

Update on energies and widths in  $^{13}\text{Be}$ 

H. T. Fortune

*Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*

(Received 26 November 2018; published 7 January 2019)

I have compared new experimental information on resonances in  $^{13}\text{Be}$  with results of theoretical calculations. For the  $1/2^+$  resonance at 0.86 MeV, the reported width of 1.70(15) MeV is considerably larger than the single-particle limit of 1.3 MeV. For the first  $5/2^+$ , the calculated width and  $2^+$  branching ratio for neutron decay are both in rough agreement with the data. I discuss the possibility that events from decay of the second  $5/2^+$  to  $^{12}\text{Be}(2^+)$  could contribute to the 0.86-MeV peak. If the 4.0-MeV resonance is indeed  $3/2^+$ , then its width should be considerably larger than reported.

DOI: [10.1103/PhysRevC.99.014304](https://doi.org/10.1103/PhysRevC.99.014304)

## I. INTRODUCTION

A recent experiment [1] has greatly improved our understanding of resonances in  $^{13}\text{Be}$ , which has no bound states. Three other relatively recent experiments had served to both clarify and confuse the issue [2–4]. Ribeiro *et al.* [1] used proton knockout from a 400 MeV/nucleon  $^{14}\text{B}$  beam incident on a  $\text{CH}_2$  target to produce  $^{13}\text{Be}$ , and they detected  $^{12}\text{Be} + n$  in coincidence. The experiment also had the ability to detect coincident  $^{12}\text{Be} \gamma$ s. They report a  $1/2^+$  resonance at an energy of 0.86(4) MeV with a width of 1.70(15) MeV and a  $5/2^+$  resonance at an energy of 2.11(5) MeV. For the latter, they took the width of 0.4 MeV from earlier heavy-ion-induced transfer experiments [5,6]—the reaction  $^{13}\text{C}(^{14}\text{C}, ^{14}\text{O}) ^{13}\text{Be}$  at  $E_{\text{Lab}} = 337$  MeV [5] and the  $^{14}\text{C}(^{11}\text{B}, ^{12}\text{N}) ^{13}\text{Be}$  reaction at  $E_{\text{lab}} = 190$  MeV [6]. Ribeiro *et al.* [1] observed low-energy neutrons in coincidence with  $\gamma$ s of energy  $\sim 2$  MeV, which they interpreted as evidence that the  $5/2^+$  resonance also decayed to the  $2^+$  of  $^{12}\text{Be}$ , in addition to the ground state (g.s.). Their branching ratio (BR) was  $2^+/\text{g.s.} = 0.1/(0.24) = 0.42$  [1,7].

Reference [1] adopted a low-lying  $1/2^-$  resonance near 0.5 MeV from Ref. [8]. They argue correctly that a negative-parity state should not be produced in proton removal from  $^{14}\text{B}$ , but that it should be populated in neutron removal from  $^{14}\text{Be}$ —which was the procedure used in Ref. [2,8]. Ribeiro *et al.* also included in their analysis a second  $5/2^+$  and first  $3/2^+$  resonances, with both  $J^\pi$  assignments tentative. They took both energies (2.92 and 4.0 MeV, respectively) and widths (both 0.4 MeV) from the heavy-ion work [5,6]. Both are weak. The resolution width (FWHM) in Ref. [1] was about 0.7 MeV at  $E = 2$  MeV, and it increases as  $E^{0.75}$ . Thus, all the widths, except for  $1/2^+$ , are significantly less than the resolution width, and that is the justification for the use of earlier widths.

Here, I examine the new experimental evidence in comparison with model calculations.

## II. CALCULATIONS AND RESULTS

I have calculated single-particle (sp) widths in a potential model, using a Woods-Saxon shape with geometrical parameters  $r_0$ ,  $a = 1.26$ , 0.60 fm. Well depth was adjusted to reproduce resonance energy, and the width was computed from the phase shift. For  $\ell = 2$ , the width calculation is straightforward. The absence of a barrier for an  $s$ -wave neutron resonance is somewhat of a complication, but I have used the relationship  $\Gamma_{\text{sp}}(\ell = 0) = (2E)^{1/2}$ , where both energy and width are in MeV. From these  $\ell = 0$  and 2 sp widths, I have computed expected widths with the expression  $\Gamma_{\text{calc}} = S \Gamma_{\text{sp}}$ , where  $S$  is the relevant  $^{12}\text{Be} + n$  spectroscopic factor, given previously [9]. Relevant information is displayed in Table I.

The  $1/2^+$  resonance could have a spectroscopic factor near unity, but even if so, the reported width of 1.70(15) MeV is significantly larger than the sp width of 1.3 MeV. I return to this discrepancy below. Even though the first theoretical  $5/2^+$  state has a (sd) $^3$  component that is larger than the  $1d_{5/2}$  component [10], its spectroscopic factor to  $^{12}\text{Be}$  (g.s.) is quite large—0.94 in my calculations. This happens because of the large (sd) $^3$  component in  $^{12}\text{Be}$  (g.s.). With my  $S$  and  $\Gamma_{\text{sp}}$ , the computed width for decay of the first  $5/2^+$  state to the g.s. is 0.63 MeV, slightly larger than the old experimental width of 0.4 MeV. However, it appears that the newer data [1–4] could easily accommodate a larger width. For decay of this state to the  $2^+$  of  $^{12}\text{Be}$ , Ref. [1] quotes a neutron energy of 0.1 MeV. With this energy, my computed width is 0.13 MeV (Table I) for  $s$ -wave decay (at this low energy,  $d$ -wave decay is weak enough to ignore), giving an expected  $\text{BR}(2^+/\text{g.s.}) = 0.13/0.63 = 0.21$ —to be compared with the experimental BR quoted above of 0.42. The authors quote a relative yield of 0.1 (no uncertainty given) for decay through the  $2^+$  and 0.24(4) for g.s. decay. (Just in passing, I note that a neutron energy of 0.4 MeV for decay to the  $2^+$  would provide exact agreement with the experimental BR.)

TABLE I. Energies, decay modes, and widths in  $^{13}\text{Be}$  (energies and widths in MeV).

$^{13}\text{Be}$ state	Decay	$E_n$	$S^c$	$\Gamma_{\text{sp}}$	$\Gamma_{\text{calc}}^d$	$\Gamma_{\text{exp}}$
$1/2^+$	to g.s.	$0.86^a$	$\sim 1$	1.3	1.3	$1.70(15)^a$
$5/2_1^+$	to g.s.	$2.11^a$	0.94	0.67	0.63	$0.4^b$
	to $2^+ + s$	$(0.1)^a$	0.29	(0.45)	(0.13)	
$5/2_2^+$	to g.s.	$2.92^b$	0.0004	1.4	0.0006	$0.4^b$
	to $2^+ + s$	0.8	0.15	1.25	0.19	
	to $2^+ + d$	0.8	0.005	0.082	0.0004	
	To exc. $0^+$	0.68	0.85	0.066	0.056	
$3/2^+$	to g.s.	$4.0^b$	$\sim 0$	2.82	small	$0.4^b$
	to $2^+ + s$	1.9	0.19	1.9	0.37	
	to $2^+ + d$	1.9	1.32	0.52	0.69	

<sup>a</sup>Reference [1].<sup>b</sup>References [5,6].<sup>c</sup>Reference [9].<sup>d</sup> $\Gamma_{\text{calc}} = S\Gamma_{\text{sp}}$ .

The second  $5/2^+$  state is predicted to have an extremely small decay branch to  $^{12}\text{Be}(\text{g.s.})$ , with the largest decay to the first excited  $0^+$  state of  $^{12}\text{Be}$  [11]. The experiment of Ref. [1] was not sensitive to this excited  $0^+$  decay, but they appear to have observed some g.s. decays. Other than the excited  $0^+$  decay, the other important branch should be  $s$ -wave decay to the  $2^+$ , for which the computed width is 0.19 MeV, considerably smaller than the supposed experimental width of 0.4 MeV. Reference [1] did not report observation of this decay, but I note that such a decay would have a neutron energy near 0.8 MeV. The presence of such decays might account for the fact that the reported width for the 0.86-MeV resonance is significantly larger than the sp limit. If this second  $5/2^+$  state does indeed also decay to the excited  $0^+$ , that would add about 0.06 MeV to the computed width.

A decrease in  $S$  for the first  $5/2^+$  would require an increase in  $S$  for the second  $5/2^+$ . Such changes would move both calculated widths closer to the experimental values.

The  $3/2^+$  resonance is predicted to have a very small g.s. branch, but reasonably strong decays to the  $2^+$ , with both

$s$  and  $d$ . Reference [1] observed 2-MeV  $\gamma$ s in coincidence with 2-MeV neutrons, indicating decay of the 4-MeV resonance to  $^{12}\text{Be}(2^+)$ . With my predictions, the width of this resonance should be considerably larger than 0.4 MeV.

### III. SUMMARY

I have compared new experimental results for resonances in  $^{13}\text{Be}$  to previous and new model calculations. The reported width of 1.70(15) MeV for the  $1/2^+$  resonance is considerably larger than the sp limit of 1.3 MeV, perhaps implying another contribution to that peak in the energy spectrum—for which one possibility is decay of the second  $5/2^+$  state to the  $2^+$  of  $^{12}\text{Be}$ . The calculated width for the first  $5/2^+$  resonance is in reasonable agreement with (but slightly larger than) the experimental value. Reference [1] was the first to positively identify decays of this state to  $^{12}\text{Be}(2^+)$ . Their  $2^+/\text{g.s.}$  branching ratio is in rough agreement with my calculations. If the 4.0-MeV resonance is indeed  $3/2^+$ , it should be considerably wider than currently thought.

A recent review [12] identified a few unanswered questions in  $^{13}\text{Be}$ . One of them was: Is the lowest resonance near 0.5 MeV  $1/2^+$  or  $1/2^-$ , or are the two unresolved? Reference [1] was unable to answer this question. They state, “To promote one of them as the ground state  $^{13}\text{Be}$  is not within the scope of the present paper but certainly a challenge for theory.”

Another unanswered question [12] was: Can better evidence be found for decays of  $^{13}\text{Be}$  resonances to excited states of  $^{12}\text{Be}$ ? Reference [1] has provided convincing evidence for this question as it relates to the first  $5/2^+$  state. Another question dealt with events near 1 MeV, and the extent to which they correspond to g.s. decays versus decays of an excited state to an excited state. I referred to this question above, in relation to the width of the  $1/2^+$  resonance and the question of decays of the second  $5/2^+$  resonance to  $^{12}\text{Be}(2^+)$ .

Perhaps the most important unanswered question concerns the possibility of decays of the second  $5/2^+$  state to the excited  $0^+$  state. Reference [1] states that this “is indeed an experimental challenge.”

[1] G. Ribeiro *et al.*, *Phys. Rev. C* **98**, 024603 (2018).  
 [2] Y. Aksyutina *et al.*, *Phys. Rev. C* **87**, 064316 (2013).  
 [3] G. Randisi *et al.*, *Phys. Rev. C* **89**, 034320 (2014).  
 [4] B. R. Marks *et al.*, *Phys. Rev. C* **92**, 054320 (2015).  
 [5] A. N. Ostrowski *et al.*, *Z. Phys. A* **343**, 489 (1992).  
 [6] A. V. Belozyorov *et al.*, *Nucl. Phys. A* **636**, 419 (1998).

[7] E. Nacher (private communication, 2018).  
 [8] Y. Kondo *et al.*, *Phys. Lett. B* **690**, 245 (2010).  
 [9] H. T. Fortune, *Phys. Rev. C* **93**, 054327 (2016).  
 [10] H. T. Fortune, *Phys. Rev. C* **90**, 064305 (2014).  
 [11] H. T. Fortune and R. Sherr, *Phys. Rev. C* **82**, 064302 (2010).  
 [12] H. T. Fortune, *Eur. Phys. J. A* **54**, 51 (2018).