

Precise measurement of α_K for the 88.2-keV $M4$ transition in ^{127}Te : Test of internal-conversion theory

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We have measured the K -shell internal conversion coefficient, α_K , for the 88.2-keV $M4$ transition from the 106-day isomer to the ground state in ^{127}Te to be 484(6). When compared with Dirac-Fock calculations of α_K , this result agrees well with the version of the theory that incorporates the effect of the K -shell atomic vacancy and disagrees with the one that does not. As a byproduct of this measurement, we have determined the β branching from the isomer to be 2.14(3)%.

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I. INTRODUCTION

Starting in 2004, we have been publishing a sequence of papers [1–8], in which we report measurements of K -shell internal conversion coefficients (ICCs) for $E3$ and $M4$ transitions with a precision of $\pm 2\%$ or better. The motivation has been to test ICC theory, in particular its treatment of the K -shell vacancy left behind by the emitted electron. When we began, the then-current survey of world data on ICCs [9] included only six α_K values for high-multipolarity transitions ($E3$, $M3$, and above) measured to two-percent precision or better. At the time, the surveyors concluded that the data favored the ICC calculation that ignored the atomic vacancy. Combining our work since then with the result reported here, we have now nearly doubled the number of precisely measured α_K values and have changed that conclusion.

What makes such precise measurements possible for us is our having a high-purity germanium (HPGe) detector whose relative efficiency is known to $\pm 0.15\%$ ($\pm 0.20\%$ absolute) over a wide range of energies; see, for example, Ref. [10]. By detecting both the K x rays and the γ rays from a transition of interest in the same well-calibrated detector at the same time, we can avoid many sources of error.

By 2008, our early results from this program influenced a reevaluation of ICCs by Kibédi *et al.* [11], who also developed BrIcc, a new database obtained from the basic code by Band *et al.* [12]. In conformity with our conclusions, BrIcc employed a version of the code that incorporates the vacancy in the “frozen orbital” approximation. The BrIcc database has been adopted by the National Nuclear Data Center (NNDC) and is available online for the determination of ICCs. Our experimental results obtained since 2008 continue to support

that decision and now include transitions in seven nuclei that cover the range $48 \leq Z \leq 78$.

The measurement we report here is of the 88.23-keV, $M4$ transition from the 106-day isomeric state in ^{127}Te to its ground state. The calculation of α_K for this transition depends appreciably on whether the atomic vacancy is accounted for, there being a 3.7% difference between the values obtained with and without inclusion of the vacancy. The only previous reported measurement of the α_K value [13] had an uncertainty of $\pm 5\%$, which overlaps both calculated values. Our new result has $\pm 1.2\%$ precision and clearly distinguishes between the two options.

II. MEASUREMENT OVERVIEW

We have described our measurement techniques in detail in previous publications [1,3] so only a summary will be given here. If a decay scheme is dominated by a single transition that can convert in the atomic K shell, and a spectrum of K x rays and γ rays is recorded for its decay, then the K -shell internal conversion coefficient for that transition is given by

$$\alpha_K \omega_K = \frac{N_K \epsilon_\gamma}{N_\gamma \epsilon_K}, \quad (1)$$

where ω_K is the fluorescence yield; N_K and N_γ are the total numbers of observed K x rays and γ rays, respectively; and ϵ_K and ϵ_γ are the corresponding photopeak detection efficiencies.

The fluorescence yield for tellurium has been measured several times, with a weighted average quoted to $\pm 3.1\%$ [14]. Furthermore, world data for fluorescence yields have also been evaluated systematically as a function of Z [15] for all elements with $10 \leq Z \leq 100$, and ω_K values have been recommended for each element in this range. The recommended value for tellurium, $Z = 52$, is 0.875(4), which is consistent with the average measured value but has a smaller relative uncertainty, $\pm 0.5\%$. We use this value.

The decay scheme of the 106-day isomer in ^{127}Te is shown in Fig. 1. With a single electromagnetic decay path, directly feeding the ground state, it clearly satisfies the condition

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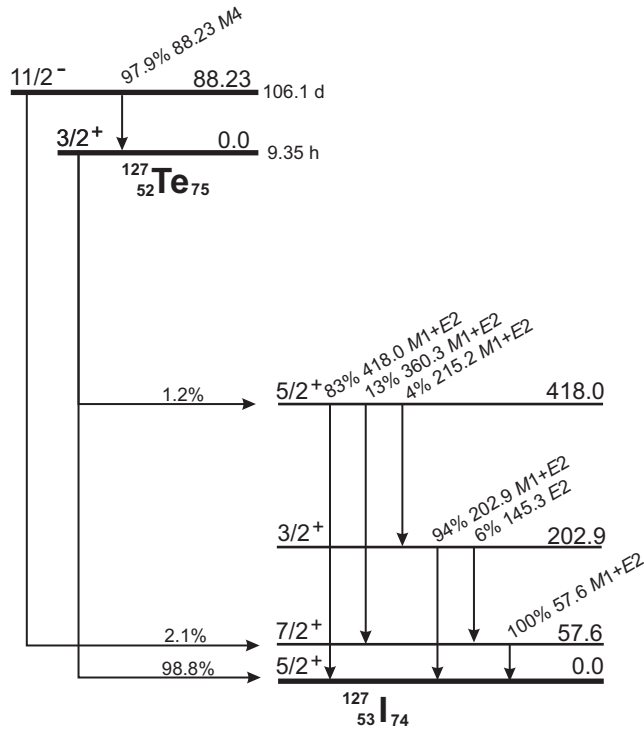


FIG. 1. Decay scheme for the 106-day isomer in ^{127}Te , illustrating the channels important to this measurement. The data are taken from Ref. [16] except for the β branching ratio from the isomeric state and the γ branching of the 418.0 keV level, which reflect new results from this work.

required for the validity of Eq. (1). The only complication is that the isomer also has a 2.1% β branch to an excited state in ^{127}I , which decays by a 57.6-keV transition with an α_K value of 3.16 [16]. This leads unavoidably to the presence of iodine K x rays. Though their intensity is only a few percent that of the predominant tellurium x rays, the two groups are unresolved from one another in our detector, so the iodine component must be carefully accounted for. The ^{127}Te ground state β decays as well, but only weakly populates an excited state, the decay of which produces considerably less conversion.

In our experiment, the HPGe detector we used to observe both γ rays and K x rays has been meticulously calibrated [10,17,18] for efficiency to sub-percent precision, originally over an energy range from 50 to 3500 keV but more recently extended [6] with $\pm 1\%$ precision down to 22.6 keV, the average energy of silver K x rays. Over this whole energy region, precise measured data were combined with Monte Carlo calculations from the CYLTRAN code [19] to yield a very precise and accurate detector efficiency curve. In our present study, the γ ray of interest at 88.2 keV is well within the energy region for which our efficiencies are known to a relative precision of $\pm 0.15\%$. The tellurium K x rays lie between 27 and 32 keV, comfortably within our extended region of calibration, so the detector efficiency for them can be quoted to a relative precision of $\pm 1\%$.

III. EXPERIMENT

A. Source preparation

We obtained tellurium powder enriched to $\geq 98\%$ in ^{126}Te from Trace Sciences International Corp. As delivered, the grain size was $25 \pm 5 \mu\text{m}$, but by grinding it between glass slides we reduced it to a relatively uniform grain size of about $1 \mu\text{m}$ as determined with a microscope. A small quantity of this powder was sprinkled over a 1-cm diameter disk of adhesive Mylar, $25 \mu\text{m}$ thick, and another layer of the same adhesive Mylar was placed on top, sealing the powder between. A second “dummy” sample, identical except for the absence of any tellurium, was prepared too.

The tellurium and dummy samples were activated together for 24 h in a neutron flux of $\sim 7.5 \times 10^{12} \text{ n}/(\text{cm}^2 \text{ s})$ at the TRIGA reactor. The tellurium and dummy samples were activated together for 24 h in a neutron flux of $\sim 7.5 \times 10^{12} \text{ n}/(\text{cm}^2 \text{ s})$ at the TRIGA reactor in the Texas A&M Nuclear Science Center. After removal from the reactor, both samples were found to have curled up. As we carefully unrolled it, the tellurium sample cracked, requiring us to secure it between two fresh Mylar sheets, $12.5 \mu\text{m}$ thick. The final source thus comprised a $4 \pm 2\text{-}\mu\text{m}$ -thick layer of tellurium enclosed between layers of Mylar $37.5 \mu\text{m}$ thick. We then stored the samples for one month to allow short-lived activities to die out.

B. Radioactive decay measurements

We acquired spectra with our precisely calibrated HPGe detector and with the same electronics used in its calibration [10]. Our analog-to-digital converter was an Ortec TRUMPTM 8k/2k card controlled by MAESTROTM software. We acquired 8k-channel spectra at a source-to-detector distance of 151 mm, the distance at which our calibration is well established. Each spectrum covered the energy interval 10–2000 keV with a dispersion of about 0.25 keV/channel.

After energy-calibrating our system with a ^{152}Eu source, we recorded sequential decay spectra from the tellurium sample in two one-week periods. Before, after and between these two periods, we interspersed measurements of the dummy sample and of room background. These were used to aid us in identifying the origin of all peaks in the spectrum.

IV. ANALYSIS

A. Peak fitting

We summed the spectra recorded from the tellurium sample, and summed the background spectra. The latter sum was then normalized to the same live time as the former and was subtracted from it. A portion of the background-subtracted spectrum recorded from the tellurium source is presented in Fig. 2: It includes the x- and γ -ray peaks of interest from the decay of ^{127m}Te , as well as a number of peaks from contaminant activities. In our analysis of the data, we followed the same methodology as we did with previous source measurements [1–8]. We first extracted areas, not only for the ^{127m}Te peaks, but also for essentially all the other x- and γ -ray peaks in the spectrum. Our procedure was to determine the areas with GF3,

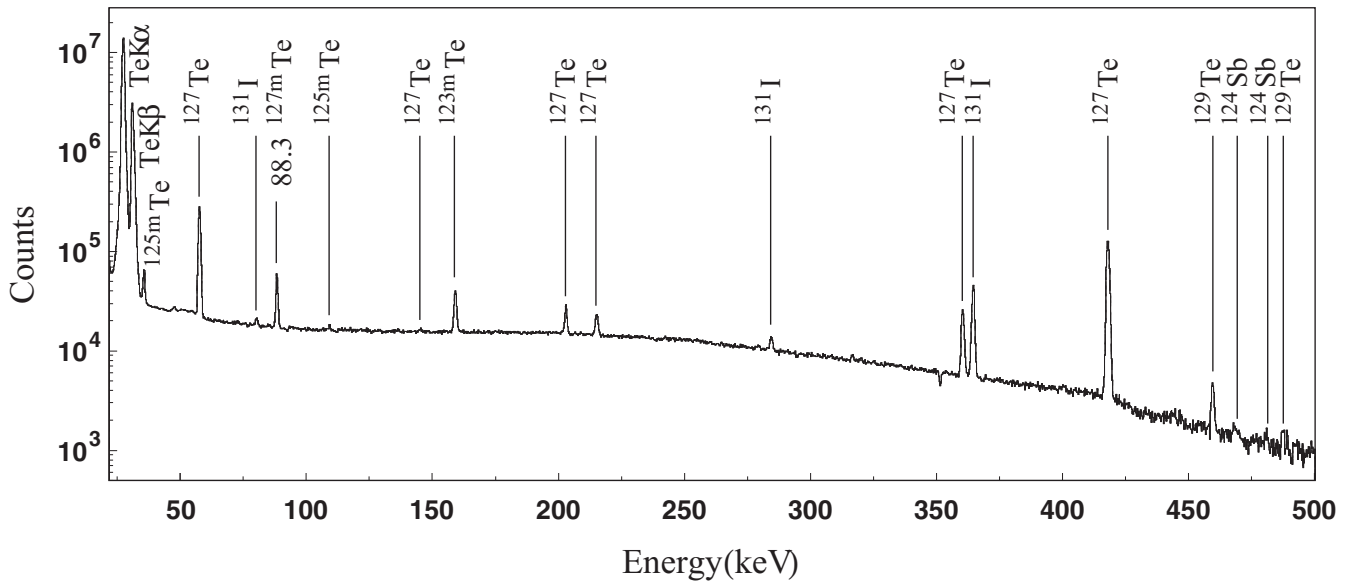


FIG. 2. Portion of the background-subtracted x- and γ -ray energy spectrum recorded over a period of two weeks, more than a month after activation of enriched ^{126}Te . Peaks are labeled by their β -decay parent. The inverted “peak” at 352 keV results from a subtracted background γ ray present in the decay of ^{214}Pb , which is a granddaughter of environmental ^{222}Rn and consequently is not constant in time.

the least-squares peak-fitting program in the RADWARE series [20]. In doing so, we used the same fitting procedures as were used in the original detector-efficiency calibration [10,17,18].

Figure 3 shows expanded versions of the two energy regions of interest for this measurement: one encompassing

the tellurium K x rays and the other, the γ ray at 88.2 keV. The tellurium x-ray peaks, which of course include some unresolved impurity x rays, lie cleanly on a low flat background; the γ ray is much weaker relative to the local background but is nevertheless easily analyzed.

In what follows we do not distinguish between K_α and K_β x rays. Scattering effects are quite pronounced at these energies and they are difficult to account for with an HPGe detector when peaks are close together, so we have chosen as before to use only the sum of the K_α and K_β x-ray peaks. For calibration purposes, we consider each sum to be located at the intensity-weighted average energy of the component peaks¹—28.03 keV for tellurium and 29.20 keV for iodine.

The total number of counts recorded in the combined x-ray peaks and in the 88.2-keV γ -ray peak must next be corrected for impurities and for other effects. These corrections are described in the following sections.

B. Impurities

Once the areas (and energies) of all γ -ray peaks had been established, we could identify every impurity in the ^{127m}Te spectrum and carefully check to see if any were known to produce x or γ rays that might interfere with the tellurium K x rays or the γ -ray peak of interest at 88.2 keV. As is evident from Fig. 2, even the weakest peaks were identified. In all, we found 10 contaminant activities that contribute in some way to the tellurium x-ray region; these are listed in Table I, where the contributions are given as percentages of the total x rays recorded. Only ^{125m}Te contributes more than a percent; most of the rest are extremely weak. As to the γ -ray peak, only one impurity, ^{123m}Te , interferes with it in any way and the impact is very slight.

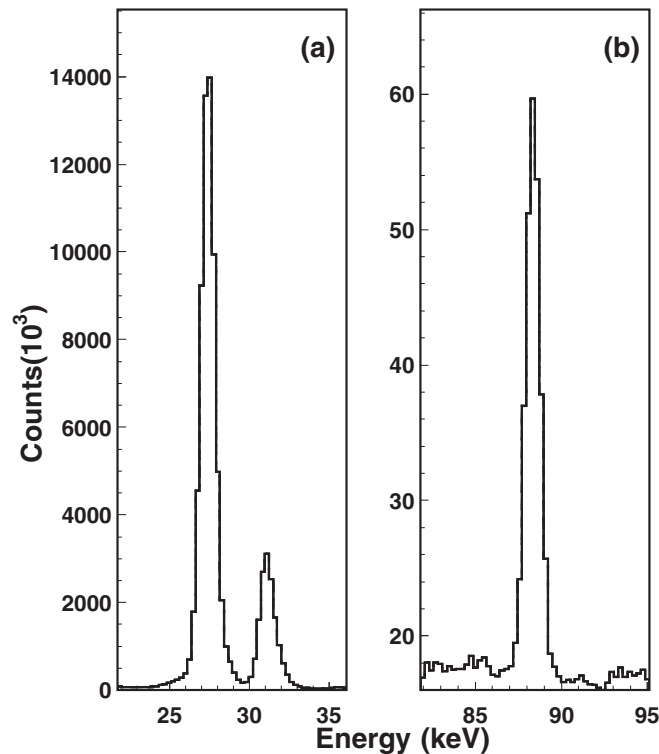


FIG. 3. Spectra for the two energy regions of interest in this measurement, the one on the left including the tellurium K x rays and the one on the right, the γ -ray peak at 88.2 keV. These correspond to the full spectrum presented in Fig. 2.

¹To establish the weighting, we used the intensities of the individual x-ray components from Table 7a in Ref. [21].

TABLE I. The contributions of identified impurities to the energy regions of the tellurium K x-ray peaks and the 88.2-keV γ -ray peak.

| Source | Contaminant | Contaminant contribution (%) |
|--------------------|------------------------|------------------------------|
| To K x-ray peaks | | |
| ^{121}Te | Sb K x rays | 0.0162(12) |
| ^{110m}Ag | Cd+Ag K x rays | 0.000458(10) |
| ^{121m}Te | Te+Sb K x rays | 0.0022(3) |
| ^{123m}Te | Te K x rays | 0.095(3) |
| ^{124}Sb | Te K x rays | 0.0135(3) |
| ^{125m}Te | Te K x rays | 2.77(8) |
| ^{125m}Te | 35.5-keV γ ray | 0.176(2) |
| ^{129}Te | 27.8-keV γ ray | 0.11(4) |
| ^{129m}Te | Te K x rays | |
| | +27.8-keV γ ray | 0.289(11) |
| ^{131}I | Xe K x rays | 0.0346(8) |
| ^{131m}Xe | Xe K x rays | 0.005(5) |
| To 88.2-keV peak | | |
| ^{123m}Te | 88.5-keV γ ray | 0.078(2) |

The count totals recorded for the combined K x-ray peaks and for the 88.2-keV peak both appear in Table II, with their respective impurity components listed immediately below them. The impurity totals correspond to the percentage breakdowns given in Table I.

C. Efficiency ratios

In order to determine α_K for the 88.2-keV $M4$ transition in ^{127}Te , we require the efficiency ratio, $\epsilon_{\gamma 88.2}/\epsilon_{K 28.0}$ as can

TABLE II. Corrections to the ^{127}Te K x rays and the 88.2-keV γ ray as well as the additional information required to extract a value for α_K .

| Quantity | Value | Source |
|---|--------------------------|------------|
| Te ($K_\alpha + K_\beta$) x rays | | |
| Total counts | $7.8923(12) \times 10^7$ | Sec. IV A |
| Impurities | $-2.77(7) \times 10^6$ | Sec. IV B |
| β -decay contribution | $-2.78(5) \times 10^6$ | Sec. IV D |
| Lorentzian correction | +0.12(2)% | Sec. IV E |
| Net corrected counts, $N_{K 28.0}$ | $7.346(9) \times 10^7$ | |
| ^{127}Te 88.2-keV γ ray | | |
| Total counts | $1.760(13) \times 10^5$ | Sec. IV A |
| Impurities | $-1.38(4) \times 10^2$ | Sec. IV B |
| Net corrected counts, $N_{\gamma 88.2}$ | $1.759(13) \times 10^5$ | |
| Efficiency calculation | | |
| $\epsilon_{\gamma 88.0}/\epsilon_{K 22.6}$ | 1.069(8) | Ref. [6] |
| $\epsilon_{\gamma 88.2}/\epsilon_{\gamma 88.0}$ | 1.0011(1) | Ref. [10] |
| $\epsilon_{K 22.6}/\epsilon_{K 28.0}$ | 0.940(5) | Ref. [10] |
| $\epsilon_{\gamma 88.2}/\epsilon_{K 28.0}$ | 1.006(9) | |
| Evaluation of α_K | | |
| $N_{K 28.0}/N_{\gamma 88.2}$ | 418(3) | This table |
| Relative attenuation | +0.9(4)% | Sec. IV F |
| ω_K | 0.875(4) | Ref. [15] |
| α_K for 88.2-keV transition | 484(6) | Eq. (1) |

be seen in Eq. (1). Following the same procedure as the one we used in analyzing the decay of ^{119m}Sn [6], we employ as low-energy calibration the well-known decay of ^{109}Cd , which emits 88.0-keV γ rays and silver K x rays at a weighted average energy of 22.57 keV. Both are close in energy to the respective γ and x rays observed in the current measurement.

We obtain the required ratio, $\epsilon_{\gamma 88.2}/\epsilon_{K 28.0}$ from the following relation:

$$\frac{\epsilon_{\gamma 88.2}}{\epsilon_{K 28.0}} = \frac{\epsilon_{\gamma 88.0}}{\epsilon_{K 22.6}} \frac{\epsilon_{\gamma 88.2}}{\epsilon_{\gamma 88.0}} \frac{\epsilon_{K 22.6}}{\epsilon_{K 28.0}}. \quad (2)$$

We take the ^{109}Cd ratio $\epsilon_{\gamma 88.0}/\epsilon_{K 22.6}$ from our previously reported measurement [6]. The ratio $\epsilon_{\gamma 88.2}/\epsilon_{\gamma 88.0}$ is very nearly unity and determined with high precision from our known detector efficiency curve calculated with the CYLTRAN code [10], while $\epsilon_{K 22.6}/\epsilon_{K 28.0}$ comes from a CYLTRAN calculation as well but in an energy region with higher relative uncertainty. Nevertheless, the energy span is not large so the uncertainty is only $\pm 0.5\%$. The values of all four efficiency ratios from Eq. (2) appear in the third block of Table II.

D. β -decay contribution

Although the decay of ^{127m}Te occurs predominantly through the 88.2-keV $M4$ electromagnetic transition, it also has a weak β -decay branch that populates the 57.6-keV level in ^{127}I . That level is additionally populated by another weak β -decay branch from the ground state of ^{127}Te , the decay of which is in secular equilibrium with the isomeric-state decay. This is illustrated in Fig. 1. The transition from the 57.6-keV state to the ground state of ^{127}I has $M1 + E2$ character with a measured mixing ratio and an $\alpha_{K 57.6}$ value of 3.16(5) [16]. This gives rise to iodine K x rays, which are unresolved from the tellurium x rays and thus must be corrected for.

If $N_{\gamma 57.6}$ is the total number of counts in the 57.6-keV peak, then the corresponding number of iodine K x rays that appear in the spectrum, $N_{K 29.2}$, is given by

$$N_{K 29.2} = N_{\gamma 57.6} \frac{\epsilon_{K 29.2}}{\epsilon_{\gamma 57.6}} \alpha_{K 57.6} \omega_{Ki}, \quad (3)$$

where $\omega_{Ki} = 0.882(4)$ is the fluorescence yield for iodine [15].

Using the same approach we took for Eq. (2), we obtain the required efficiency ratio, $\epsilon_{K 29.2}/\epsilon_{\gamma 57.6}$, from the following relation:

$$\frac{\epsilon_{K 29.2}}{\epsilon_{\gamma 57.6}} = \frac{\epsilon_{K 22.6}}{\epsilon_{\gamma 88.0}} \frac{\epsilon_{\gamma 88.0}}{\epsilon_{\gamma 57.6}} \frac{\epsilon_{K 29.2}}{\epsilon_{K 22.6}}, \quad (4)$$

in which our previously measured ^{109}Cd ratio [6] again plays a key role. Here it is inverted compared to Eq. (2) and has the value $\epsilon_{K 22.6}/\epsilon_{\gamma 88.0} = 0.935(7)$. Using our known detector efficiency curve calculated with the CYLTRAN code, we determine the other two ratios to be $\epsilon_{\gamma 88.0}/\epsilon_{\gamma 57.6} = 0.9796(11)$ and $\epsilon_{K 29.2}/\epsilon_{K 22.6} = 1.073(5)$. The final result for $\epsilon_{K 29.2}/\epsilon_{\gamma 57.6}$ thus becomes 0.983(9).

Substituting this result into Eq. (3), together with the measured number of counts in the 57.6-keV γ -ray peak, $N_{\gamma 57.6} = 1.0146(15) \times 10^6$, we obtain the contribution of the β -decay channels to the x-ray spectrum to be $N_{K 29.2} =$

TABLE III. Comparison of the measured α_K values for the 88.23(7)-keV $M4$ transition from ^{127m}Tm with calculated values based on three different theoretical models for dealing with the K -shell vacancy. Shown also are the percentage deviations, Δ , from the experimental value calculated as (experiment – theory)/theory. For a description of the various models used to determine the conversion coefficients, see text and Ref. [1].

| Model | α_K | Δ (%) |
|--------------------------|------------|--------------|
| Experiment | 484(6) | |
| Theory | | |
| No vacancy | 468.6(17) | +3.3(13) |
| Vacancy, frozen orbitals | 486.4(17) | –0.5(13) |
| Vacancy, SCF of ion | 483.1(17) | +0.2(13) |

$2.78(5) \times 10^6$. This result also appears in the first block of Table II.

E. Lorentzian correction

As explained in our previous papers (see, for example, Ref. [1]) we use a special modification of the GF3 program that allows us to sum the total counts above background within selected energy limits. To account for possible missed counts outside those limits, the program adds an extrapolated Gaussian tail. This extrapolated tail does not do full justice to x-ray peaks, whose Lorentzian shapes reflect the finite widths of the atomic levels responsible for them. To correct for this effect we compute simulated spectra using realistic Voigt functions to generate the x-ray peaks, and we then analyze them with GF3, following exactly the same fitting procedure as is used for the real data, to ascertain how much was missed by this approach. The resultant correction factor appears as a percentage in the first block of Table II.

F. Attenuation in the sample

As described in Sec. III A, our source was a thin layer of tellurium enclosed between Mylar sheets. We obtained the attenuation both of the tellurium x rays and of the 88.2-keV γ ray using standard tables of attenuation coefficients [22]. Since we are striving to evaluate α_K from Eq. (1), what matters in that context is the attenuation for the x rays relative to that for the 88.2-keV γ ray. We determine that the x rays suffered 0.9(4)% more attenuation and it is this number that appears in the bottom block of Table II.

V. RESULTS AND DISCUSSION

A. α_K for the 88.2-keV transition

The third and fourth blocks of Table II contain all the data required to determine the value of α_K for the 88.2-keV transition from Eq. (1). The result, $\alpha_K = 484(6)$, agrees very closely with the only previous measurement [13] of this ICC, 484(23), but has replaced a value having $\pm 5\%$ precision with one having $\pm 1.2\%$.

Our measured α_K value is compared with three different theoretical calculations in Table III. All three calculations

TABLE IV. Relative intensities of γ rays deexciting the 418.0-keV level in ^{127}I .

| Energy (keV) | Relative γ -ray intensity | | |
|-----------------|----------------------------------|-----------|-----------|
| | Ref. [23] | Ref. [24] | This work |
| 418.0 | 100 | 100 | 100 |
| 360.3 | 14.8(1) | 13.6(4) | 15.2(1) |
| 215.2 | 3.9(2) | 3.91(17) | 4.44(7) |

were made within the Dirac-Fock framework, but one ignores the presence of the K -shell vacancy while the other two include it using different approximations: the frozen orbital approximation, in which it is assumed that the atomic orbitals have no time to rearrange after the electron's removal; and the SCF approximation, in which the final-state continuum wave function is calculated in the self-consistent field (SCF) of the ion, assuming full relaxation of the ion orbitals. To obtain these results we used the value 88.23(7) keV [16] for the ^{127m}Te transition energy. The experimental uncertainty in this number is reflected in the uncertainties quoted on the theoretical values of α_K in the table.

The percentage deviations given in Table III indicate excellent agreement between our measured result and the two calculations that include some provision for the atomic vacancy. Our measurement disagrees by 2.5 standard deviations with the calculation that ignores the vacancy. This outcome is consistent with our previous six precise α_K measurements on $E3$ and $M4$ transitions in ^{111}Cd [8], ^{119}Sn [6,7], ^{134}Cs [3,4], ^{137}Ba [3,4], ^{193}Ir [1,2], and ^{197}Pt [5], all of which agreed well with calculations that included the vacancy, and disagreed—some by many standard deviations—with the no-vacancy calculations.

B. β -decay branching ratio from ^{127m}Te

As remarked in Sec. IV D and illustrated in Fig. 1, the 57.6-keV level in ^{127}I is populated by two β -decay branches, one from the isomeric state and the other from the ground state of ^{127}Te . Since the ground state has a half-life of 9.35 hours and our spectrum was acquired more than a month after activation, the two decays were certainly in secular equilibrium for our measurement. Under these conditions, we see from the decay scheme that the 57.6-keV state is fed ~ 13 times more strongly from the isomer than it is from the ground state. As a consequence, it is possible to use the ratio of intensities of the γ -ray peaks at 57.6 and 88.2 keV to extract a rather precise value for the strength of the stronger β branch even though the strength of the weaker branch remains relatively imprecise.

If we represent the β branching ratio from the isomer by R , then we can write

$$\frac{R}{1-R} = \frac{N_{\gamma 57.6} \epsilon_{\gamma 88.2} (1 + \alpha_{T 57.6})}{N_{\gamma 88.2} \epsilon_{\gamma 57.6} (1 + \alpha_{T 88.2})} - 0.00155(21), \quad (5)$$

where the α_T values are the total ICCs for the indicated transitions, $\alpha_{T 57.6} = 3.72(3)$ and $\alpha_{T 88.2} = 1138(5)$, for which we have used the Dirac-Fock “frozen orbital” calculated

values. The numerical term, 0.00155(21) is the product of the β branching ratio for the transition from the ground state of ^{127}Te to the 418.0-keV state in ^{127}I , and the branching ratio for the subsequent electromagnetic decay of the 418.0-keV state to populate the 57.6-keV state. We take the former, 0.0119(16), from Ref. [16], while the latter we obtain from our own γ -ray spectrum.

Table IV lists the relative intensities we measure for the three γ rays that deexcite the 418-keV state in ^{127}I . The two previous measurements, with which our results are compared, were published a half century ago, with now-untraceable efficiency calibrations and likely rather optimistic uncertainties. We consider our results to be more reliable and use them exclusively. After adjusting for internal conversion, we determine the electromagnetic branching ratio for the 360.3-keV transition, which feeds the 57.6-keV state, to be 0.128(1). In addition, the 215.2-keV transition has a small probability for populating the 57.6-keV state (see Fig. 1) so the total feeding of the 57.6-keV state from the decay of the 418.0-keV state becomes 0.130(1). We use this result in determining the numerical term in Eq. (5).

Solving Eq. (5) for R , we determine the β branching of the isomeric state, ^{127m}Te , to be 2.14(3)%. This compares favorably with, but is 7 times more precise than, the only previous measurement of this quantity, 2.4(2)%, which was published in 1970 [24].

VI. CONCLUSIONS

Our measurement of the K -shell internal conversion coefficient for the 88.2-keV $M4$ transition from ^{127m}Te has yielded a value, $\alpha_K = 484(6)$, which agrees with a version of the Dirac-Fock theory that includes the atomic vacancy. It disagrees (by $\sim 2.5\sigma$) with theory if the vacancy is ignored. We have now made seven precise α_K measurements for $E3$ and $M4$ transitions in nuclei with a wide range of Z values. Their corresponding conversion-electron energies also ranged widely, from ~ 4 keV in ^{193}Ir to ~ 630 keV in ^{137}Ba . These measurements together present a consistent pattern that supports the Dirac-Fock theory for calculating internal conversion coefficients provided that it takes account of the atomic vacancy.

As a byproduct of the current measurement, we have also improved considerably the β -branching ratio for the 88.2-keV isomeric state in ^{127}Te .

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