

Investigation of Branching Ratios Between Isomeric States of  $^{154}\text{Tb}$ 

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Half-lives and branching ratios of the three isomeric states of  $^{154}\text{Tb}$  have been measured from irradiation of gadolinium with 45- and 65-MeV protons. The spin assignments  $7^\pm$ ,  $3^\pm$ , and  $0^\pm$  previously suggested were confirmed as representing the  $m_2$ ,  $m_1$ , and  $g$  states, respectively. The branching from the  $m_2$  to  $m_1$  isomers was found to be weak, approximately  $1.8 \pm 0.6\%$ . Branching from the  $m_1$  to  $g$  state was quite substantial, measured to be  $21.8 \pm 0.7\%$ .

## INTRODUCTION

In the past three years, much research has been directed to studying the decay characteristics of the neutron-deficient isotopes of terbium, especially  $^{154}\text{Tb}$ ,<sup>1-7</sup> which are at the lower limit of the highly deformed region. The existence of three isomeric states of  $^{154}\text{Tb}$  was first suggested by Riedinger *et al.*<sup>4</sup> The spins and parities of these states are now thought to be  $7^-$ ,  $3^\pm$ , and  $0^{+4,6,7}$  with respective half-lives of 22.5, 9.0, and 21.8 h.<sup>2-4</sup> It is the purpose of the present work to confirm the existence and suggested ordering<sup>4,6</sup> of these three states, and to measure the branching between them by following their decay characteristics.

## EXPERIMENTAL

The target arrangement was identical to that used in previous nuclear reaction studies in the McGill Radiochemistry Laboratory.<sup>8</sup> Gadolinium oxide of natural abundance was mounted in a thin-walled aluminum tube and irradiated in the internal beam of the McGill University synchrocyclotron with protons of 45 to 65 MeV, for periods of 1.5 to 2 min. The target was slit open and the irradiated samples were deposited on a card suitable for counting. Counting began as early as two hours after irradiation making use of a 30-cm<sup>3</sup> Ge(Li) detector with associated 4096-channel pulse-height analyzer. Samples were counted at intervals of 0.5-3 h for as much as 120 h after irradiation. The intensity of individual peaks in each spectrum was determined after subtraction of the underlying Compton background. There were no problems with resolution of  $\gamma$  rays of interest in this work (see Fig. 1). Transitions of excited levels of  $^{154}\text{Gd}$  populated following  $\beta$  decay of the three isomeric states were measured. The work reported includes 11 measurements of the decay from the  $m_1$  to  $g$  states and 4 of the decay from the  $m_2$  to  $m_1$  state.

## RESULTS AND DISCUSSION

Four  $\gamma$ -ray transitions at 226.1, 540.1, 2064.1, and 2187.8 keV, depopulating levels in the daughter nucleus  $^{154}\text{Gd}$ , were followed for periods up to 120 h. Previous work<sup>6</sup> suggests that these transitions are populated *only* following decay of the  $m_2$ ,  $m_1$ ,  $g$ , and  $g$  states, respectively.

## A. Half-Life Measurements

Analysis of the decay curve of the 226-keV transition, one of the most intense peaks associated with decay of the  $7^-$  isomer,<sup>6</sup> shows a single component decay (Fig. 2). The statistically weighted average of a series of four measurements was found to be  $22.6 \pm 0.6$  h in excellent agreement with previous measurements<sup>5,6</sup> and with a somewhat reduced limit of error. There seems no question that the 226.1-keV transition arises following decay of the  $m_2$  state.

Measurement of the 540-keV transition, associated with decay of the  $3^\pm$  isomer exhibits the

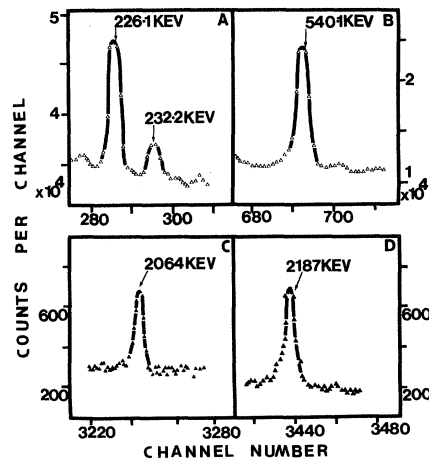


FIG. 1. Observed activities in 40 channels around peaks of interest in this work. (Scales on left and right are for A, C and B, D, respectively.)

characteristics of a long half-life parent feeding a short half-life daughter (Fig. 3).

The activity of each component was resolved, using the CLSQ decay curve analysis program<sup>9</sup> and the half-life measured for the  $m_2$  transition. Eight measurements, statistically weighted, yield a half-life of  $9.0 \pm 0.5$  h. Again, agreement with work of Sousa *et al.*<sup>6</sup> is excellent; however, we find a 10% shorter half-life than Ageev, Mitrovich, and Feoktistov<sup>3</sup> ( $9.9 \pm 0.1$  h) from their spallation work of 680-MeV proton bombardment of  $^{182}\text{Ta}$  followed by chemical separation of the rare earths. It is to be noted that the possibility that the 21.4-h component (rather than the 22.6-h  $m_2$  state) is the parent of the 540-keV transition is eliminated in the paragraph following.

Measurements of the 2064- and 2187-keV transitions (associated with the  $0^+$  isomer)<sup>6</sup> indicate that there is growth from a shorter lived component (see Fig. 4). This establishes that the  $0^+$  state must lie below the  $3^+$  in energy and cannot be the parent of the 9.0-h activity. The half-life of the  $0^+$  isomer determined by statistical weighting of 11 measurements was found to be  $21.4 \pm 0.5$  h, consistent with previous literature values.<sup>2,4</sup> It is noteworthy that this represents the first measurement of a statistically significant difference in the half-lives of the  $m_2$  and  $g$  states. Furthermore, the relative intensities of the 2064- and 2187-keV transitions, both exhibiting growth, are independent of proton bombarding energy while the 226-keV peak showing no growth varies with energy relative to the 2064- and 2187-keV peaks. Together these facts indicate absolute confirmation of two isomers with similar half-lives.

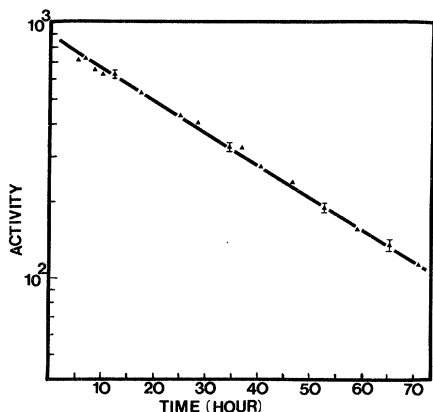


FIG. 2. Experimental decay curve for the peak at 226.1 keV (45-MeV bombarding energy; activities are measured in counts per minute).

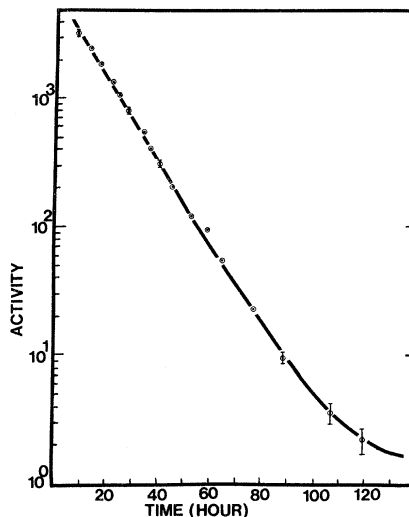


FIG. 3. Experimental decay curve for the peak at 540.1 keV (45-MeV bombarding energy; activities are measured in counts per minute).

#### B. Branching Ratios

The relative  $\gamma$ -ray intensities for transitions following the  $\beta$  decay of the three states are known from the comprehensive studies of Riedinger *et al.*<sup>4</sup> and Sousa *et al.*<sup>6</sup> The fraction of decays from one isomeric state to another has been determined by careful analysis of the growth and/or decay of these states.

From the 540-keV transition growth, one can

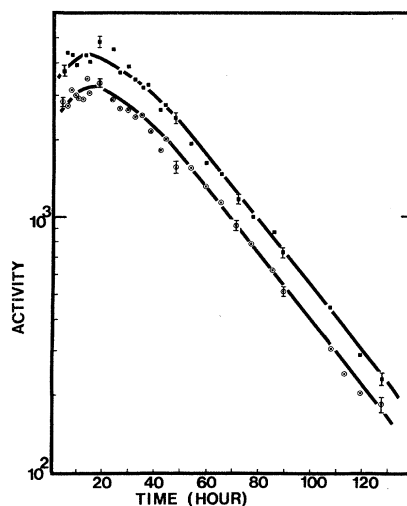


FIG. 4. Experimental decay curves for the peaks at 2187 keV (top curve) and 2064 keV (bottom curve). (45-MeV bombarding energy; activities are measured in counts per minute).

measure the apparent intensity of the  $m_2 \rightarrow m_1$  transition; this measurement, which is independent of the conversion coefficient for the  $m_2 \rightarrow m_1$  transition, may be compared to the intensity of the 226-keV transition, which represents a known<sup>6</sup> fraction of  $m_2 \rightarrow {}^{154}\text{Gd}$  transition. Four measurements of this branching fraction indicate that the decay of  $m_2 \rightarrow m_1$  accounts for  $1.8 \pm 0.6\%$  of all decays of  $m_2$ . This is, as suggested in Ref. 6, a branch of less than 2%.

The growth and decay of the peaks at 2064 and 2187 keV give a measure of the  $m_1$  to  $g$  transition. In conjunction with the data on decay of the 540-keV peak, the fraction of decays of  $m_1$  to the ground state may be calculated. Decay of the  $m_2$  state to  $m_1$  and thence from  $m_1$  to  $g$  is a second-order correction to decay of the 2064 and 2187 peaks but by a process of iteration, it has been made. Results of the analysis of 11 measurements show a branch to the ground state of  $21.8 \pm 0.7\%$ , larger than might be expected,<sup>6</sup> but not inconsistent with an  $E3$  transition.

### C. Discussion

It was previously suggested that the ground-state Nilsson configuration for  ${}^{154}\text{Tb}$  is  $\frac{3}{2}^+[411]_{\pi} + \frac{3}{2}^+[651]_{\nu} \rightarrow 0^+$ ,  $\Sigma=0$ , and for the  $m_2$  isomer is  $\frac{3}{2}^+[411]_{\pi} + \frac{11}{7}^-[505]_{\nu} \rightarrow 7^-$ ,  $\Sigma=1$ .<sup>4</sup> Because of the presence of two neutron states of differing parity,  $\frac{3}{2}^+[651]$  and  $\frac{3}{2}^-[521]$ , the parity of the spin-3  $m_1$  isomer was not known. However, many isomers attributed to the  $\frac{3}{2}^-[521]$  orbital have been observed among rare-

earth isotopes of the same region.<sup>10</sup> It is suggested, based upon the branching fractions measured in this work, that the parity of this state is negative, i.e., that the  $m_1$  configuration is likely to be  $\frac{3}{2}^+[411]_{\pi} + \frac{3}{2}^-[521]_{\nu} \rightarrow 3^-$  and the  $m_2 \rightarrow m_1$  transition is  $E4$ . With a 1.8% branch, the partial half-life for the  $m_2 \rightarrow m_1$  transition is  $4.52 \times 10^6$  sec (52.3 day); it is doubtful that an  $M4$  transition of reasonable energy could occur with such rapidity. Assuming the  $m_1$  is  $3^-$ , the partial half-life of the  $3^- \rightarrow 0^+$ ,  $m_1 \rightarrow g$  transition is  $1.48 \times 10^5$  sec (1.7 day), which is longer than normally expected for an  $E3$  transition of the order of 100 keV. One must conclude then that the ground and first excited states differ in energy by only a very small amount. It is significant in this respect that the neighboring isotones of  $N=89$ ,  ${}^{153}\text{Gd}$  and  ${}^{155}\text{Dy}$ , have ground states  $\frac{3}{2}^+$ ,<sup>11</sup> and  $\frac{3}{2}^-$ ,<sup>12</sup> respectively, illustrating the similarity in energies of the [651] and [521] Nilsson states.

### ACKNOWLEDGMENT

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