

**$\alpha$  decay of  $^{105}\text{Te}$** 

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The  $\alpha$  decay of the neutron-deficient nuclide  $^{105}\text{Te}$  was observed. The  $^{50}\text{Cr}(^{58}\text{Ni},3n)$  reaction was used to produce  $^{105}\text{Te}$  nuclei. The  $^{105}\text{Te}$  residues were selected with the Argonne Fragment Mass Analyzer and implanted into a double-sided Si strip detector where their subsequent  $\alpha$  decay was detected. An  $\alpha$ -decay  $Q$  value of  $Q_\alpha = 4900(50)$  keV and a half life of  $T_{1/2} = 0.70(-0.17 + 0.25)\mu\text{s}$  were measured for  $^{105}\text{Te}$  and a reduced  $\alpha$ -decay width of  $W_\alpha = 3.3(-1.7 + 2.1)$  was deduced. The decay properties of  $^{105}\text{Te}$  are compared with those of heavier Te isotopes and theoretical predictions.

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The doubly magic, self-conjugate  $^{100}\text{Sn}$  nucleus and its neighbors have been a subject of intense research. It culminated in the production of a handful of  $^{100}\text{Sn}$  nuclei via fragmentation reactions [1,2]. Neutron-deficient nuclei near the  $N = Z$  line, just above the  $N = 50$  shell closure, exhibit large  $\alpha$ -decay branches [3]. In addition, the proximity of the proton drip line manifests itself by the presence of proton emitters [4]. The  $\alpha$ - and proton-decay  $Q$  values constrain mass models and the decay widths provide information about the structure of low-lying states. It has also been speculated that the occupation of the same orbitals by the valence nucleons might enhance the  $\alpha$ -particle preformation probability, leading to faster  $\alpha$  decays. The  $N = Z + 2$  ( $T_z = 1$ )  $\alpha$  emitters along the  $^{114}\text{Ba}$ - $^{110}\text{Xe}$ - $^{106}\text{Te}$  decay chain have been recently studied in Refs. [5,6]. The observation of  $\alpha$ -decay fine structure in  $^{107}\text{Te}$  yielded the energy separation between the  $d_{5/2}$  and  $g_{7/2}$  neutron single-particle states in  $^{103}\text{Sn}$  [7]. The same measurement could also be done for  $^{105}\text{Te}$ . This Rapid Communication reports the results of a search for the main  $\alpha$ -decay branch in  $^{105}\text{Te}$ .

A  $^{58}\text{Ni}$  beam from the Argonne Tandem Linac Accelerator System impinged on a  $^{50}\text{Cr}$  target to produce  $^{105}\text{Te}$  nuclei following the evaporation of three neutrons. Three beam energies, namely 224, 214, and 204 MeV, were used during the experiment to optimize the  $^{105}\text{Te}$  yield. Reaction products recoiling from the target were separated from the beam and dispersed according to their mass over charge state ratio in the Argonne Fragment Mass Analyzer (FMA) [8]. The calculated time of flight through the FMA for  $^{105}\text{Te}$  residues was about 650 ns. Mass slits were used to select only  $A = 105$  residues with charge states  $Q = 25^+$  and  $Q = 26^+$  ( $A = 106/Q = 25^+$  and  $A = 106/Q = 26^+$  residues were partially allowed as well). After passing through a position-sensitive parallel-grid avalanche counter at the focal plane of the FMA, the recoils were implanted into a double-sided Si strip detector (DSSD). The  $32 \times 32$  mm<sup>2</sup>, 60- $\mu\text{m}$ -thick DSSD was divided into 80 horizontal front strips and 80 vertical back strips,

forming 6400 pixels. Subsequent  $\alpha$  decays took place in the same pixel as the implantation. In the analysis, the implants were correlated with their subsequent  $\alpha$  decays using spatial and temporal relations.

Because of the expected short half-life of  $^{105}\text{Te}$ , special care was taken to optimize the detection of fast  $\alpha$ -decay events. First, delay-line amplifiers, which recovered within 0.6  $\mu\text{s}$  after the implantation signal, were used to detect  $\alpha$  decays. Second, only implant events followed by decay events within 8  $\mu\text{s}$  triggered the data acquisition system. This reduced the dead time drastically and allowed running with high rates in the DSSD.

Figure 1 shows the  $\alpha$  spectra collected during the experiment. The top panel contains decays correlated with  $A = 105$  recoils, whereas the bottom panel corresponds to  $A = 106$  residues. Because of the trigger used, only decays faster than 8  $\mu\text{s}$  were collected. A line containing 7 counts around 4 MeV is visible in the  $A = 106$  spectrum. The properties of these 7 events are consistent with the known  $^{106}\text{Te}$   $\alpha$  decay [5,6,9]. The  $A = 105$  spectrum contains a group of 13 counts just above 4.5 MeV. The wide energy distribution of these events is associated with the significant radiation damage caused by a large number of ions implanted in the DSSD. These events were interpreted as the  $\alpha$  decay from  $^{105}\text{Te}$ ; all other known  $A = 105$  nuclei produced in this reaction do not exhibit such a decay. Their time distribution is given in the inset. The observed yield for the  $^{105}\text{Te}$  events corresponds to a cross section  $\approx 10$  nb (for beam energies of 214 and 204 MeV; no events were observed at 224 MeV), if an FMA efficiency of 5% is assumed together with a factor of 4 loss because of the short half-life (see below).

Individual DSSD strips were gain matched using  $^{244}\text{Cm}$  and  $^{240}\text{Pu}$  long-lived  $\alpha$  emitters. To avoid systematic errors associated with implant-decay pileup, the  $^{105}\text{Te}$  energy was determined relative to the  $^{106}\text{Te}$  events. A value of 4128(9) keV from Ref. [9] was adopted for  $E_\alpha(^{106}\text{Te})$ . The  $^{105}\text{Te}$  half-life was obtained from the decay times of the

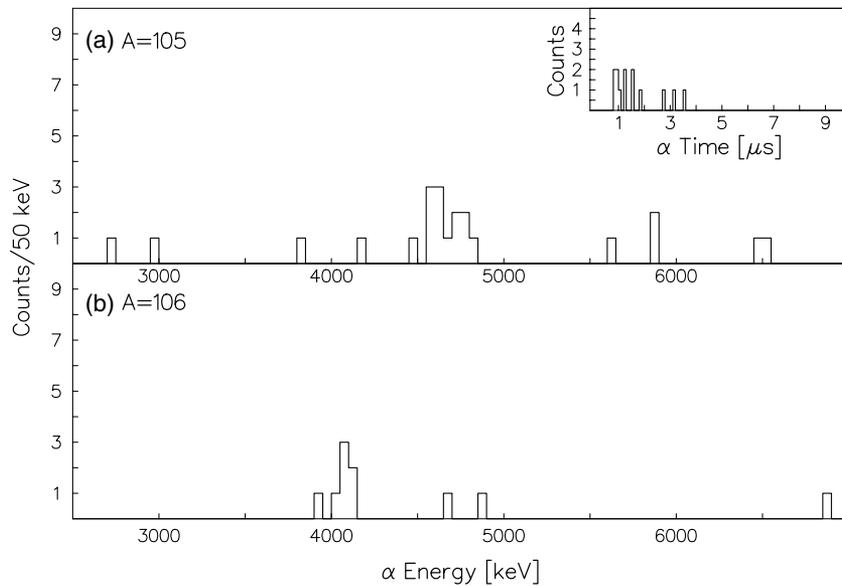


FIG. 1. Alpha spectrum for (a)  $A = 105$  and (b)  $A = 106$  residues. The inset in the top panel contains the decay time spectrum for the 13 events corresponding to the full energy  $^{105}\text{Te}$   $\alpha$  decay. Only decays with decay times longer than  $0.6 \mu\text{s}$  and shorter than  $8 \mu\text{s}$  were detected.

$^{105}\text{Te}$  events using the maximum likelihood method. As a result, an energy of  $E_\alpha = 4720(50)$  keV, corresponding to  $Q_\alpha = 4900(50)$  keV, and a half-life of  $T_{1/2} = 0.7(-0.17 + 0.25)\mu\text{s}$  were extracted. This is the shortest half-life ever observed using the implantation-decay correlation technique. Only statistical fluctuations were considered in error determination.

The  $\alpha$ -decay constant can be factored into the reduced width,  $\delta^2$ , and the barrier penetration factor,  $P$ , according to the formula  $\lambda b_\alpha = \delta^2 P/h$  [10], where  $b_\alpha$  is the  $\alpha$ -decay branching and  $h$  is Planck's constant. Assuming a 100%  $\alpha$  branch and calculating the penetration factor using the WKB approximation (a screening correction of 17 keV was added to the  $Q_\alpha$  value), a reduced width of  $\delta^2 = 230(-140 + 100)$  keV was obtained for  $^{105}\text{Te}$ , after taking into account the errors associated with the energy and half-life. The reduced width, relative to the 70-keV reduced width in  $^{212}\text{Po}$ , was

calculated to be  $W_\alpha = 3.3(-1.7 + 2.1)$ . This should be compared to  $W_\alpha = 4.6(-1.3 + 0.7)$  [6],  $W_\alpha = 1.46(0.64)$ , and  $W_\alpha = 2.57(0.24)$  for  $^{106}\text{Te}$ ,  $^{107}\text{Te}$ , and  $^{108}\text{Te}$ , respectively. The  $^{105}\text{Te}$  and  $^{106}\text{Te}$  widths appear to increase modestly compared to these of  $^{107}\text{Te}$  and  $^{108}\text{Te}$ , supporting the idea of an increasing  $\alpha$ -particle formation probability toward the  $N = Z$  line, although the errors are rather large.

Figure 2 contains  $Q_\alpha$  values for the neutron-deficient  $\alpha$  emitters above  $^{100}\text{Sn}$ . The  $Q_\alpha$  value for  $^{105}\text{Te}$  obtained in this work is larger by about 600 keV than the corresponding value for  $^{106}\text{Te}$ . It compares very well with a prediction of  $Q_\alpha(^{105}\text{Te}) = 4.69$  MeV from the Liran-Zeldes semiempirical formula [11], which is known to give good results far from the line of stability. A much larger value,  $Q_\alpha(^{105}\text{Te}) = 6.31$  MeV, was calculated using the FRDM model [12]. The  $Q_\alpha$  values increase by about 300 keV between  $^{109}\text{Te}$  and  $^{108}\text{Te}$  and between  $^{107}\text{Te}$  and  $^{106}\text{Te}$ . Thus, the extrapolation from  $^{105}\text{Te}$  to  $^{104}\text{Te}$  suggests a  $Q_\alpha$  value of about 5.2 MeV for the latter nucleus. If the same reduced width as in  $^{106}\text{Te}$  is assumed, a half-life of about 20 ns can be calculated for  $^{104}\text{Te}$ . Unfortunately, this implies that the direct observation of  $^{104}\text{Te}$ , produced in a fusion-evaporation reaction or via fragmentation, using a recoil separator is very difficult if not impossible, because of inherently long flight times. However, the observed  $^{105}\text{Te}$  yield makes the search for a 1%  $\alpha$  decay branch to the neutron  $g_{7/2}$  excited state in  $^{101}\text{Sn}$  possible.

In conclusion, the  $^{105}\text{Te}$   $\alpha$  decay was observed. It is the fastest  $\alpha$  emitter observed directly using the implantation-decay correlation technique. The  $Q_\alpha$  value measured for  $^{105}\text{Te}$  is reproduced well by the semiempirical formula of Liran and Zeldes [11]. The  $^{105}\text{Te}$   $\alpha$ -decay reduced width supports a modest enhancement of  $\alpha$ -decay rates toward the  $N = Z$  line.

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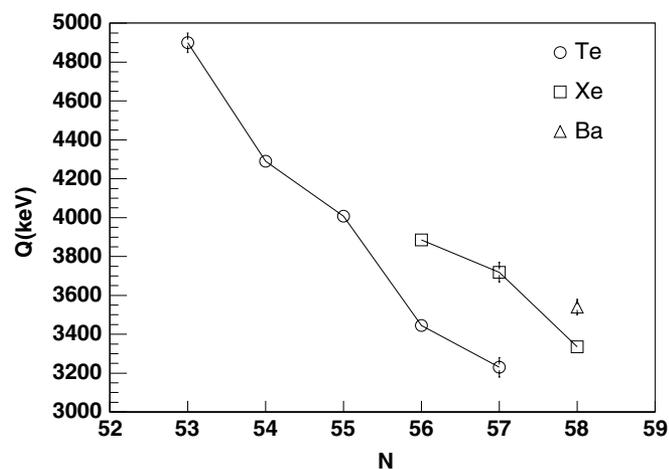


FIG. 2. Alpha-decay  $Q$  values for nuclei in the island of  $\alpha$  activity above  $^{100}\text{Sn}$ . The line connects the experimental points to guide the eye.

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