# Thermal neutron capture cross sections and neutron separation energies for ${}^{23}Na(n,\gamma)$

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Prompt thermal neutron capture  $\gamma$ -ray cross sections  $\sigma_{\gamma}$  were measured for the <sup>23</sup>Na( $n,\gamma$ ) reaction with guided cold neutron beams at the Budapest Reactor. The <sup>24</sup>Na  $\gamma$ -ray cross sections were internally standardized with a stoichiometric NaCl target by using standard <sup>35</sup>Cl( $n,\gamma$ )<sup>36</sup>Cl  $\gamma$ -ray cross sections. Transitions were assigned to levels in <sup>24</sup>Na based primarily upon the known nuclear structure information from the literature, producing a nearly complete neutron capture decay scheme. The total radiative thermal neutron cross section  $\sigma_0$  was determined from the sum of prompt  $\gamma$ -ray cross section populating the ground state as 0.540 (3) b, and from the activation  $\gamma$ -ray cross sections for the decay of <sup>24</sup>Na as 0.542 (3) b. The isomer cross section  $\sigma_0$  (<sup>23</sup>Na<sup>m</sup>,  $t_{1/2} = 20.20$  ms) = 0.501 (3) b and the <sup>24</sup>Na neutron separation energy  $S_n = 6959.352$  (18) keV were also determined in these experiments. New level spins and parities were proposed on the basis of new transition assignments and the systematics of reduced transition probabilities for the primary  $\gamma$  rays.

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## I. INTRODUCTION

Precise prompt thermal neutron capture  $\gamma$ -ray cross sections  $\sigma_{\nu}$  have been measured for all elements with Z = 1-83, 90, and 92, except for He and Pm, at the Budapest Reactor [1]. These data have been evaluated together with additional information from the literature and compiled into the Evaluated Gamma-ray Activation File (EGAF) [2] and published in the Handbook of Prompt Gamma Activation Analysis [3]. For low-Z isotopes, where the decay schemes are completely measured, the total radiative thermal neutron capture cross section  $\sigma_0$  can be determined by  $\sigma_0 = \Sigma \sigma_{\gamma}(GS) = \Sigma \sigma_{\gamma}(CS)$ using the transitions either feeding the ground state or deexciting the capture state. The total radiative thermal neutron capture cross sections can also be determined in these experiments from the activation  $\gamma$ -ray cross sections and their decay emission probabilities  $P_{\gamma}$  if the half-lives are sufficiently short. In this paper we report our determination of  $\sigma_0$  for the <sup>23</sup>Na( $n, \gamma$ ) reaction.

#### **II. EXPERIMENT**

Neutron capture  $\gamma$ -ray cross sections were measured with cold guided neutron beams at the 10 MW Budapest Reactor [1]. Neutrons enter the evacuated target holder and continue to the beam stop at the rear wall of the guide hall. The target station is located  $\approx$ 35 m from the reactor where both primary and secondary  $\gamma$  rays can be measured in low-background conditions. At the time of the experiment, the thermal-equivalent neutron flux was 5 × 10<sup>7</sup> *n* cm<sup>-2</sup> s<sup>-1</sup>.

Prompt  $\gamma$  rays from the target were measured with an *n*type high-purity, 25% relative efficiency, germanium (HPGe) detector with closed-end coaxial geometry located 23.5 cm from the target. The detector is Compton suppressed by a bismuth-germanate (BGO) scintillator guard detector annulus surrounded by 10-cm-thick lead shielding. Energy resolution was 1.5 keV at 91 keV, 2.1 keV at 1.3 MeV, and 4.3 keV at 6.4 MeV. Counting efficiency was calibrated from 50 keV to 10 MeV with radioactive sources and  $(n,\gamma)$  reaction  $\gamma$  rays to a precision of better that 1% from 500 keV to 6 MeV and better than 3% at all other energies [4]. The  $\gamma$ -ray spectra were analyzed using the HYPERMET PC program [4,5].

The <sup>24</sup>Na thermal neutron capture  $\gamma$ -ray cross sections were internally calibrated with a stoichiometric, >99.5% pure NaCl target assuming the  $\sigma_{\nu}({}^{36}\text{Cl}, 1951.1 \text{ keV}) = 6.51(2) \text{ b} [6]$  and using the relative <sup>36</sup>Cl  $\gamma$ -ray emission probabilities of Molnar et al. [7]. Besides the usual elements (H, B, C, N, F, Al, and Pb), whose concentrations did not exceed background levels, no other elements could be observed in the prompt  $\gamma$ -ray spectrum at the 0.1% level. For the homogenous target the measured cross section does not depend on the spatial distribution of the neutron flux. No target impurities were observed in the prompt  $\gamma$ -ray calibration spectrum. Both the chlorine and sodium isotopes have a 1/v cross-section energy dependence [2] so the respective  $\gamma$ -ray intensity ratios used in these cross-section calibrations are independent of neutron energy distribution. No fast neutrons were present in the guided neutron beam.

The total radiative neutron capture  $\gamma$ -ray cross section  $\sigma_0$  for <sup>23</sup>Na was also determined by using the activation  $\gamma$ -ray cross sections, observed in both the prompt  $\gamma$ -ray data and after bombardment, with the  $\gamma$ -ray transition probabilities  $P_{\gamma}$  from the Decay Database Evaluation Project (DDEP) [8] database as shown in Eq. (1). The observed cross section is the corrected for

$$\sigma_0 = \frac{\sigma_\gamma}{P_\gamma b},\tag{1}$$

the in-beam saturation b, as shown in Eq. (2), where t is

$$b = 1 - \frac{1 - e^{-\lambda t}}{\lambda t},\tag{2}$$

the bombardment time and the decay constant  $\lambda = \ln(2)/t_{1/2}$ .

TABLE I. <sup>23</sup>Na( $n, \gamma$ ) thermal neutron capture  $\gamma$ -ray energies and cross sections measured in this work.

$E_{\gamma}$ (keV)	$\sigma_{\gamma}$ (b)	Placement (initial $\rightarrow$ final)	$E_{\gamma}$ (keV)	$\sigma_{\gamma}$ (b)	Placement (initial $\rightarrow$ final)
91 004 (15)	0 244 (3)	$563 \rightarrow 472$	2414 33 (3)	0.02519(23)	$2978 \rightarrow 563$
$242.30(9)^{a}$	0.000110(11)	$3656 \rightarrow 3413$	$2426.44(22)^{b}$	0.000143(23)	$4940 \rightarrow 2513$
373 24 (6)	0.000075(7)	$3745 \rightarrow 3372$	2431 74 (8)	0.000563 (23)	$2904 \rightarrow 472$
$387.98(18)^{b}$	0.000028 (6)	$3977 \rightarrow 3589$	$2482.9(5)^{b}$	0.000103 (20)	$6073 \rightarrow 3589$
390 51 (15) <sup>b</sup>	0.000044 (6)	$2904 \rightarrow 2513$	2505 26 (4)	0.01735(19)	$2978 \rightarrow 472$
$464\ 47\ (12)^{b}$	0.000011(0)	$4442 \rightarrow 3977$	2503.20(1) 2512.43(13)	0.00107 (6)	$2513 \rightarrow 0$
472 205 (14)	0.50010(3)	$472 \rightarrow 0$	2512.43(13) 2517.54(3)	0.00107(0)	$CS \rightarrow 4442$
499 363 (22)	0.01474(22)	$1846 \rightarrow 1347$	2517.54(5) 2521.92(21)	0.0068 (6)	$6111 \rightarrow 3589$
501 35 (4)	0.01474(22) 0.00308(12)	$1846 \rightarrow 1347$	2521.92(21) 2524 02(7) <sup>a</sup>	0.00008(0)	$3866 \rightarrow 1341$
504.55 (5)	0.00508(12) 0.00140(8)	$1846 \rightarrow 1341$	2524.02(7) 2546.00(19)	0.00038(3)	$5060 \rightarrow 1541$ $5060 \rightarrow 2513$
542 04 (12) <sup>b</sup>	0.00140(0)	$10+0 \rightarrow 10+1$ $4751 \rightarrow 4207$	2540.00(1))	0.000145(10)	$\frac{1000}{1442} \rightarrow \frac{1005}{1885}$
551 21 (4)	0.000042(7) 0.000348(13)	$4731 \rightarrow 4207$ $4207 \rightarrow 2656$	2555.00(7)	0.00040(3)	$4442 \rightarrow 1003$
552 721 (25)	0.000346(13)	$4207 \rightarrow 5030$	2505.2(3)	0.00008(3)	$2303 \rightarrow 0$
552.721(25)	0.0000003(14)	$CS \rightarrow 0407$	2574.9 (5)	0.000100(23)	$0237 \rightarrow 3082$
505.1/1 (11) (05.51 (2)	0.00925 (8)	$303 \rightarrow 0$	2588.58 (5)	0.00154 (5)	$3933 \rightarrow 1343$
003.31(3)	0.000159 (10)	$3977 \rightarrow 3372$	2591.55 (5)	0.00181 (3)	$5955 \rightarrow 1541$
$617.84(5)^{\circ}$	0.000143 (10)	$4207 \rightarrow 3589$	2595.19 (4)	0.00520 (6)	$4442 \rightarrow 1846$
685.54 (12)	0.000229 (17)	$3589 \rightarrow 2904$	2623.3 (4)	0.000081 (19)	$6251 \rightarrow 3628$
696.570 (20) <sup>a</sup>	0.00015 (17)	$4442 \rightarrow 3745$	2630.31 (5)	0.00296 (4)	$3977 \rightarrow 1347$
702.14 (3)	0.000312 (10)	$CS \rightarrow 6257$	2635.36 (12) <sup>b</sup>	0.000208 (18)	$3977 \rightarrow 1341$
708.051 (15)	0.001266 (16)	$CS \rightarrow 6251$	2657.88 (19)	0.000156 (17)	$6247 \rightarrow 3589$
711.865 (15)	0.00469 (4)	$CS \rightarrow 6247$	2661.55 (14) <sup>b</sup>	0.000211 (16)	$6251 \rightarrow 3589$
737.15 (7) <sup>d</sup>	0.00010 (3)	$CS \rightarrow 6222$	2701.8 (14)	0.000259 (18)	$4049 \rightarrow 1347$
773.685 (23)	0.000552 (12)	$4751 \rightarrow 3977$	2715.71 (5)	0.00303 (4)	$4562 \rightarrow 1846$
778.142 (13)	0.00602 (5)	$1341 \rightarrow 563$	2752.13 (4)	0.0593 (11)	$CS \rightarrow 4207$
781.295 (13)	0.01736 (14)	$1345 \rightarrow 563$	2763.19 (4)	0.00234 (4)	$CS \rightarrow 4196$
783.14 (4) <sup>c</sup>	0.00058 (4)	$1347 \rightarrow 563$	2808.32 (4)	0.01644 (16)	$3372 \rightarrow 563$
	0.00031 (4)	$CS \rightarrow 6172$	2833.7 (3) <sup>b</sup>	0.000080 (12)	$6247 \rightarrow 3413$
785.58 (7)	0.000206 (12)	$4442 \rightarrow 3656$	2849.97 (5)	0.00137 (3)	$3413 \rightarrow 563$
793.767 (15)	0.002119 (22)	$4207 \rightarrow 3413$	2860.26 (4)	0.01821 (22)	$4207 \rightarrow 1347$
813.36 (9)	0.000082 (10)	$4442 \rightarrow 3628$	2865.4 (5) <sup>c</sup>	0.0121 (9)	$4207 \rightarrow 1341$
835.259 (14)	0.01164 (10)	$4207 \rightarrow 3372$		0.0029 (3)	$4751 \rightarrow 4751$
852.25 (3)	0.000392 (13)	$4442 \rightarrow 3589$	2875.55 (17)	0.000112 (13)	$6247 \rightarrow 3372$
852.25 (3)		$5060 \rightarrow 4207$	2899.410 (15) <sup>a</sup>	0.000110 (13)	$3372 \rightarrow 472$
858.431 (24) <sup>b</sup>	0.00008 (3)	$3372 \rightarrow 2513$	2903.819 (22) <sup>a,c</sup>	0.00062 (10)	$2904 \rightarrow 0$
863.21 (5) <sup>a</sup>	0.00011 (5)	$5060 \rightarrow 4196$	2904.55 (4) <sup>c</sup>	0.00543 (7)	$4751 \rightarrow 4751$
869.173 (13)	0.1120 (15)	$1341 \rightarrow 472$	2910.6 (8)	0.000332 (16)	$CS \rightarrow 4049$
874.353 (13)	0.0796 (13)	$1347 \rightarrow 472$	2940.74 (4)	0.00340 (4)	$3413 \rightarrow 472$
886.729 (17)	0.00408 (4)	$CS \rightarrow 6073$	2977.35 (7)	0.000554 (23)	$2978 \rightarrow 0$
$906\ 04\ (4)^{a}$	0.00036 (9)	$4562 \rightarrow 3656$	2981 98 (4)	0.01385(14)	$CS \rightarrow 3977$
$943.4(5)^{b}$	0.000046(12)	$6251 \rightarrow 5309$	2993 62 (11)	0.00214(12)	$6407 \rightarrow 3413$
$992.40(14)^{\circ}$	0.000010(12)	$4621 \rightarrow 3628$	3015 79 (10)	0.000217(12) 0.000437(16)	$CS \rightarrow 3943$
<i>yy</i> 2.10(11)	0.000010(10)	$CS \rightarrow 5967$	3025 69 (8)	0.000137(10)	$3589 \rightarrow 563$
999 33 (15)	0.000087(13)	$3977 \rightarrow 2978$	5025.07 (0)	0.0094(4)	$CS \rightarrow 3933$
$1005 \ 914 \ (21)^{\circ}$	0.000000(11)	$3777 \rightarrow 2978$ $4751 \rightarrow 3745$	3002 64 (6) <sup>c</sup>	0.0030(4)	$C3 \rightarrow 5935$
1005.914 (21)	0.00100(3)	$-131 \rightarrow 5745$	5092.04 (0)	0.00075(4)	$CS \rightarrow 3866$
1012 05 (5)	0.00135(3)	$C3 \rightarrow J9JJ$	2004 202 (12)	0.00090(9)	$CS \rightarrow 3800$
1012.93(3)	0.000037(13)	$3433 \rightarrow 4442$	2006 45 (4)	0.0029(3)	$4442 \rightarrow 1347$
1018.40 (13)	0.000077(11) 0.000260(12)	$2904 \rightarrow 1883$	3090.43 (4) 2000.74 (4)	0.01903(21)	$4442 \rightarrow 1343$
1020.23 (4)	0.000500(12) 0.001525(21)	$4442 \rightarrow 5415$	3099.14 (4) 2116 96 (5)b	0.01319(17)	$4442 \rightarrow 1341$
1041.122 (23)	0.001535 (21)	$CS \rightarrow 5918$	5110.80 (5)° 2169.57 (11)	0.00470(6)	$3389 \rightarrow 4/2$
1057.00 (4)	0.000032 (10)	$2303 \rightarrow 1314$	5108.57 (11) 2172.91 (0)	0.000244 (14)	$0073 \rightarrow 2904$
1057.00 (4)	0.000355 (12)	$2904 \rightarrow 1846$	51/5.81 (9)	0.000313 (15)	$5060 \rightarrow 1885$
1092.02 (3)	0.001667 (19)	$29/8 \rightarrow 1885$	3181.77 (6)	0.00133 (4)	$3/45 \rightarrow 563$
1094.86 (3)	0.000696 (12)	$4/51 \rightarrow 3656$	3183.52 (26)*	0.00029 (3)	$3656 \rightarrow 4/2$
1097.2 (3) <sup>a</sup>	0.00020 (8)	$CS \rightarrow 5963$	3198.75 (5)	0.001111 (22)	$5045 \rightarrow 1846$
1108.12 (7)	0.000158 (10)	$CS \rightarrow 5851$	3209.31 (5)	0.00378 (5)	$3682 \rightarrow 472$
1131.35 (12)	0.000155 (22)	$2978 \rightarrow 1846$	3214.2 (5)	0.00539 (7)	$CS \rightarrow 3745$

$E_{\gamma}$	$\sigma_{\gamma}$	Placement	$E_{\gamma}$	$\sigma_{\gamma}$	Placement
(KeV)	(b)	$(initial \rightarrow final)$	(KeV)	(b)	$(initial \rightarrow final)$
1143.09 (14)	0.00026 (3)	$3656 \rightarrow 2513$	3231.68 (13)	0.000164 (14)	$5117 \rightarrow 1885$
1149.748 (19)	0.00580 (8)	$CS \rightarrow 5809$	3244.31 (7) <sup>a</sup>	0.000027 (10)	$6222 \rightarrow 2978$
1171.869 (21) <sup>a</sup>	0.00014 (4)	$2513 \rightarrow 1341$	3270.47 (6)	0.000564 (15)	$5117 \rightarrow 1846$
1183.8 (9) <sup>b</sup>	0.00007 (1)	$CS \rightarrow 5775$	3277.45 (5)	0.00382(5)	$CS \rightarrow 3682$
1208.34 (7) <sup>a</sup>	0.00023 (10)	$4621 \rightarrow 3413$	3295.878 (24) <sup>a</sup>	0.000026 (11)	$5809 \rightarrow 2513$
1219.95 (22)	0.000110 (14)	$4196 \rightarrow 2976$	3303.26 (6)	0.001251 (22)	$CS \rightarrow 3656$
1225.0 (6)	0.000087 (22)	$5918 \rightarrow 4694$	3330.6 (5)	0.00134 (3)	$CS \rightarrow 3628$
1229.26 (4)	0.00153 (3)	$4207 \rightarrow 2978$	3343.02 (5)	0.001168 (24)	$6247 \rightarrow 2904$
1231.65 (15)	0.000191 (21)	$3745 \rightarrow 2513$	3369.84 (6)	0.0008 (4)	$3933 \rightarrow 563$
1247.37 (5)	0.001186 (20)	$5809 \rightarrow 4562$		0.0140 (2)	$CS \rightarrow 3589$
1218.59 (5) <sup>a</sup>	0.000129 (22)	$4196 \rightarrow 2978$	3409.15 (5)	0.00237 (4)	$4751 \rightarrow 1341$
1282.714 (21)	0.00551 (5)	$1846 \rightarrow 563$	3413.74 (5)	0.00449 (6)	$3977 \rightarrow 563$
1292.4 (3) <sup>b</sup>	0.000037 (11)	$4196 \rightarrow 2904$	3471.27 (7)	0.00136 (6)	$3943 \rightarrow 472$
1314.55 (5)	0.000308 (11)	$5060 \rightarrow 3744$	3492.8 (6)	0.000454 (13)	$5339 \rightarrow 1846$
1322.202 (19)	0.00632 (5)	$1885 \rightarrow 563$	3504.78 (5)	0.00718 (8)	$3977 \rightarrow 472$
1330.52 (19) <sup>b</sup>	0.000075 (11)	$3216 \rightarrow 1885$	3545.89 (5)	0.00492 (6)	$CS \rightarrow 3413$
1337.646 (19)	0.00332 (3)	$4751 \rightarrow 3413$	3576.4 (3)	0.00025 (4)	$4049 \rightarrow 472$
1340.98 (22)	0.000098 (15)	$1341 \rightarrow 0$	3587.4 (5)	0.0619 (17)	$CS \rightarrow 3372$
1344.475 (19)	0.02119 (17)	$1345 \rightarrow 0$	3627.97 (7)	0.000766 (19)	$3628 \rightarrow 0$
1373.654 (20)	0.00818 (9)	$1846 \rightarrow 472$	3632.62 (8) <sup>c</sup>	0.00055 (7)	$4196 \rightarrow 563$
1378.80 (12)	0.00038 (4)	$4751 \rightarrow 3372$		0.000430 (18)	$5479 \rightarrow 5479$
1387.83 (8)	0.000278 (22)	$CS \rightarrow 5571$	3643.55 (6)	0.00735 (8)	$4207 \rightarrow 563$
1412.4 (8)	0.000075 (17)	$1885 \rightarrow 472$	3697.8 (11)	0.00100 (3)	$5045 \rightarrow 1347$
1415.8 (10)	0.000045 (17)	$5045 \rightarrow 3628$	3703.07 (12)	0.00224 (4)	$5045 \rightarrow 1341$
1420.00 (18) <sup>b</sup>	0.000117 (18)	$3933 \rightarrow 2513$	3713.47 (25)	0.000198 (20)	$5060 \rightarrow 1345$
1455.65 (3) <sup>a</sup>	0.00019 (11)	$5045 \rightarrow 3589$	3723.47 (12)	0.00165 (4)	$4196 \rightarrow 472$
1470.0 (3) <sup>b</sup>	0.000043 (9)	$5060 \rightarrow 3589$	3734.17 (15)	0.00055 (3)	$4207 \rightarrow 472$
1477.14 (15)	0.000089 (13)	$5455 \rightarrow 3977$	3744.04 (24)	0.000217 (21)	$3745 \rightarrow 1$
1480.32 (4)	0.001802 (22)	$CS \rightarrow 5479$	3770.51 (17)	0.000529 (17)	$5117 \rightarrow 1345$
1486.05 (4) <sup>b</sup>	0.001271 (17)	$3372 \rightarrow 1885$	3775.7 (3)	0.000091 (12)	$5117 \rightarrow 1341$
1504.76 (4)	0.00255 (3)	$CS \rightarrow 5455$	3866.20 (14) <sup>b</sup>	0.000324 (17)	$3866 \rightarrow 0$
1514.7 (4)	0.000101 (20)	$1514 \rightarrow 0$	3878.1 (5)	0.02316 (25)	$4442 \rightarrow 563$
1525.79 (15)	0.000093 (12)	$3372 \rightarrow 1846$	3934.34 (25)	0.000096 (13)	$6448 \rightarrow 2513$
1559.22 (4)	0.001488 (24)	$2904 \rightarrow 1345$	3942.84 (8)	0.000434 (16)	$3943 \rightarrow 0$
1562.36 (4)	0.00257 (4)	$2904 \rightarrow 1341$	3964.13 (15) <sup>b</sup>	0.000204 (15)	$5309 \rightarrow 1345$
1567.10 (5)	0.000537 (15)	$3413 \rightarrow 1846$	3969.07 (6)	0.00235 (4)	$4442 \rightarrow 472$
1578.0 (6) <sup>a</sup>	0.000064 (10)	$4143 \rightarrow 2565$	3981.35 (5)	0.0723 (8)	$CS \rightarrow 2978$
1584.05 (14)	0.000260 (11)	$4562 \rightarrow 2978$	3997.57 (6)	0.00170 (3)	$5339 \rightarrow 1341$
1620.39 (3)	0.00297 (3)	$CS \rightarrow 5339$	4055.23 (6)	0.00340 (18)	$CS \rightarrow 2904$
1631.19 (6) <sup>b</sup>	0.00086 (6)	$2978 \rightarrow 1347$	4058.47 (19)	0.00033 (4)	$4621 \rightarrow 563$
1633.223 (24)	0.00834 (9)	$2978 \rightarrow 1345$	4089.36 (7)	0.00216 (11)	$4562 \rightarrow 472$
1636.21 (3)	0.02511 (22)	$2978 \rightarrow 1341$	4106.92 (4) <sup>a</sup>	0.000158 (22)	$5953 \rightarrow 1885$
1646.16 (13) <sup>b</sup>	0.000130 (12)	$5060 \rightarrow 3413$	4108.1 (3)	0.000064 (10)	$5455 \rightarrow 1347$
1651.08 (12) <sup>b</sup>	0.00136 (13)	$CS \rightarrow 5309$	4132.55 (19)	0.000171 (13)	$5479 \rightarrow 1345$
1683.1 (3) <sup>b</sup>	0.000101 (14)	$4196 \rightarrow 2513$	4137.03 (7)	0.000836 (22)	$5479 \rightarrow 1341$
1685.7 (3) <sup>b</sup>	0.000079 (15)	$6247 \rightarrow 4562$	4142.6 (6)	0.000064 (10)	$4143 \rightarrow 0$
1693.677 (22) <sup>a</sup>	0.00064 (7)	$4207 \rightarrow 2513$	4187.39 (6) <sup>c</sup>	0.00795 (9)	$4751 \rightarrow 563$
1711.16 (16) <sup>b</sup>	0.00028 (3)	$5918 \rightarrow 4207$		0.00014 (6)	$6073 \rightarrow 6073$
1714.2 (4) <sup>a</sup>	0.00043 (3)	$4694 \rightarrow 2978$	4226.17 (13)	0.000174 (11)	$6073 \rightarrow 1846$
1743.26 (8)	0.000444 (17)	$3589 \rightarrow 1846$	4278.59 (24)	0.000084 (10)	$4751 \rightarrow 472$
1766.92 (5)	0.000321 (11)	$CS \rightarrow 5192$	4361.38 (25)	0.000097 (11)	$6247 \rightarrow 1885$
1770.18 (6)	0.000332 (12)	$3656 \rightarrow 1885$	4376.05 (12)	0.00060 (3)	$4940 \rightarrow 563$
1773.02 (3)	0.001007 (20)	$4751 \rightarrow 2978$	4394.9 (5)	0.000055 (10)	$CS \rightarrow 2565$
1809.96 (9)	0.000189 (13)	$3656 \rightarrow 1846$	4405.06 (17)	0.000305 (19)	$6251 \rightarrow 1846$
1831.87 (5)	0.00092 (7)	$5809 \rightarrow 3977$	4433.8 (3)	0.000069 (9)	$5776 \rightarrow 1341$

## TABLE I. (Continued.)

4445.69 (7)

0.00249 (3)

 $\text{CS} \rightarrow 2513$ 

 $CS \rightarrow 5117$ 

1841.92 (3)

0.00246 (19)

$\overline{E_{\gamma}}$	$\sigma_{\gamma}$ (b)	Placement $(initial \rightarrow final)$	$E_{\gamma}$	$\sigma_{\gamma}$ (b)	Placement $(initial \rightarrow final)$
(Ke V)	(0)	$(\min a) \rightarrow \min a)$	(KC V)	(0)	$(\min a) \rightarrow \min a)$
1846.67 (8)	0.00041 (4)	$4751 \rightarrow 2904$	4462.43 (8)	0.000747 (16)	$5809 \rightarrow 1347$
1859.2 (8)	0.00034 (4)	$3745 \rightarrow 1885$	4466.96 (7)	0.001558 (24)	$4940 \rightarrow 472$
1876.24 (5) <sup>a</sup>	0.00034 (3)	$6073 \rightarrow 4196$	4481.4 (7)	0.001005 (18)	$5045 \rightarrow 563$
1885.35 (3)	0.00347 (4)	$1885 \rightarrow 0$	4496.02 (7)	0.00216 (3)	$5060 \rightarrow 563$
1899.15 (3) <sup>c</sup>	0.0021 (3)	$3745 \rightarrow 1846$	4521.3 (3)	0.000070 (8)	$5863 \rightarrow 1341$
	0.0047 (3)	$CS \rightarrow 5060$	4553.79 (15)	0.000389 (24)	$5117 \rightarrow 563$
1914.31 (3)	0.00602 (6)	$CS \rightarrow 5045$	4562.57 (25)	0.000125 (12)	$6448 \rightarrow 1885$
1928.15 (3) <sup>b</sup>	0.00471 (5)	$4442 \rightarrow 2513$	4571.76 (10) <sup>c</sup>	0.00044 (7)	$5045 \rightarrow 472$
1950.03 (3)	0.00894 (8)	$2513 \rightarrow 563$		0.00070 (7)	$5918 \rightarrow 5918$
1964.3 (4)	0.000026 (12)	$6407 \rightarrow 4442$	4586.99 (11)	0.000808 (20)	$5060 \rightarrow 472$
2009.76 (9) <sup>a</sup>	0.00093 (10)	$5953 \rightarrow 3943$	4619.1 (18) <sup>a</sup>	0.000012 (4)	$5967 \rightarrow 1347$
2019.74 (3) <sup>b,c</sup>	0.0023 (2)	$CS \rightarrow 4940$	4628.72 (14)	0.000226 (13)	$5192 \rightarrow 563$
	0.0009 (2)	$3866 \rightarrow 1846$	4644.47 (18)	0.000296 (13)	$5117 \rightarrow 472$
2024.97 (3)	0.0318 (3)	$3372 \rightarrow 1347$	4725.94 (12)	0.000625 (22)	$6073 \rightarrow 1347$
2027.02 (4)	0.00415(9)	$3372 \rightarrow 1345$	4727.601 (13) <sup>a</sup>	0.00012 (4)	$6073 \rightarrow 1345$
2030.15 (3)	0.02094 (18)	$3372 \rightarrow 1341$	4730.75 (11)	0.00177 (4)	$6073 \rightarrow 1341$
2040.9 (3) <sup>b</sup>	0.000078 (19)	$5455 \rightarrow 3413$	4775.3 (13)	0.000919 (21)	$5339 \rightarrow 563$
2047.91 (9)	0.000534 (20)	$4562 \rightarrow 2513$	4829.68 (18)	0.000151 (12)	$6176 \rightarrow 1347$
2051.29 (26)	0.000098 (17)	$CS \rightarrow 4908$	4836.06 (24) <sup>b</sup>	0.000113 (11)	$5309 \rightarrow 472$
2066.4 (3)	0.001300 (21)	$3413 \rightarrow 1347$	4866.0 (9)	0.00008 (5)	$5339 \rightarrow 472$
2071.58 (3)	0.00567(6)	$3413 \rightarrow 1341$	4875.18 (7) <sup>a</sup>	0.000032 (5)	$6222 \rightarrow 1347$
2087.44 (4)	0.000707(15)	$3933 \rightarrow 1846$	4890.69 (19)	0.0017 (7)	$5455 \rightarrow 563$
2107.7 (3)	0.000062 (13)	$5479 \rightarrow 5479$	4900.356 (13) <sup>a</sup>	0.0007 (2)	$6247 \rightarrow 1347$
2108.12 (7) <sup>a</sup>	0.00013 (5)	$4621 \rightarrow 2513$	4902.0 (3)	0.0010 (4)	$6247 \rightarrow 1345$
2117.91 (12) <sup>a</sup>	0.00009 (4)	$5863 \rightarrow 3744$	4905.517 (13) <sup>a</sup>	0.000016 (5)	$6247 \rightarrow 1341$
2130.92 (12)	0.000167 (15)	$3977 \rightarrow 1846$	4914.7 (4)	0.00020 (8)	$5479 \rightarrow 563$
2139.22 (9)	0.000210 (16)	$5117 \rightarrow 2978$	4981.79 (21)	0.00035 (11)	$5455 \rightarrow 472$
2208.25 (3)	0.02682 (24)	$CS \rightarrow 4751$	5006.32 (21)	0.00024 (8)	$5479 \rightarrow 472$
2219.84 (6)	0.00080 (5)	$5809 \rightarrow 3589$	5059.75 (20)	0.000113 (9)	$5060 \rightarrow 0$
2237.55 (7)	0.00121 (5)	$4751 \rightarrow 2513$	5073.33 (8)	0.00237 (3)	$CS \rightarrow 1885$
2243.27 (15)	0.0041 (3)	$3589 \rightarrow 1345$	5112.79 (8)	0.00289 (4)	$CS \rightarrow 1846$
2247.38 (12)	0.00059 (3)	$3589 \rightarrow 1341$	5191.5 (15)	0.000037 (9)	$5192 \rightarrow 0$
2265.70 (21)	0.00022 (3)	$CS \rightarrow 4694$	5245.5 (5)	0.000051 (10)	$5809 \rightarrow 563$
2271.0 (3) <sup>c</sup>	0.00012 (3)	$5953 \rightarrow 3682$	5336.7 (3)	0.000068 (9)	$5809 \rightarrow 472$
	0.00012 (3)	$6247 \rightarrow 3977$	5445.28 (12)	0.000855 (18)	$5918 \rightarrow 472$
2279.5 (3) <sup>a</sup>	0.000035 (4)	$6257 \rightarrow 3977$	5493.8 (3)	0.000066 (9)	$5967 \rightarrow 472$
2281.65 (24) <sup>a</sup>	0.000261 (25)	$3628 \rightarrow 1347$	5508.4 (6)	0.000041 (9)	$6073 \rightarrow 563$
2286.33 (7)	0.00086 (4)	$3628 \rightarrow 1341$	5599.73 (10)	0.000694 (18)	$6073 \rightarrow 472$
2297.04 (17) <sup>b</sup>	0.000293 (27)	$4143 \rightarrow 1846$	5612.75 (17)	0.00132 (12)	$CS \rightarrow 1347$
2309.17 (3) <sup>a</sup>	0.00018 (4)	$3656 \rightarrow 1347$	5617.14 (8)	0.0228 (3)	$CS \rightarrow 1341$
2310.79 (5) <sup>a</sup>	0.00019 (4)	$4196 \rightarrow 1885$	5658.41 (7) <sup>a</sup>	0.000032 (5)	$6222 \rightarrow 563$
2314.33 (3) <sup>a</sup>	0.000163 (24)	$3656 \rightarrow 1343$	5683.592 (12) <sup>a</sup>	0.00005 (3)	$6247 \rightarrow 563$
2334.71 (21)	0.000132 (17)	$3682 \rightarrow 1347$	5703.91 (18)	0.000162 (12)	$6176 \rightarrow 472$
2337.83 (8)	0.000570 (20)	$CS \rightarrow 4621$	5774.59 (13)	0.001367 (24)	$6247 \rightarrow 472$
2340.48 (18)	0.000178 (20)	$2904 \rightarrow 563$	5784.16 (25)	0.000125 (11)	$6257 \rightarrow 472$
2349.78 (12)	0.000288 (17)	$4196 \rightarrow 1846$	6111.1 (3)	0.000082 (12)	$6111 \rightarrow 0$
2360.89 (3)	0.00834 (8)	$4207 \rightarrow 1846$	6246.641 (12) <sup>a</sup>	0.000014 (5)	$6247 \rightarrow 0$
2376.8 (13) <sup>a</sup>	0.000009 (3)	$5967 \rightarrow 3589$	6395.48 (9)	0.1005 (11)	$CS \rightarrow 563$
2397.24 (4)	0.00586 (9)	$CS \rightarrow 4562$	6406.4 (7)	0.000042 (9)	$6407 \rightarrow 0$
2400.18 (6)	0.00126 (3)	$3745 \rightarrow 1345$	6486.35 (11)	0.00233 (3)	$CS \rightarrow 472$
2403.31 (4)	0.00154 (3)	$3745 \rightarrow 1341$	6958.8 (5)	0.000015 (3)	$CS \rightarrow 0$

TABLE I. (Continued.)

<sup>a</sup>From <sup>24</sup>Na Adopted Levels Gammas [9], intensity normalized to adopted branching ratios.

<sup>b</sup>Placed on the basis of energy sums and consistency with the level scheme.

<sup>c</sup>Transition multiply placed. Intensity divided on the basis of adopted  $\gamma$ -ray branching ratios [9] and level scheme.

<sup>d</sup>Transition adopted from Hungerford *et al.* [10].

TABLE II. Cross sections populating and depopulating levels in <sup>24</sup>Na from the  $(n,\gamma)$  reaction. Level energies are calculated from a least-squares fit of the  $\gamma$ -ray energies to the level scheme. The  $J^{\pi}$  values are from the Evaluated Nuclear Structure Data File (ENSDF) [11], modified as discussed in the text.

Level energy keV	$J^{\pi}$	In (b)	Out (b)	Net <sup>a</sup> (b)	Level energy eV	$J^{\pi}$	In (b)	Out (b)	Net <sup>a</sup> (b)
0	$4^{+}$	0.540 (3)	0.0	0.540 (3)	4561.94 (4)	$2^{-}$	0.0071(1)	0.0063 (2)	0.0008 (2)
472.208 (14)	$1^{+}$	0.500 (4)	0.501 (3)	-0.001(5)	4621.50 (7)	$(2)^{+}$	0.00057 (2)	0.00070 (12)	0.00013 (12)
563.189 (12)	$2^{+}$	0.252 (2)	0.254 (3)	0.002 (3)	4693.4 (3)	(3-)	0.00031 (4)	0.00011 (4)	0.00020 (6)
1341.426 (17)	$2^{+}$	0.121 (1)	0.118 (2)	-0.003(2)	4750.876 (24)	2-	0.0268 (2)	0.0274 (3)	-0.0005(4)
1344.512 (17)	3+	0.0386 (5)	0.0386 (2)	-0.0000(5)	4908.022 (16)	2+,3	0.00010 (2)	0.0	000010 (2)
1346.591 (17)	$1^{+}$	0.0814 (6)	0.0802 (13)	-0.0012 (14)	4939.55 (4)	1-	0.0023 (2)	0.0023 (1)	0.0000 (2)
1514.5 (5)	$5^{+}$	0.00003 (1)	0.00010(2)	0.00007 (2)	5044.96 (3)	$(1)^{-}$	0.00602 (6)	0.00603 (14)	-0.0.00001 (15)
1845.978 (22)	$2^{+}$	0.0342 (4)	0.0329 (3)	0.0013 (5)	5059.55 (5)	3-	0.0047 (3)	0.0047 (1)	0.0000 (3)
1885.438 (24)	3+	0.0087 (3)	0.0099 (1)	0.0012 (3)	5117.34 (5)	$(2^{-})$	0.0025 (2)	0.0022 (1)	0.0002 (2)
2513.276 (22)	3+	0.0108 (1)	0.0102 (1)	0.0006 (2)	5192.37 (8)	3-	0.00032(1)	0.0026 (2)	0.00006 (2)
2564.6 (3)	$4^{+}$	0.00012(1)	0.00022 (4)	-0.00010 (4)	5308.6 (3)	$2(^{+})$	0.00018 (2)	0.00032 (2)	-0.00014 (3)
2903.81 (4)	3+	0.0055 (2)	0.0059(1)	-0.0004(2)	5339.02 (8)	$2^{-}$	0.00297 (3)	0.00315 (6)	-0.00018 (7)
2977.681 (25)	$2^{+}$	0.0757 (8)	0.0752 (4)	0.0005 (9)	5454.52 (5)	$1^{-}, 2^{-}$	0.0025 (1)	0.0023 (7)	-0.0002(7)
3216.0 (3)	$(4^{+})$	0.0	0.000075 (11)	0.000075 (11)	5478.95 (5)	$1^{-}$	0.00180 (2)	0.0019 (1)	-0.0001 (1)
3371.696 (20)	$2^{-}$	0.0743 (17)	0.0748 (4)-	0.0005 (17)	5571.48 (9)	$(2^{+})$	0.00028 (2)	0.0	0.00028 (2)
3413.195 (23)	$1^{+}$	0.0116 (1)	0.0123 (1)	-0.0007(2)	5775.7 (4)	(3+)	0.00007 (8)	0.00007 (1)	0.00000 (8)
3589.26 (3)	$1^{+}$	0.0161 (2)	0.0169 (4)	0.0008 (5)	5809.40 (4)	2-	0.0058 (1)	0.0038 (1)	0.0020(1)
3628.29 (8)	3+	0.00156 (4)	0.00189 (5)	0.00033 (7)	5851.26 (8)	$(2^{+})$	0.00016(1)	0.0	0.00016(1)
3655.89 (3)	$(1)^{+}$	0.0029(1)	0.0022 (1)	0.0007(1)	5862.85 (17)	$(2^{+})$	0.00020 (8)	0.00016 (4)	0.00004 (9)
3681.71 (6)	$0^+$	0.00405 (6)	0.00391 (5)	-0.00013 (8)	5918.21 (4)	$(2^{-})$	0.00154 (2)	0.00192 (8)	-0.00039 (8)
3744.94 (3)	3-	0.00700 (9)	0.0071 (3)	0.0001 (3)	5953.39 (3)	$(1^{-})$	0.00155 (3)	0.00121 (11)	0.00034 (11)
3865.81 (5)	(3-)	0.0009(1)	0.0016 (2)	-0.0007(2)	5966.88 (20)	$0^+$	0.00009 (2)	0.00009(1)	0.00000 (2)
3933.33 (7)	$1^+, 2^+, 3$	0.0050 (4)	0.0050 (4)	0.0000 (4)	6072.611 (22)	$1^{+}$	0.00408 (4)	0.00406 (10)	0.00002 (11)
3943.38 (9)	$(2,3)^+$	0.00137 (10)	0.00179 (6)	0.00043 (12)	6176.20 (6)	$(1^{-},2^{-})$	0.00031 (4)	0.00031 (2)	0.0000 (4)
3977.18 (3)	$2^{+}$	0.0158 (2)	0.0153 (1)	0.0005 (2)	6222.26 (6)	$(1^+, 2^+)$	0.00010 (3)	0.000091 (12)	0.000009 (3)
4048.58 (11)	$0^{-}$	0.00033 (2)	0.00051 (4)	-0.00018 (5)	6247.492 (14)	2-	0.00469 (4)	0.0050 (5)	-0.0003 (5)
4143.09 (25)	$(4^{+})$	0.0	0.00042 (3)	-0.00042 (3)	6251.29 (3)	$(2^{-})$	0.00127 (2)	0.00064 (3)	0.00062 (4)
4196.28 (5)	$(2^+)$	0.0026 (1)	0.0028 (1)	-0.0002(1)	6257.19 (5)	(2 <sup>-</sup> )	0.00031 (1)	0.00027 (3)	0.00005 (3)
4206.990 (20)	$2^{-}$	0.0600 (10)	0.0630 (9)	-0.0030 (14)	6406.62 (4)	(2 <sup>-</sup> )	0.00060(1)	0.00033 (2)	0.00027 (3)
4441.532 (24)	$2^{-}$	0.0737 (7)	0.0751 (5)	-0.0014 (9)	6959.352 (18)	1 <sup>+b</sup>	0.0	0.537 (3)	0.537 (3)

<sup>a</sup>Net level feeding  $\sigma_{\gamma}$  (In)- $\sigma_{\gamma}$  (Out).

<sup>b</sup>The composite spin is  $\approx 96.5\% 1^+ + \approx 3.5\% 2^+$ .

## **III. RESULTS**

A pure <sup>23</sup>Na( $n,\gamma$ )  $\gamma$ -ray spectrum was acquired using an 0.16 g Na metal target contained in an evacuated target chamber. The  $\gamma$ -ray intensities were renormalized to cross sections using the NaCl standardization described above. The Na  $\gamma$ -ray energies and cross sections measured in this work and placed in the <sup>24</sup>Na level scheme are listed in Table I. A total of 326  $\gamma$  rays were placed in the level scheme depopulating 59 levels including the neutron capture state. Level energies have been calculated by a least-squares fit to the connecting  $\gamma$ -ray energies observed in this experiment.

## A. Prompt $\gamma$ -ray cross section

The cross-section balance through the level scheme is shown in Table II. The level scheme is nearly complete with 0.538(3) barns observed deexciting the neutron capture state and 0.540(3) barns populating the ground state. A total cross section of  $\Sigma \frac{E_{\gamma}}{S_n} \sigma_{\gamma} = 0.537(6)$  barns was calculated suggesting that the decay scheme is  $\approx 99\%$  complete. The population and

depopulation of intermediate levels is generally well balanced although clearly numerous weak transitions are missing.

The cross section calculated from the transition deexciting the short-lived isomeric first-excited state in <sup>24</sup>Na at 472.214 keV, with  $t_{1/2} = 20.18(10)$  ms, is 0.501(3) barns, as measured in the prompt  $\gamma$ -ray spectrum. This value is significantly higher than than previous measurements of 0.39(6) b by Groshev [12] and 0.40(3) b by Alexander [13] and adopted by Mughabghab [40], but more consistent with 0.476(11) b by Matsue and Yonezawa [14] and 0.504(11) by Szentmiklosi *et al.* [15].  $P_{\gamma}(472.2) = 0.904 \pm 0.006$ from ENSDF for <sup>23</sup>Na( $n,\gamma$ ) would give a total cross section of  $\sigma_0(^{24}\text{Na}^{m+g}) = 0.44$  b from Mughabghab's adopted isomeric cross section which is inconsistent with all later measured values, indicating that the early measurements are significantly discrepant for unknown reasons.

#### **B.** Activation $\gamma$ -ray cross section

The <sup>23</sup>Na( $n,\gamma$ ) total radiative cross section was also determined from the <sup>24</sup>Na<sup>m+g</sup> $\beta$ <sup>-</sup>decay scheme, shown in Fig. 1,



FIG. 1. (Color online) The <sup>24</sup>Na<sup>m+g</sup> decay scheme from the Table of Isotopes [41]. The  $\gamma$ -ray cross sections have been added from this work. Transition and level energies are in keV.

using the decay  $\gamma$ -ray cross sections, after correction for saturation, and the DDEP transition probabilities [8]. For the 1368.626(5) keV  $\gamma$  ray with  $P_{\gamma} = 0.999\,935\,3\,(5)$  we get  $\sigma_0 =$ 0.544(5) barns in the prompt spectrum and  $\sigma_0 = 0.542(3)$ barns in the decay spectrum. For the 2754.007(11) keV  $\gamma$ ray with  $P_{\gamma} = 0.99872(8)$  we get  $\sigma_0 = 0.540(3)$  barns in the prompt spectrum and  $\sigma_0 = 0.542(3)$  barns in the decay spectrum. The average total radiative thermal neutron cross section from decay,  $\sigma_0 = 0.542(3)$  b is consistent with the value determined from the prompt  $\gamma$  rays and we recommend  $\sigma_0 = 0.541(3)$  b from the weighted average of all values. This value is compared with earlier measurements in Table III. The new measurements are consistent with most previous values but inconsistent with Mughabghab's [40] recommended value of 0.517(4) barns. The weighted average value from all previous experiments gives  $\sigma_0 = 0.527(20)$  b which is consistent with our new value. We note that most previous measurements were performed by neutron activation which can be subject to significant summing corrections for the two intense coincident decay  $\gamma$  rays that would lead to an underestimating the total cross section when the source or detector distance is too small.

### C. Neutron separation energy

The HPGe detector was calibrated offline with  $\gamma$  rays from the <sup>14</sup>N( $n,\gamma$ ) [42] and <sup>35</sup>Cl( $n,\gamma$ ) [7] reactions. A total of 55 primary  $\gamma$  rays with energies from 553 to 6959 keV, which were observed to deexcite the capture state, were used to determine a neutron separation energy  $S_n = 6959.352(18)$ keV after correction for nuclear recoil. The consistency of  $S_n$ over a wide range of primary  $\gamma$ -ray energies demonstrates the stability of our energy calibration over the time of the experiment. Our  $S_n$  value is slightly lower than the recommended values of 6959.42(4) keV from the AME2012 Atomic Mass Evaluation [43].

TABLE III.	Comparison of <sup>23</sup> N	$\operatorname{Ia}(n,\gamma)$ total	lradiative	cross sectio
$\sigma_0$ measurement	its from this work.			

Author (year)	$\sigma_0$ (b) $^{24}$ Na $^m$	Method
Groshev (1955) [12]	$0.39 \pm 0.06^{a}$	Activation
Alexander (1963) [13]	$0.40 \pm 0.03$	Activation
Matsue (2004) [14]	$0.476 \pm 0.11$	PGAA
Szentmiklosi (2010) [15]	$0.504 \pm 0.011$	PGAA
Mughabghab (2006)	$0.40 \pm 0.03$	
This work	$0.501 \pm 0.003$	
	$^{24}$ Na <sup><math>m+g</math></sup>	
Meadows (1961) [16]	$0.47 \pm 0.06$	Pile oscillator
Pomerance (1951) [17]	$0.49 \pm 0.05^{b}$	Pile oscillator
Brooksbank (1955) [18]	$0.50 \pm 0.05$	Activation
Koehler (1963) [19]	$0.50 \pm 0.02$	Activation
Yamamuro (1970) [20]	$0.50 \pm 0.03$	Transmission
Harris (1953) [21]	$0.511 \pm 0.005^{b}$	Pile oscillator
De Corte (2003) [22]	$0.513 \pm 0.006$	Activation
Kennedy (2003) [23]	$0.515 \pm 0.021$	Activation
Grimeland (1955) [24]	$0.52\pm0.03$	Activation
Heft (1978) [25]	$0.523 \pm 0.005$	Activation
Ryves (1970) [26]	$0.527 \pm 0.005$	Activation
Szentmiklosi (2006) [27]	$0.527 \pm 0.008$	PGAA
Jozefowicz (1963) [28]	$0.53\pm0.01$	Activation
Wolf (1960) [29]	$0.531 \pm 0.008$	Activation
Cocking (1957) [30]	$0.536 \pm 0.006$	Activation
Rose (1958) [31]	$0.536 \pm 0.008$	Pile oscillator
Kappe (1966) [32]	$0.536 \pm 0.027$	Activation
Popovic (1955) [33]	$0.538 \pm 0.027^{b}$	Activation
Gleason (1975) [34]	$0.539 \pm 0.020$	Activation
Bartholomew (1953) [35]	$0.562 \pm 0.034^{b}$	Activation
Grimeland (1952) [36]	$0.57 \pm 0.04^{b}$	Activation
Kaminishi (1982) [37]	$0.577 \pm 0.008$	Activation
Coltman (1946) [38]	$0.60 \pm 0.05^{b}$	Transmission
Seren (1947) [39]	$0.65 \pm 0.13^{b}$	Activation
Average <sup>c</sup>	$0.527 \pm 0.020$	
Mughabghab (2006)	$0.517 \pm 0.004$	
This work	$0.541\pm0.003$	

<sup>a</sup>Measured assuming  $\sigma_0(^{24}\text{Na}^g = 530)$  mb.

<sup>b</sup>Corrected for standardization value from Mughabghab [40].

<sup>c</sup>Weighted average of all values with external error.

#### **IV. NUCLEAR STRUCTURE**

The level spins and parities shown in Table II are mainly derived from the <sup>24</sup>Na ENSDF Adopted Levels [9] and the (d, p) reactions assignments of Tomandl *et al.* [44]. We have revised some of these spin and parity assignments on the basis of newly assigned  $\gamma$ -ray transitions and the primary  $\gamma$ -ray reduced transition probabilities shown in Fig. 2. Transitions to levels with  $J^{\pi} = 1^-$ ,  $2^-$  are consistently about an order of magnitude stronger than all other transitions making this indicator a strong argument for these spin assignments. Although *s*-wave neutron capture on the <sup>23</sup>Na  $J^{\pi} = 3/2^+$ target can populate a capture state with an admixture of spins  $1^+$  and  $2^+$ , the dominant observed spin is  $1^+$ , predominantly due to a  $1^+$  resonance at 2.787 keV [40]. This is consistent with the weak feeding observed to  $3^-$  resonances. The admixture



FIG. 2. (Color online) Reduced transition probabilities,  $\sigma_{\gamma}/E^3$ , for thermal neutron capture primary  $\gamma$ -ray transitions in <sup>24</sup>Na. Transitions to levels with  $J^{\pi} = 1^{-}, 2^{-}$  are significantly stronger than transitions to other states. The lines connecting states of the same spin and parity are provided to guide the eye.

of  $J^{\pi} = 2^+$  to the capture state can be estimated from the average strength populating 3<sup>-</sup> levels which is  $\approx 3.5\%$  of that populating the 1<sup>-</sup>, 2<sup>-</sup> states. This population is larger than  $\approx 1.6\%$  estimated feeding the 2<sup>+</sup> bound state near  $\approx 176$  keV below the neutron separation energy, reported by Mughabghab [40], but consistent with contributions from the higher-lying 2<sup>+</sup> resonances at 35.295 and 52.973 keV.

3933.33 level. This level has no spin assignment in ENSDF and is loosely restricted to J = 1-4 by Tomandl. It decays only to levels with  $J^{\pi} = 2^+, 3^+$  and is weakly populated by a primary  $\gamma$  ray limiting likely spins to  $J^{\pi} = 1^+, 2^+, 3$ .

4048.58 level. This level is assigned as  $J^{\pi} = 0^{-}$  and is expected to be strongly populated by a primary transition, yet it is populated nearly two orders of magnitude more weakly than expected for a favored transition. Although this transition may be anomalous due to nuclear structure considerations, the spin assignment of this state should be reexamined.

4143.09 level. Assignment as  $(4^-)$  in ENSDF is inconsistent with  $\gamma$  rays deexciting this level to a  $2^+$  state. Tomandl assigns this level as  $J^{\pi} = 3^+ - 6^+$  and since no primary  $\gamma$  ray is observed feeding this level, which is weakly populated, a  $(4^+)$ spin assignment has been adopted.

4196.28 level. The spin proposed by Tomandl is  $J^{\pi} = 1^{-}, 2^{-}$ . The 1<sup>-</sup> spin is unlikely since this level deexcites to a 3<sup>+</sup> state. This level is only weakly populated by a primary  $\gamma$  ray, suggesting the parity is more likely positive.  $J^{\pi} = (2^{+})$  has been adopted for this level.

4206.991 level. This level is assigned  $J^{\pi} = 2^+$  in ENSDF and  $J^{\pi} = 1^-, 2^-$  by Tomandl. The  $\gamma$  ray deexciting this level to a  $3^+$  state eliminates the  $1^-$  assignment, and very strong population by a primary  $\gamma$  ray favors the  $2^-$  assignment.

4561.94 level. This level is assigned  $J^{\pi} = 1^{-}$  in ENSDF and  $J^{\pi} = 1^{+}, 2, 3^{+}$  by Tomandl. The  $\gamma$  ray deexciting this level to a 3<sup>+</sup> state eliminates the 1<sup>-</sup> assignment, and strong population by a primary  $\gamma$  ray favor a 2<sup>-</sup> assignment. 4908.022 level. This level is assigned  $J^{\pi} = 2^+ - 6^+$  by Tomandl. Weak population by a primary  $\gamma$  ray is most consistent with  $J^{\pi} = 2^+, 3$ 

5044.96, 5809.40 *levels*. These levels are assigned  $J^{\pi} = 1^{-}, 2^{-}$  by Tomandl. The  $1^{-}$  assignments are eliminated by transitions deexciting these levels to  $3^{+}$  states.

5117.34 level. This level is assigned  $J^{\pi} = 1^{-}$  in ENSDF and  $J^{\pi} = 1^{+}, 2^{+}$  by Tomandl. The  $\gamma$  ray deexciting this level to a  $3^{+}$  state eliminates the  $1^{-}$  assignment; however, the strong population by a primary  $\gamma$  ray favors a  $2^{-}$  assignment.

5308.6 *level*. This level is assigned  $J^{\pi} = 1^{-}-3^{-}$  by Tomandl.  $\gamma$  rays populating  $1^{+}$  and  $3^{+}$  eliminate the  $1^{-}$  and  $3^{-}$  assignments. Weak population by a primary  $\gamma$  ray favors a  $2(^{+})$  assignment.

5454.52 *level*. This level is assigned  $J^{\pi} = 1, 2, (3^+)$  by Tomandl. Strong population by a primary  $\gamma$  ray favors  $J^{\pi} = 1^-, 2^-$ .

5571.48 *level*. This level is assigned  $J^{\pi} = 2^+ - 4^+$  by Tomandl. Population by a weak primary  $\gamma$  ray is consistent with an assignment of  $J^{\pi} = (2^+)$ .

5775.5, 5851.32 levels. This level is assigned  $J^{\pi} = 1^{-}$ -3<sup>-</sup> by Tomandl. Very weak feeding by a primary  $\gamma$  ray is consistent with a (2<sup>+</sup>) assignment.

5862.85 level. No spin assignment has been previously recommended for this level. Weak population of this level by a primary  $\gamma$  ray and deexcitation to a 3<sup>-</sup> level suggest a (2<sup>+</sup>) assignment.

5918.21 level. This level is assigned  $J^{\pi} = 2^+, (1^+)$  by Tomandl. 1<sup>+</sup> is eliminated by a  $\gamma$  ray feeding a 3<sup>-</sup> level. While this would suggest the 2<sup>+</sup> assignment, this level is very strongly fed by a primary  $\gamma$  ray and has been tentatively assigned as (2<sup>-</sup>) in this work.

5953.39 *level*. No spin assignment has been previously recommended for this level. Strong population by a primary  $\gamma$  ray and deexcitation to a 0<sup>+</sup> level favors a (1<sup>-</sup>) assignment.

6072.611 level. This level is assigned  $J^{\pi} = 1^+$  in ENSDF and  $J^{\pi} = 2^+, (3^+)$  by Tomandl; however, it is very strongly fed by a primary  $\gamma$  ray suggesting negative parity. Depopulation to a 3<sup>+</sup> level suggests that (2<sup>-</sup>) is a more likely assignment.

6222.26 level. No spin assignment has been previously recommended for this level. Weak population by a primary  $\gamma$  ray suggests a  $(1^+, 2^+)$  assignment.

6251.29 level. This level is assigned  $J^{\pi} = 0^+, 1, 2$  by Tomandl. Strong feeding by a primary  $\gamma$  ray and deexcitation to a 3<sup>+</sup> level suggest a (2<sup>-</sup>) assignment.

6257.19 *level*. This level is assigned  $J^{\pi} = 1^{-}, 2^{-}$  by Tomandl. A  $\gamma$  ray to a 0<sup>+</sup> level eliminates the 2<sup>-</sup> possibility.

6406.62 *level*. New level assigned in this work. Strong feeding by a primary  $\gamma$  ray and a weak (M2)  $\gamma$  ray deexciting to a 4<sup>+</sup> level suggests a 2<sup>-</sup> assignment.

## **V. DISCUSSION**

A new, nearly complete <sup>23</sup>Na $(n,\gamma)^{24}$ Na thermal neutron capture  $\gamma$ -ray decay scheme has been measured in this work. Self-consistent total radiative neutron cross section,  $\sigma_0 = 0.541(3)$  b was determined from both prompt and delayed  $\gamma$ -ray measurements. This value is 4.4% higher than the recommended value of Mughabghab [40] but consistent with numerous earlier measurements. The neutron separation energy,  $S_n$ -6959.352(18) is slightly lower and more precise than the recently adopted value in the latest Atomic Mass Evaluation publication [43]. New spin and parity assignments have been made here on the basis of a more complete level scheme and the systematics of primary  $\gamma$ -ray transition probabilities.

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