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<sup>1</sup>J. C. Hardy, J. E. Esterl, R. G. Sextro, and J. Cerny, *Phys. Rev. C* **3**, 700 (1971).

<sup>2</sup>J. C. Hardy, R. I. Verrall, R. Barton, and R. E. Bell, *Phys. Rev. Letters* **14**, 376 (1965).

<sup>3</sup>K. Gul, B. H. Armitage, and B. W. Hooton, *Nucl. Phys.* **A153**, 390 (1970); A. S. Clough, C. J. Batty, B. E. Bonner, and L. E. Williams, *Nucl. Phys.* **A143**, 385

(1970); and references therein.

<sup>4</sup>S. Cohen and D. Kurath, *Nucl. Phys.* **73**, 1 (1965); F. C. Barker, *Nucl. Phys.* **83**, 418 (1966).

<sup>5</sup>J. E. Esterl, R. G. Sextro, J. C. Hardy, G. J. Ehrhardt, and J. Cerny, *Nucl. Instr. Methods* **97**, 229 (1971).

<sup>6</sup>J. M. Mosher, R. W. Kavanagh, and T. A. Tombrello, *Phys. Rev. C* **3**, 438 (1971).

<sup>7</sup>Y. S. Chen, T. A. Tombrello, and R. W. Kavanagh, *Nucl. Phys.* **A146**, 136 (1970).

<sup>8</sup>D. H. Wilkinson, J. T. Sample, and D. E. Alburger, *Phys. Rev.* **146**, 662 (1966).

<sup>9</sup>J. D. Anderson, C. Wong, B. A. Pohl, and J. W. McClure, *Phys. Rev. C* **2**, 319 (1970); and R. J. Slobodrian, H. Bichsel, J. S. C. McKee, and W. F. Tivol, *Phys. Rev. Letters* **19**, 595 (1967).

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## Comment on a Possible $J^\pi = 0^+$ , $T = 2$ Resonance in $\text{Be}^9(\text{He}^3, \gamma\gamma)\text{C}^{12\dagger}$

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The reaction  $\text{Be}^9(\text{He}^3, \gamma\gamma)\text{C}^{12}$  has been reexamined near a previously reported resonance at  $E_{\text{He}^3} = 1.739 \pm 0.007$  MeV, which was ascribed to the lowest  $T = 2$  state in  $\text{C}^{12}$ . No resonance was observed and an upper limit  $\Gamma_{\text{He}^3\gamma\gamma}/\Gamma < 1.5$  meV is established for the  $T = 2$  resonance strength (assuming  $\Gamma \leq 1.5$  keV) which is  $\frac{1}{5}$  of the previously reported strength.

Several unsuccessful efforts have been made in recent years to observe the lowest  $T = 2$  level in  $\text{C}^{12}$  as an isospin-forbidden resonance in proton<sup>1</sup> and deuteron<sup>2</sup>-induced reactions. This level is known to have an excitation energy  $E_x = 27.595 \pm 0.020$  MeV from a  $\text{C}^{14}(p, t)\text{C}^{12}$  measurement.<sup>3, 4</sup> Recently Black, Caelli, and Watson<sup>2</sup> reported the observation of a strong candidate for this level as a resonance in the reaction  $\text{Be}^9(\text{He}^3, \gamma\gamma)\text{C}^{12}$  at an excitation energy of  $27.585 \pm 0.005$  MeV corresponding to a bombarding energy of  $1.739 \pm 0.007$  MeV. An upper limit of  $\Gamma < 1.5$  keV for the total width and a value for the capture strength of  $\Gamma_{\text{He}^3\gamma\gamma}/\Gamma = 8 \pm 5$  meV were given. We present the results of a reinvestigation of the same reaction in the region  $E_{\text{He}^3} = 1.721$  to 1.764 MeV, in which no resonance was observed.

In this experiment, thin metallic  $\text{Be}^9$  targets evaporated on polished Au backings were bombarded with the  $\text{He}^{3(+)}$  beam of the Brookhaven National Laboratory 3.5-MV Van de Graaff accelerator, and high-energy  $\gamma$  rays were detected in a  $10 \times 10$ -

in. NaI(Tl) detector at  $0^\circ$ . The accelerator beam analyzing magnet was calibrated by use of the resonance  $\text{Mg}^{24}(\alpha, \gamma)\text{Si}^{28}$  at  $E_\alpha = 3.1998 \pm 0.0010$  MeV,<sup>5</sup> the  $\text{C}^{13}(p, \gamma)\text{N}^{14}$  resonance at  $E_p = 1.7476 \pm 0.0009$  MeV,<sup>6</sup> and the  $\text{Be}^9(p, \gamma)\text{B}^{10}$  resonance at  $1.0832 \pm 0.0004$  MeV.<sup>7</sup> The internal consistency of the various calibrations was equivalent to  $\pm 1$  keV at  $E_{\text{He}^3} = 1.74$  MeV. To prevent energy shifts from target contamination, carbon buildup on the target surface was kept to a negligible level by the use of a liquid-nitrogen cold trap with a cold finger  $\sim 2$  mm from the target. The thicknesses of the thin targets were measured in two steps: First, the thickness of a  $33\text{-}\mu\text{g}/\text{cm}^2$   $\text{Be}^9$  target was determined from the observed width of the narrow  $\text{Be}^9(p, \gamma)\text{B}^{10}$  resonance at  $E_p = 1.083$  MeV; secondly, the thicknesses of the  $1.3\text{-}$  and  $3.2\text{-}\mu\text{g}/\text{cm}^2$  targets were obtained from a comparison of relative yields of the reaction  $\text{Be}^9(d, p)\text{Be}^{10}$ . The thicknesses of the latter two targets correspond to energy losses of 1.7 and 4.1 keV, respectively, for the  $\text{He}^3$  beam at 1.74 MeV.

The  $J^\pi = 0^+$ ,  $T=2$  level in  $C^{12}$  is expected to undergo  $\gamma$  decay strongly to the  $(1^+, 1)$  level at 15.1 MeV, which in turn decays predominantly to the ground state. Black, Caelli, and Watson<sup>2</sup> searched for resonances in the coincidence yield of two high-energy  $\gamma$  rays detected in two large NaI crystals placed  $180^\circ$  apart. The present experiment was designed to detect the  $\gamma$ -cascade coincidences as a summed peak at  $E_\gamma = 27.6$  MeV in a single  $10 \times 10$ -in. NaI detector placed with its front face from 6 to 10 mm from the target spot at  $\theta = 0^\circ$ . The crystal had a plastic anticoincidence shield for cosmic-ray rejection and lineshape improvement and was operated with antipileup electronics similar to previously described arrangements.<sup>3</sup> The sum-coincidence technique has the advantage of producing a high-energy signal which lies above the strong background below  $E_\gamma \sim 20$  MeV. This signal has the same energy as the (nonresonant) background from direct radiative capture,  $Be^9(He^3, \gamma_0)C^{12}$  to the ground state of  $C^{12}$ , which served as a useful monitor during the experiment.

The high-energy portion of a typical run taken near 1.74 MeV is shown in Fig. 1. The 27.6-MeV ground-state transition is clearly resolved, and the transition to the 4.44-MeV state is apparent as a shoulder. The area of the ground-state peak was obtained using a background subtraction indicated by the dashed lines.  $\gamma$ -ray spectra were recorded with both thin targets for bombarding energies ranging more than 2 standard deviations above and below the reported resonance energy

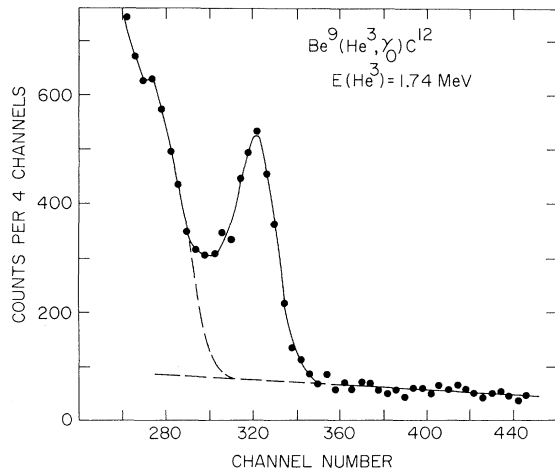


FIG. 1. High-energy portion of the  $\gamma$  spectrum from the reaction  $Be^9(He^3, \gamma)C^{12}$  at  $E_{He^3} = 1.74$  MeV observed in the  $10 \times 10$ -in. NaI(Tl) crystal. The peak near channel 320, which has an energy of 27.6 MeV, contains the ground-state transition and possible contributions from summed cascades.

of  $1.739 \pm 0.007$  MeV. The step size was 1.9 keV for the 4.1-keV target and 1.3 keV for the 1.7-keV target and runs were typically made for an accumulated charge of about 25 000  $\mu C$  with a beam current of about 4  $\mu A$ . The observed strength of the 27-MeV  $\gamma$  peak is plotted in Fig. 2 as a function of bombarding energy. The assigned errors contain statistical errors as well as relative uncertainties in the evaluation of the area. No resonance is apparent in either curve.

A comparison of this result with the published resonance strength requires knowledge of the detector efficiency  $\epsilon$ , which enters quadratically in the present measurement. The efficiency  $\epsilon$  contains solid angle, absorption, and electronic-acceptance-ratio<sup>8</sup> factors, and was determined experimentally by observing the yields of three known reactions which produce  $\gamma$  rays with energies comparable to those expected in the decay of the  $T=2$  state. The reaction  $C^{13}(p, \gamma_0)N^{14}$  was measured at the narrow resonance<sup>5</sup> at  $E_x = 9.17$  MeV, and the reaction<sup>9</sup>  $B^{11}(p, \gamma)C^{12}$  was measured at  $E_p = 1.42$  MeV. Also, the present  $Be^9(He^3, \gamma_0)C^{12}$  measurements were compared at  $E_{He^3} = 2.5$  MeV with absolute measurements made previously.<sup>10</sup> Care was taken in these comparisons to account properly for cascade-summing effects in the present measurements. These three measurements, which were internally consistent within  $\pm 5\%$ ,

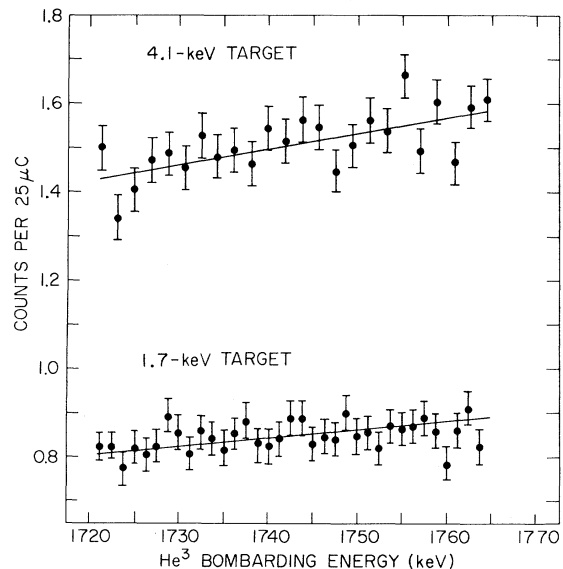


FIG. 2. Observed full-energy  $\gamma$  yield from the reaction  $Be^9(He^3, \gamma)C^{12}$  in the vicinity of the  $T=2$  state in  $C^{12}$  expected at  $E_{He^3} = 1739 \pm 7$  keV. The upper curve was measured using a 4.1-keV target and a step size of 1.9 keV, and the lower curve with a 1.7-keV target and a step size of 1.3 keV. The yields do not scale according to the thickness, because of slightly different solid angles.

yielded a photopeak efficiency  $\epsilon = 0.214$  ( $\pm 15\%$ ) at a detector distance of 6 mm, where the uncertainty is dominated by the absolute errors quoted for the calibration reactions. The same  $\epsilon$  was used for both  $\gamma$  rays in the cascade, since the efficiency is only weakly energy-dependent.

The previously reported capture strength of  $\Gamma_{\text{He}^3}\Gamma_\gamma/\Gamma = 8 \pm 5$  meV along with the  $0^+ \rightarrow 1^+ \rightarrow 0^+$  angular correlation and the measured efficiency quoted above lead to a predicted yield of  $0.95 \pm 0.60$  counts/25  $\mu\text{C}$  for an infinitely thick target, which is an order of magnitude greater than the apparent fluctuations in the lower curve of Fig. 2. The present data for  $E_{\text{He}^3}$  between 1.721 and 1.764 MeV result in an upper limit of  $\Gamma_{\text{He}^3}\Gamma_\gamma/\Gamma < 1.5$  meV corresponding to 2 standard deviations (95% confidence level), assuming  $\Gamma \leq 1.5$  keV as previously quoted.<sup>2</sup> For a less restrictive upper

limit,  $\Gamma \leq 10$  keV, the upper limit on the capture strength becomes 13 meV. The radiative width  $\Gamma_\gamma$  of this level is expected to be strong, based on the shell model. If  $\Gamma_\gamma$  is equal to 1 Weisskopf unit (41 eV), then the present results yield  $\Gamma_{\text{He}^3}/\Gamma < 3.6 \times 10^{-4}$  for  $\Gamma \leq 10$  keV. We note that these limits are dependent on the assumption of  $J^\pi = 0^+$  for the resonance, since they are dependent on the angular correlations of the two  $\gamma$  rays with each other and with the beam axis.

In summary, all attempts to detect the lowest  $T = 2$  level in  $\text{C}^{12}$  as an isospin-forbidden compound-nuclear resonance have failed. This is in contrast to the other known  $0^+$ ,  $T = 2$  states in light even-even self-conjugate nuclei, which all have an appreciable ground-state decay width in at least one of the energetically open (but isospin-forbidden) particle channels.

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<sup>1</sup>K. A. Snover, Ph.D. thesis, Stanford University, 1969 (unpublished).

<sup>2</sup>J. L. Black, W. J. Caelli, and R. B. Watson, Phys. Rev. Letters 25, 877 (1970).

<sup>3</sup>P. H. Nettles, C. A. Barnes, D. C. Hensley, and C. D. Goodman, Bull. Am. Phys. Soc. 16, 489 (1971).

<sup>4</sup>J. Cerny, Ann. Rev. Nucl. Sci. 18, 27 (1968).

<sup>5</sup>A. Rytz, H. H. Staub, H. Winkler, and F. Zamboni, Nucl. Phys. 43, 229 (1963).

<sup>6</sup>F. Ajzenberg-Selove, Nucl. Phys. A152, 1 (1970).

<sup>7</sup>T. Lauritsen and F. Ajzenberg-Selove, Nucl. Phys. 78, 1 (1968).

<sup>8</sup>E. M. Diener, J. F. Amann, S. L. Blatt, and P. Paul, Nucl. Instr. Methods 83, 115 (1970).

<sup>9</sup>R. E. Segel, S. S. Hanna, and R. G. Allas, Phys. Rev. 139, B818 (1965).

<sup>10</sup>S. L. Blatt *et al.* (private communication) report  $d\sigma/d\Omega(90^\circ) = 0.25$  ( $\pm 20\%$ )  $\mu\text{b}/\text{sr}$  at  $E_{\text{He}^3} = 2.5$  MeV, and  $W(\Theta) = 1 - 0.8P_2$  at  $E_{\text{He}^3} = 3.5$  MeV (assumed to hold for  $E_{\text{He}^3} = 2.5$  MeV) for  $\text{Be}^9(\text{He}^3, \gamma_0)\text{C}^{12}$ . The much smaller cross section reported in J. L. Black, G. A. Jones, and P. B. Treacy, Nucl. Phys. 54, 689 (1964), appears to be in error.