

**Erratum: High-spin states in the five-valence-particle nucleus  $^{213}\text{Po}$  [Phys. Rev. C **83**, 034302 (2011)]**

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As input of the computation of the states drawn in Fig. 3(b) of this paper, we erroneously used the numerical values of the residual interactions of two neutrons in the  $g_{9/2}$  orbit instead of those of two protons in the  $h_{9/2}$  orbit. The corrected Fig. 3 is shown here.

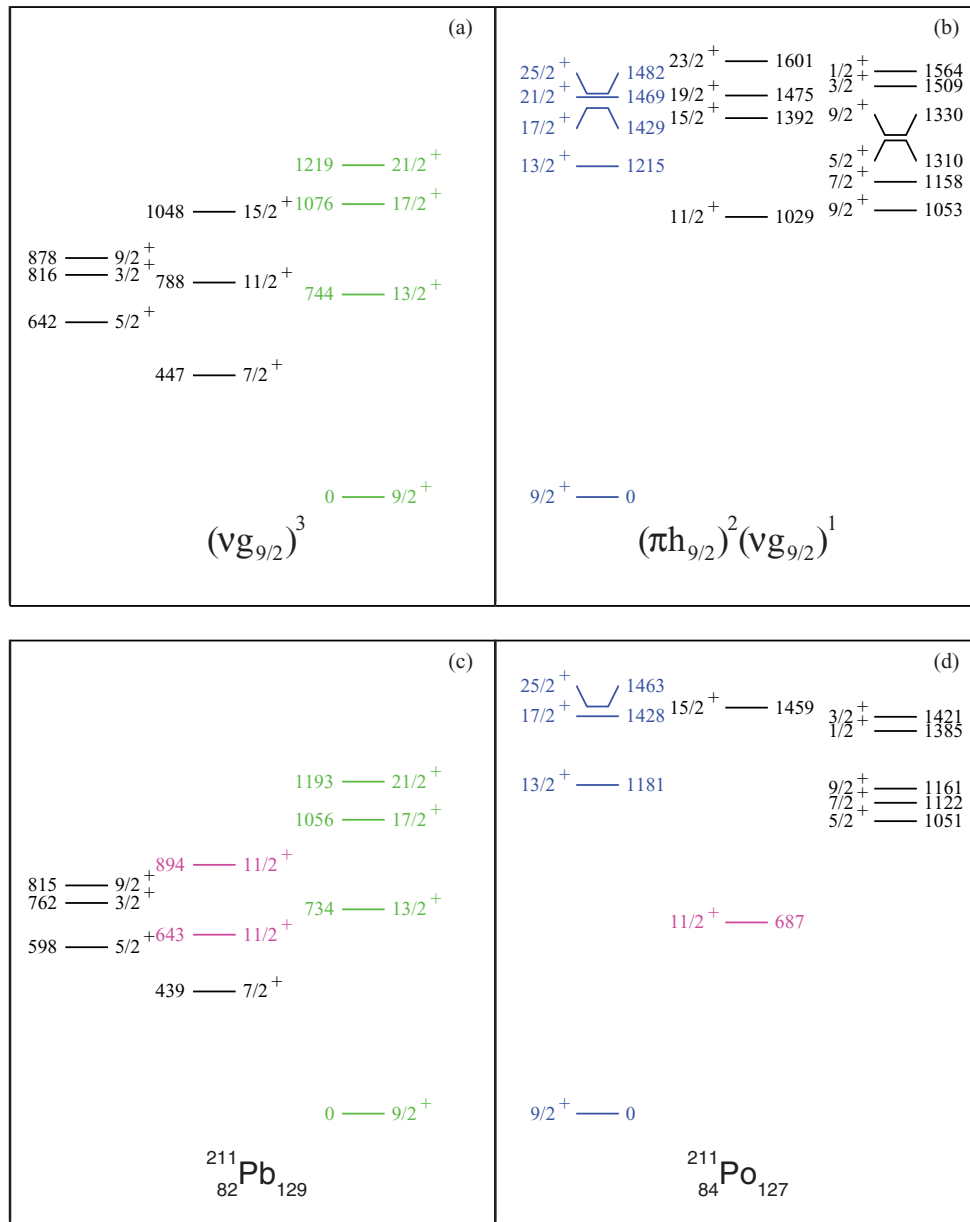


FIG. 3. (Color online) Top: Theoretical states belonging to (a) the spherical  $(vg_{9/2})^3$  configuration and (b) the  $(\pi h_{9/2})^2(vg_{9/2})^1$  configuration, computed using empirical interactions extracted from the two-nucleon neighboring nuclei ( $^{210}\text{Pb}$ ,  $^{210}\text{Po}$ ,  $^{210}\text{Bi}$ ). The yrast states of the two sets are drawn in green and blue, respectively. Bottom: First excited states measured (c) in the three-neutron nucleus,  $^{211}\text{Pb}$  [1,19], and (d) in the two-proton one-neutron nucleus,  $^{211}\text{Po}$  [20, 21]. The yrast states of the two nuclei are drawn in green and blue, respectively. The wave functions of the  $11/2^+$  states (drawn in magenta) contain a large component of the  $\nu i_{11/2}$  orbital located just above the  $vg_{9/2}$  one.

The agreement between the predictions of the  $(\pi h_{9/2})^2 \otimes (vg_{9/2})^1$  states [Fig. 3(b)] and the experimental results of  $^{211}\text{Po}$  [Fig. 3(d)] has obviously improved, with the yrast states (drawn in blue) being particularly well reproduced.

The conclusions of the paper are unchanged.

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