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Recent α decay half-lives and analytic expression predictions including superheavy nuclei

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New recent experimental α decay half-lives have been compared with the results obtained from previously proposed formulas depending only on the mass and charge numbers of the α emitter and the Q_α value. For the heaviest nuclei they are also compared with calculations using the Density-Dependent M3Y (DDM3Y) effective interaction and the Viola-Seaborg-Sobiczewski (VSS) formulas. The correct agreement allows us to make predictions for the α decay half-lives of other still unknown superheavy nuclei from these analytic formulas using the extrapolated Q_α of G. Audi, A. H. Wapstra, and C. Thibault [Nucl. Phys. **A729**, 337 (2003)].

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The α decay process was described in 1928 [1,2] in terms of a quantum tunneling through the potential barrier separating the mother nucleus energy and the total energy of the separated α particle and daughter nucleus. To describe the α emission two different approaches have been developed. The cluster-like theories suppose that the α particle is preformed in the nucleus with a certain preformation factor while the fission-like approaches consider that the α particle is formed progressively during the very asymmetric fission of the parent nucleus. The experimental investigation cannot unambiguously distinguish these two formation modes. However, the possible one-body configurations play a minor role because in the quasi-molecular decay path investigated in the α decay process the potential barrier is governed by the balance between the repulsive Coulomb forces and the attractive proximity forces and the Q_α value; consequently the barrier top is more external and lower than the pure Coulomb barrier and corresponds to two separated fragments. The difference between the two approaches appears mainly in the way the decay constant is determined. In the unified fission models [3,4] the decay constant λ is the product of the constant assault frequency ν_0 and the barrier penetrability P while in the preformed cluster models [5,6] a third factor is introduced: the cluster preformation probability P_0.

Before the theoretical explanation and description of the α decay process, Geiger and Nuttall [7] observed a dependence of the α decay partial half-life $T^{\text{expt}}_{1/2,\alpha}$ on the mean α particle range for a fixed radioactive family and Geiger-Nuttall plots are now an expression of $\log_{10}T_\alpha$ as a function of $ZQ^{-1/2}$, because different new relations have been proposed [4,8–12] to calculate $\log_{10}T_\alpha$ from the measured kinetic energy of the α particle via $E_\alpha = Q_\alpha A_d/(A_\alpha + A_d)$ or from Q_α given or extrapolated from mass formulas or tables.

Recently, isotopes of the elements 112, 113, 114, 115, 116, and 118 have been synthetized in fusion-evaporation reactions using ^{209}Bi, ^{233,238}U, ^{242,244}Pu, ^{243}Am, ^{245,248}Cm, and ^{249}Cf targets with ^{48}Ca and ^{70}Zn beams and observed via their α decay cascades [13–19]. These recent experimental results have led to new theoretical studies on α decay, for example, within the relativistic mean field theory [20], the Desity-Dependent M3Y (DDM3Y) interaction [21,22], the generalized liquid drop model (GLDM) [23,24], and the Skyrme-Hartree-Fock mean-field model [25]. The predicted half-lives against α decay of these transuranium nuclei obtained with a semiempirical formula taking into account the magic numbers have also been compared with the analytical superasymmetric fission model results and the universal curves and the experimental data [12].

In previous studies [4,26] both theoretical description and analytical formulas were presented for the α emission. Within a generalized liquid drop model including the proximity effects between the α particle and the daughter nucleus and adjusted to reproduce the experimental Q value, the α emission half-lives were deduced from the WKB barrier penetration probability as for a spontaneous asymmetric fission. The RMS deviation between the theoretical and experimental values of $\log_{10}T_\alpha$ was 0.63 for a data set of 373 emitters having an α branching ratio close to one and 0.35 for the subset of 131 even-even nuclides. A fitting procedure led to the following empirical formulas, respectively, for the 131 even(Z)-even(N), 106 even-odd, 86 odd-even, and 50 odd-odd nuclei. A and Z are the mass and charge numbers of the mother nucleus. The rms deviations are, respectively, 0.285, 0.39, 0.36, and 0.35.

$$\log_{10}[T_{1/2}(s)] = -25.31 - 1.1629A^{1/6}\sqrt{Z} + \frac{1.5864Z}{\sqrt{Q_\alpha}}, \quad (1)$$

$$\log_{10}[T_{1/2}(s)] = -26.65 - 1.0859A^{1/6}\sqrt{Z} + \frac{1.5848Z}{\sqrt{Q_\alpha}}, \quad (2)$$

$$\log_{10}[T_{1/2}(s)] = -25.68 - 1.1423A^{1/6}\sqrt{Z} + \frac{1.592Z}{\sqrt{Q_\alpha}}, \quad (3)$$

$$\log_{10}[T_{1/2}(s)] = -29.48 - 1.113A^{1/6}\sqrt{Z} + \frac{1.6971Z}{\sqrt{Q_\alpha}}. \quad (4)$$

Because new α decays have been observed and their partial α decay half-lives $T^{\text{expt}}_{1/2,\alpha}$ have been measured [15–19,25,27–32], they are compared in Tables I and II with the calculated

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TABLE I. Comparison between recently known experimental α decay half-lives and results obtained with formulas (1)–(4).

Nucleus	Q_α^{expt} (MeV)	$T_{1/2,\alpha}^{\text{expt}}$ (s)	$T_{1/2,\alpha}^{\text{Formulas}}$ (s)	Nucleus	Q_α^{expt} (MeV)	$T_{1/2,\alpha}^{\text{expt}}$ (s)	$T_{1/2,\alpha}^{\text{Formulas}}$ (s)
^{105}Te	4.900(0.050)	$0.70_{-0.17}^{+0.25} \times 10^{-6}$	$0.37_{-0.13}^{+0.21} \times 10^{-6}$	^{156}Er	3.486	2.3×10^{10}	1.6×10^{10}
^{158}Yb	4.172	4.3×10^{6}	2.7×10^{6}	^{160}Hf	4.902	1.9×10^{3}	1.9×10^{3}
^{174}Hf	2.497	6.3×10^{22}	4.55×10^{23}	^{158}W	$6.612_{-0.003}^{+0.003}$	$1.5_{-2}^{+2} \times 10^{-3}$	$1.21_{-0.03}^{+0.03} \times 10^{-3}$
^{168}W	4.507	1.6×10^{6}	3.1×10^{6}	^{162}Os	$6.767_{-0.003}^{+0.003}$	$1.9_{-2}^{+2} \times 10^{-3}$	$2.33_{-0.05}^{+0.06} \times 10^{-3}$
^{164}Os	6.475	4.2×10^{-2}	2.2×10^{-2}	^{166}Pt	7.286	3.0×10^{-4}	2.85×10^{-4}
^{168}Pt	6.997	2.0×10^{-3}	2.2×10^{-3}	^{170}Pt	6.708	1.4×10^{-2}	2.0×10^{-2}
^{172}Hg	7.525	4.2×10^{-4}	2.7×10^{-4}	^{174}Hg	7.233	2.1×10^{-3}	2.0×10^{-3}
^{188}Hg	4.705	5.3×10^{8}	2.0×10^{8}	^{178}Pb	7.790	2.3×10^{-4}	2.1×10^{-4}
^{180}Pb	7.415	5.0×10^{-3}	2.75×10^{-3}	^{184}Pb	6.774	6.1×10^{-1}	3.6×10^{-1}
^{186}Pb	6.470	1.2×10^{1}	4.7×10^{0}	^{188}Pb	6.109	2.7×10^{2}	1.3×10^{2}
^{190}Pb	5.697	1.8×10^{4}	8.7×10^{3}	^{192}Pb	5.221	3.6×10^{6}	2.1×10^{6}
^{194}Pb	4.738	9.8×10^{9}	1.3×10^{9}	^{188}Po	$8.087_{-0.025}^{+0.025}$	$4.0_{-1.5}^{+2.0} \times 10^{-4}$	$1.1_{-0.17}^{+0.19} \times 10^{-4}$
^{189}Po	$7.703_{-0.020}^{+0.020}$	$5.0_{-1}^{+1} \times 10^{-3}$	$3.0_{-0.4}^{+0.4} \times 10^{-3}$	^{190}Po	7.693	2.5×10^{-3}	1.5×10^{-3}
^{192}Po	$7.319_{-0.011}^{+0.011}$	$2.9_{-0.8}^{+1.5} \times 10^{-2}$	$2.2_{-0.2}^{+0.2} \times 10^{-2}$	^{210}Po	5.407	1.2×10^{7}	1.0×10^{6}
^{196}Rn	$7.616_{-0.009}^{+0.009}$	$4.4_{-0.9}^{+1.3} \times 10^{-3}$	$1.36_{-0.09}^{+0.09} \times 10^{-2}$	^{198}Rn	7.349	6.5×10^{-2}	9.56×10^{-2}
^{202}Ra	8.020	2.6×10^{-3}	3.63×10^{-3}	^{204}Ra	7.636	5.9×10^{-2}	5.5×10^{-2}
^{210}Th	8.053	1.7×10^{-2}	1.3×10^{-2}	^{212}Th	7.952	3.6×10^{-2}	2.4×10^{-2}
^{218}U	$8.773_{-0.009}^{+0.009}$	$5.1_{-1.0}^{+1.7} \times 10^{-4}$	$4.0_{-0.2}^{+0.2} \times 10^{-4}$	^{220}U	10.30	6.0×10^{-8}	5.8×10^{-8}
^{224}U	8.620	7.0×10^{-4}	8.2×10^{-4}	^{226}U	7.701	5.0×10^{-1}	5.67×10^{-1}
^{228}Pu	7.950	2.0×10^{-1}	5.13×10^{-1}	^{230}Pu	7.180	1.0×10^{2}	2.71×10^{2}
^{238}Cm	6.62	2.3×10^{5}	3.3×10^{5}	^{258}No	8.151	1.2×10^{2}	5.4×10^{1}
^{258}Rf	9.25	9.2×10^{-2}	1.03×10^{-1}	^{260}Rf	8.901	1.0×10^{0}	1.0×10^{0}
^{266}Hs	10.34	2.3×10^{-3}	2.1×10^{-3}	^{270}Hs	9.02	2.2×10^{1}	1.03×10^{1}
^{270}Ds	11.2	1.0×10^{-4}	6.7×10^{-4}	^{282}113	$10.63_{-0.08}^{+0.08}$	$7.3_{-2.9}^{+13.4} \times 10^{-2}$	$4.27_{-1.7}^{+2.8} \times 10^{-2}$

values from the above-mentioned formulas using the measured Q_α values. Table II displays also the results obtained with the DDM3Y effective interaction [21,22], the GLDM [4,23], and the Viola-Seaborg formulas with Sobiczewski constants [8,9].

A quite good agreement appears in Table I in the whole mass range confirming the accuracy of the formulas (1)–(4) and their usefulness for new predictions. Table II focuses on the heaviest elements for which the uncertainties both on the experimental Q value and α decay half-lives are larger because only some α cascades have been observed. The results obtained with the DDM3Y effective interaction agree with the experimental data as the ones calculated from the GLDM and largely better than the VSS calculations which give systematically longer half-lives. This shows that a GLDM taking account of the proximity effects, the mass asymmetry, and the quasi-molecular shapes is sufficient to reproduce the α decay potential barriers when the experimental Q_α value is known and proves that the double folding potential obtained using M3Y effective interaction supplemented by a zero-range potential for the single-nucleon exchange is also very appropriate to describe the α decay process. The DDM3Y results are on an average slightly larger than the experimental data while the GLDM values are slightly lower than the measured values. The values obtained using formulas (1)–(4) and, then, only A, Z, and Q_α are close to the values derived from the DDM3Y interaction and in agreement with the still rough experimental data. The fact that the partial α decay half-lives of these superheavy elements follow these simple formulas seems to prove that the experimental data are consistent with the formation of a cold and relatively compact composite nuclear system. The shell effects are implicitly contained in the Q_α value but difficult to disentangle.

Thus predictions of the partial α decay half-lives of still unknown superheavy nuclei within formulas (1)–(4) seem reliable and are displayed in Table III. The values obtained using the GLDM and the VSS expressions are also given for comparison. The assumed α decay energies are calculated from the atomic mass evaluation of Audi *et al.* [33] because the agreement with the experimental data on the mass of the known heaviest elements is very satisfactory. It may be useful for future experimental assignment and identification.

TABLE II. Comparison between recent experimental α decay half-lives and results obtained with the DDM3Y effective interaction [21,22], the GLDM [4,23], the formulas (1)–(4), and the VSS expressions [8,9].

Nucleus	$Q_\alpha^{\rm expt}$ (MeV)	$T_{1/2,\alpha}^{\rm expt}$	$T_{1/2,\alpha}^{\rm DDM3Y}$	$T_{1/2,\alpha}^{\rm GLDM}$	$T_{1/2,\alpha}^{\rm Formulas}$	$T_{1/2,\alpha}^{\rm VSS}$
$^{294}118$	11.81 ± 0.06	$1.8^{+75}_{-1.3}$ ms	$0.66^{+0.23}_{-0.18}$ ms	$0.15^{+0.05}_{-0.04}$ ms	$0.39^{+0.15}_{-0.11}$ ms	$0.64^{+0.24}_{-0.18}$ ms
$^{293}116$	10.67 ± 0.06	53^{+62}_{-19} ms	206^{+90}_{-61} ms	$22.81^{+10.22}_{-7.06}$ ms	308^{+136}_{-93} ms	1258^{+557}_{-384} ms
$^{292}116$	10.80 ± 0.07	18^{+16}_{-6} ms	39^{+20}_{-13} ms	$10.45^{+5.65}_{-3.45}$ ms	27^{+14}_{-9} ms	49^{+26}_{-16} ms
$^{291}116$	10.89 ± 0.07	$6.3^{+11.6}_{-2.5}$ ms	$60.4^{+30.2}_{-20.1}$ ms	$6.35^{+3.15}_{-2.08}$ ms	89^{+46}_{-30} ms	$336.4^{+173.1}_{-113.4}$ ms
$^{290}116$	11.00 ± 0.08	15^{+26}_{-6} ms	$13.4^{+7.7}_{-5.2}$ ms	$3.47^{+1.99}_{-1.26}$ ms	$8.9^{+5.4}_{-3.3}$ ms	$15.2^{+9.0}_{-5.6}$ ms
$^{288}115$	10.61 (6)	87^{+105}_{-30} ms	409 ms	$94.7^{+41.9}_{-28.9}$ ms	582^{+278}_{-187} ms	997^{+442}_{-303} ms
$^{287}115$	10.74 (9)	32^{+155}_{-14} ms	49 ms	$46.0^{+33.1}_{-19.1}$ ms	53^{+38}_{-22} ms	207^{+149}_{-85} ms
$^{289}114$	9.96 ± 0.06	$2.7^{+1.4}_{-0.7}$ s	$3.8^{+1.8}_{-1.2}$ s	$0.52^{+0.25}_{-0.17}$ s	$6.1^{+3.0}_{-2.0}$ s	$26.7^{+13.1}_{-8.7}$ s
$^{288}114$	10.09 ± 0.07	$0.8^{+0.32}_{-0.18}$ s	$0.67^{+0.37}_{-0.27}$ s	$0.22^{+0.12}_{-0.08}$ s	$0.52^{+0.30}_{-0.19}$ s	$0.98^{+0.56}_{-0.40}$ s
$^{287}114$	10.16 ± 0.06	$0.51^{+0.18}_{-0.10}$ s	$1.13^{+0.52}_{-0.40}$ s	$0.16^{+0.08}_{-0.05}$ s	$1.79^{+0.85}_{-0.57}$ s	$7.24^{+3.43}_{-2.61}$ s
$^{286}114$	10.35 ± 0.06	$0.16^{+0.07}_{-0.03}$ s	$0.14^{+0.06}_{-0.04}$ s	$0.05^{+0.02}_{-0.02}$ s	$0.11^{+0.05}_{-0.03}$ s	$0.19^{+0.08}_{-0.06}$ s
$^{284}113$	10.15 (6)	$0.48^{+0.58}_{-0.17}$ s	$1.55^{+0.72}_{-0.48}$ s	$0.43^{+0.21}_{-0.13}$ s	$2.4^{+1.2}_{-0.80}$ s	$4.13^{+1.94}_{-1.31}$ s
$^{283}113$	10.26 (9)	100^{+490}_{-45} ms	$201.6^{+164.9}_{-84.7}$ms	222^{+172}_{-96} ms	234^{+180}_{-100} ms	937^{+719}_{-402} ms
$^{285}112$	9.29 ± 0.06	34^{+17}_{-9} s	75^{+41}_{-26} s	$13.22^{+7.25}_{-4.64}$ s	127^{+69}_{-44} s	592^{+323}_{-207} s
$^{283}112$	9.67 ± 0.06	$4.0^{+1.3}_{-0.7}$ s	$5.9^{+2.9}_{-2.0}$ s	$0.95^{+0.48}_{-0.32}$ s	$9.6^{+4.9}_{-3.2}$ s	$41.3^{+20.9}_{-13.8}$ s
$^{280}111$	9.87 (6)	$3.6\ ^{+4.3}_{-1.3}$ s	$1.9^{+0.9}_{-0.6}$ s	$0.69^{+0.33}_{-0.23}$ s	$3.1^{+1.6}_{-1.05}$ s	$5.70^{+2.74}_{-1.84}$ s
$^{279}111$	10.52(16)	170^{+810}_{-80} ms	$9.6^{+14.8}_{-5.7}$ ms	$12.4^{+19.9}_{-7.6}$ ms	$10.9^{+17.8}_{-6.7}$ ms	$45.3^{+73.1}_{-27.6}$ms
$^{279}110$	9.84 ± 0.06	$0.18^{+0.05}_{-0.03}$ s	$0.40^{+0.18}_{-0.13}$ s	$0.08^{+0.04}_{-0.02}$ s	$0.65^{+0.31}_{-0.21}$ s	$2.92^{+1.4}_{-0.94}$ s
$^{276}109$	9.85 (6)	$0.72^{+0.87}_{-0.25}$ s	$0.45^{+0.23}_{-0.14}$ s	$0.19^{+0.08}_{-0.06}$ s	$0.65^{+0.33}_{-0.22}$ s	$1.44^{+0.68}_{-0.46}$ s
$^{275}109$	10.48 (9)	$9.7^{+46}_{-4.4}$ ms	$2.75^{+1.85}_{-1.09}$ ms	$4.0^{+2.8}_{-1.6}$ ms	$3.2^{+2.3}_{-1.3}$ ms	$13.7^{+9.6}_{-5.6}$ ms
$^{275}108$	9.44 ± 0.07	$0.15^{+0.27}_{-0.06}$ s	$1.09^{+0.73}_{-0.40}$ s	$0.27^{+0.16}_{-0.10}$ s	$1.9^{+1.2}_{-0.72}$ s	$8.98^{+5.49}_{-3.38}$ s
$^{272}107$	9.15 (6)	$9.8^{+11.7}_{-3.5}$ s	$10.1^{+5.4}_{-3.4}$ s	$5.12^{+3.19}_{-1.58}$ s	$17.6^{+10.2}_{-6.4}$ s	$33.8^{+17.9}_{-11.6}$ s
$^{271}106$	8.65 ± 0.08	$2.4^{+4.3}_{-1.0}$ min	$1.0^{+0.8}_{-0.5}$ min	$0.33^{+0.28}_{-0.16}$ min	$1.8^{+1.5}_{-0.8}$ min	$8.6^{+7.3}_{-3.9}$ min

TABLE III. Predicted α decay half-lives using the GLDM, the formulas (1)–(4), and the VSS formulas. The α decay energies are taken from the extrapolated data of Audi *et al*. [33].

$^{A}_{Z}$	Q	$T_{1/2}^{\rm GLDM}$	$T_{1/2}^{\rm form.}$	$T_{1/2}^{\rm VSS}$	$^{A}_{Z}$	Q	$T_{1/2}^{\rm GLDM}$	$T_{1/2}^{\rm form.}$	$T_{1/2}^{\rm VSS}$	$^{A}_{Z}$	Q	$T_{1/2}^{\rm GLDM}$	$T_{1/2}^{\rm form.}$	$T_{1/2}^{\rm VSS}$
$^{293}_{118}$	12.30	77 μs	187 μs	592 μs	$^{292}_{117}$	11.60	1.30 ms	6.47 ms	13.33 ms	$^{291}_{117}$	11.90	0.29 ms	0.32 ms	1.23 ms
$^{291}_{115}$	10.00	4.33 s	4.8 s	21.9 s	$^{290}_{115}$	10.30	0.62 s	4.2 s	6.86 s	$^{289}_{116}$	11.70	0.43 ms	1.05 ms	3.63 ms
$^{289}_{115}$	10.60	97.4 ms	113 ms	482 ms	$^{287}_{113}$	9.34	102 s	99.4 s	461 s	$^{286}_{113}$	9.68	9.44 s	61.5 s	92.5 s
$^{285}_{114}$	11.00	5.1 ms	12 ms	44.6 ms	$^{285}_{113}$	10.02	0.99 s	1.0 s	4.35 s	$^{284}_{112}$	9.30	64.7 s	25.1 s	47.3 s
$^{283}_{111}$	8.96	6.01 min	5.5 min	25.73 min	$^{282}_{112}$	9.96	0.772 s	0.297 s	0.516 s	$^{282}_{111}$	9.38	18.6 s	99.8 s	158.4 s
$^{281}_{112}$	10.28	0.102 s	0.2 s	0.786 s	$^{281}_{111}$	9.64	3.12 s	2.72 s	11.96 s	$^{281}_{110}$	8.96	3.05 min	4.6 min	22.47 min
$^{280}_{112}$	10.62	13.3 ms	25.4 ms	8.62 ms	$^{280}_{111}$	9.98	0.335 s	1.43 s	2.79 s	$^{279}_{112}$	10.96	2.06 ms	3.88 ms	14.1 ms
$^{279}_{109}$	8.70	10.35 min	7.72 min	36.32 min	$^{278}_{112}$	11.38	0.223 ms	0.083 ms	0.121 ms	$^{278}_{111}$	10.72	3.89 ms	12.5 ms	30.9 ms
$^{278}_{110}$	10.00	148.5 ms	51.8 ms	89.8 ms	$^{278}_{109}$	9.10	31 s	143 s	240 s	$^{277}_{112}$	11.62	0.069 ms	0.12 ms	0.402 ms
$^{277}_{111}$	11.18	0.323 ms	0.28 ms	1.073 ms	$^{277}_{110}$	10.30	23.1 ms	39 ms	162 ms	$^{277}_{109}$	9.50	1.89 s	1.48 s	6.61 s
$^{277}_{108}$	8.40	49.7 min	65.25 min	330.3 min	$^{276}_{111}$	11.32	0.157 ms	0.39 ms	1.11 ms	$^{276}_{110}$	10.60	4.03 ms	1.47 ms	2.35 ms
$^{276}_{108}$	8.80	131 s	40.6 s	75 s	$^{275}_{111}$	11.55	51.5 μs	42.3 μs	152 μs	$^{275}_{110}$	11.10	0.26 ms	0.43 ms	1.65 ms
$^{274}_{111}$	11.60	41.4 μs	88.1 μs	258 μs	$^{274}_{110}$	11.40	55.5 μs	19.5 μs	28.7 μs	$^{274}_{109}$	10.50	3.67 ms	9.84 ms	26.8 ms
$^{274}_{108}$	9.50	0.92 s	0.3 s	0.51 s	$^{274}_{107}$	8.50	9.94 min	48.45 min	70.98 min	$^{273}_{111}$	11.20	0.33 ms	0.29 ms	0.96 ms
$^{273}_{110}$	11.37	0.067 ms	0.11 ms	0.39 ms	$^{273}_{109}$	10.82	0.61 ms	0.5 ms	1.96 ms	$^{273}_{108}$	9.90	69.4 ms	101 ms	441.6 ms
$^{273}_{107}$	8.90	28.8 s	21.1 s	92.8 s	$^{272}_{110}$	10.76	1.97 ms	0.697 ms	0.94 ms	$^{272}_{109}$	10.60	2.34 ms	5.74 ms	15.02 ms
$^{272}_{108}$	10.10	21.7 ms	6.9 ms	10.9 ms	$^{272}_{106}$	8.30	24.9 min	6.38 min	11.4 min	$^{271}_{110}$	10.87	1.12 ms	1.79 ms	5.86 ms
$^{271}_{109}$	10.14	37.5 ms	29.9 ms	105.6 ms	$^{271}_{108}$	9.90	79.2 ms	109.7 ms	441.7 ms	$^{271}_{107}$	9.50	0.499 s	0.338 s	1.40 s

TABLE III. (*Continued*).

${}^{A}_{Z}$	Q	$T_{1/2}^{\mathrm{GLDM}}$	$T_{1/2}^{\mathrm{form.}}$	$T_{1/2}^{\mathrm{VSS}}$	${}^{A}_{Z}$	Q	$T_{1/2}^{\mathrm{GLDM}}$	$T_{1/2}^{\mathrm{form.}}$	$T_{1/2}^{\mathrm{VSS}}$	${}^{A}_{Z}$	Q	$T_{1/2}^{\mathrm{GLDM}}$	$T_{1/2}^{\mathrm{form.}}$	$T_{1/2}^{\mathrm{VSS}}$
${}^{270}_{110}$	11.20	0.199 ms	0.067 ms	0.083 ms	${}^{270}_{109}$	10.35	10.7 ms	30 ms	65 ms	${}^{270}_{108}$	9.30	4.48 s	1.4 s	2.02 s
${}^{270}_{107}$	9.30	2.0 s	6.25 s	11.9 s	${}^{270}_{106}$	9.10	3.59 s	0.99 s	1.66 s	${}^{270}_{105}$	8.20	24.38 min	94.58 min	140.53 min
${}^{269}_{109}$	10.53	3.75 ms	3.12 ms	10.25 ms	${}^{269}_{108}$	9.63	0.48 s	0.68 s	2.52 s	${}^{269}_{107}$	8.84	55.9 s	39 s	144.5 s
${}^{269}_{106}$	8.80	32.5 s	37.5 s	167.9 s	${}^{269}_{105}$	8.40	4.96 min	3.01 min	12.93 min	${}^{268}_{110}$	11.92	6.3 μs	1.84 μs	2.1 μs
${}^{268}_{109}$	10.73	1.28 ms	3.07 ms	7.15 ms	${}^{268}_{108}$	9.90	85.7 ms	28.6 ms	37.7 ms	${}^{268}_{107}$	9.08	9.86 s	35.4 s	55.5 s
${}^{268}_{106}$	8.40	12.1 min	3.4 min	5.1 min	${}^{268}_{105}$	8.20	25.4 min	102.7 min	140.5 min	${}^{268}_{104}$	8.10	23.8 min	5.88 min	10.2 min
${}^{267}_{110}$	12.28	1.3 μs	1.57 μs	4.4 μs	${}^{267}_{109}$	10.87	0.61 ms	0.49 ms	1.49 ms	${}^{267}_{108}$	10.12	22.1 ms	32.9 ms	112.5 ms
${}^{267}_{107}$	9.37	1.33 s	0.97s	3.36s	${}^{267}_{106}$	8.64	1.9 min	2.25 min	9.3 min	${}^{267}_{105}$	7.90	330 min	205 min	787 min
${}^{267}_{104}$	7.80	315 min	306 min	1494 min	${}^{266}_{109}$	10.996	0.32 ms	0.69 ms	1.63 ms	${}^{266}_{108}$	10.336	6.26 ms	2.16 ms	2.64 ms
${}^{266}_{107}$	9.55	0.41 s	1.21 s	2.21 s	${}^{266}_{105}$	8.19	29.0 min	121.8 min	152.5 min	${}^{266}_{104}$	7.50	81.47 h	20.09 h	31.30 h
${}^{265}_{109}$	11.07	0.223 ms	0.178 ms	0.498 ms	${}^{265}_{107}$	9.77	99.7 ms	74.4 ms	241 ms	${}^{265}_{105}$	8.49	2.70 min	1.76 min	6.43 min
${}^{265}_{104}$	7.78	6.58 h	6.58 h	29.65 h	${}^{264}_{107}$	9.97	29.9 ms	74.1 ms	151 ms	${}^{264}_{106}$	9.21	1.99 s	0.60 s	0.77 s
${}^{264}_{105}$	8.66	46.1s	154 s	232 s	${}^{264}_{104}$	8.14	19.2 min	5.03 min	7.36 min	${}^{263}_{108}$	10.67	1.03 ms	1.52 ms	4.45 ms
${}^{263}_{107}$	10.08	15.5 ms	11.6 ms	34.9 ms	${}^{263}_{105}$	9.01	3.65 s	2.4 s	8.27 s	${}^{263}_{104}$	8.49	72.7 s	76.8 s	324.7 s
${}^{262}_{107}$	10.30	4.42 ms	9.51 ms	20.5 ms	${}^{262}_{106}$	9.60	160.4 ms	47.5 ms	56.7 ms	${}^{262}_{105}$	9.01	4.06 s	10.9 s	18.2 s
${}^{262}_{104}$	8.49	82.6 s	20.6 s	27.9 s	${}^{261}_{107}$	10.56	1.04 ms	0.74 ms	2.07 ms	${}^{261}_{106}$	9.80	44.8 ms	56.1 ms	183.9 ms
${}^{261}_{105}$	9.22	0.96 s	0.60 s	1.92 s	${}^{260}_{107}$	10.47	1.77 ms	3.58 ms	7.62 ms	${}^{260}_{104}$	8.90	4.09 s	1.08 s	1.35 s
${}^{259}_{106}$	9.83	39.4 ms	50.5 ms	152.3 ms	${}^{259}_{105}$	9.62	69.0 ms	45.9 ms	136.7 ms	${}^{259}_{104}$	9.12	0.89 s	0.93 s	3.38 s
${}^{258}_{106}$	9.67	114 ms	36.1 ms	36 ms	${}^{258}_{105}$	9.48	0.18 s	0.42 s	0.74 ms	${}^{258}_{104}$	9.25	380 ms	103 ms	120 ms
${}^{257}_{105}$	9.23	1.0 s	0.67 s	1.8 s	${}^{257}_{104}$	9.04	1.66 s	1.76 s	5.88 s	${}^{256}_{105}$	9.46	230 ms	522 ms	848 ms
${}^{256}_{104}$	8.93	3.78 s	1.04 s	1.09 s	${}^{255}_{105}$	9.72	42.9 ms	28.9 ms	72.4 ms	${}^{255}_{104}$	9.058	1.57 s	1.69 s	5.19 s
${}^{254}_{104}$	9.38	181 ms	51.9 ms	50.5 ms	${}^{253}_{104}$	9.55	63.1 ms	68.3 ms	195.0 ms					

In conclusion, formulas already presented to determine the partial α decay half-lives have been checked on new experimental data in the whole mass range and the correct agreement allows us to make predictions for the partial α decay half-lives of still unknown superheavy nuclei.

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