of the algebras $\mathfrak{a}_i$, i.e.,

$$\mathfrak{B}=\mathfrak{a}_1\otimes\mathfrak{a}_2\otimes\cdots\otimes\mathfrak{a}_n\,.$$

These stronger assumptions are intuitively consequences of the weaker statements in the text, but they do not follow formally. The stronger version is used in the proof.

Lorentz-Pole Structure and Duality of Some Crossing-Symmetric Amplitudes with Regge Behavior, KHALIL M. BITAR [Phys. Rev. **185**, 2032 (1969)]. This paper was inadvertently printed in the "Comments and Addenda" section. It should have been a regular Article.

ANNOUNCEMENT

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