High-Spin States in ²²Na

J. N. Hallock, H. A. Enge, and A. Sperduto

Department of Physics and Laboratory for Nuclear Science,* Massachusetts Institute of Technology, Cambridge, Massachusetts 02139

and

R. Middleton, J. D. Garrett, † and H. T. Fortune Department of Physics, ‡ University of Pennsylvania, Philadelphia, Pennsylvania 19104 (Received 19 June 1972)

Based on their selective population in the ${}^{12}C + {}^{14}N \rightarrow \alpha + {}^{22}Na$ reactions, states at $E_x = 4.466$, 4.522, and 4.708 MeV are suggested, respectively, as the 4⁻ member of the $K^{\pi} = 1^{-}$ band, the 7⁺ member of the ground-state band, and the 5⁺ member of the $K^{\pi} = 0^+$, T = 0 band. A very strongly populated state at 8.60 MeV may be the 8⁺ member of the ground-state band.

The low-lying levels of ²²Na have been described with considerable success in terms of the collective model.¹⁻⁵ Several rotational bands have been identified from γ -ray decay studies,⁶ lifetime measurements,⁷ and single-^{4, 5} and two-nucleon³ transfer reactions. Spin and parity assignments now appear definite for all but two of the levels below 4.2 MeV, and J^{π} values have been suggested^{2,3} for those two states. Limits⁵ have been placed on the spins and parities of all known final states up to 5.7 MeV, except for those at 4.294, 4.466, 4.522, 5.099, and 5.117 MeV (see Table I). Since states of low spins $(J \leq 5)$ would be preferentially populated by the reactions that have been performed earlier, it is not surprising that only one state having J > 5 (the suggested² 6⁺ level at 3.708) has been identified.

Positive-parity states having $J^{\pi} \leq 4^+$ can be populated in a direct-stripping reaction on ²¹Ne $(J^{\pi} = \frac{3^+}{2})$ or in a direct-pickup reaction on ²³Na $(J^{\pi} = \frac{3^{+}}{2})$ by l=0 or 2, or both. An observed low yield for a particular positive-parity state in both reactions is thus strong evidence that the state has J > 4 or that the state is of an unusual configuration. Twoparticle transfer on ²⁰Ne can populate positiveparity states with $J \leq 5$ by depositing both nucleons into the *sd* shell and states with $J \leq 7$ by depositing both nucleons into the $f_{7/2}$ orbit. However, these $(f_{7/2})^2$ states are expected to lie quite high in ²²Na. Thus, a positive-parity state not observed to be directly populated in any of the three reactions either has a very unusual configuration or has J > 5. Similar arguments suggest J > 3 for negative-parity states not observed in pickup and stripping.

Recent studies of $({}^{12}C, \alpha)$ reactions^{8, 9} have shown that the states selectively populated in this reaction either have high spin^{8, 10} or possess α clusterlike configurations.⁹ We have, thus, used this reaction to search for the missing high-spin states that are expected to exist in ²²Na. In the present work, the states of ²²Na were investigated with the reaction ¹⁴N + ¹²C $\rightarrow \alpha$ + ²²Na at center-of-mass energies of 21.0, 18.5, 16.2, and 13.9 MeV. For the highest energy experiment, a ¹⁴N target was bombarded with a ¹²C beam at a laboratory energy of 39 MeV. For the other three runs, a ¹²C target was bombarded with a ¹⁴N beam of energies 40, 35, and 30 MeV. In both reactions, data were obtained at forward angles. The combined studies then yield both forward- and backward-angle data, albeit at somewhat different center-of-mass energies.

The ¹⁴N(¹²C, α) reaction was studied at a laboratory energy of 39 MeV using the University of Pennsylvania's tandem accelerator and multiangle spectrograph. The target was 8.5 Torr (12.6 $\mu g/$ cm²) of natural nitrogen contained in a gas cell with no entrance window.¹¹ An α -particle spectrum obtained at 7.5° is shown in Fig. 1. The energy loss¹² of approximately 42 keV for 39-MeV ¹²C ions in the target accounts for most of the observed width of the α groups. The 30-MeV ¹²C- $(^{14}N, \alpha)$ study was also performed at the University of Pennsylvania, using a $13 - \mu g/cm^2$ self-supporting natural C foil target. Angular distributions were measured in that study.¹³ Data from the $^{12}C(^{14}N, \alpha)$ reactions at 35 and 40 MeV were obtained using the Brookhaven National Laboratory's MP tandem accelerator and the Massachusetts Institute of Technology's multiple-gap spectrograph recently moved to Brookhaven National Laboratory (BNL). The target was a natural carbon foil of thickness 9 $\mu g/cm^2$, corresponding to energy losses¹² of approximately 54 and 51 keV for 35and 40-MeV incident ¹⁴N ions. A 35-MeV spectrum of this reaction obtained at a lab angle of

6

7.5° is shown in Fig. 2. Absolute differential cross sections measured at 7.5° in the 39- and 35-MeV runs are presented in Table I. Also tabulated are absolute differential cross sections at 3.75° and total cross sections σ_T (obtained by integrating from 0 to 90°) measured in the 30-MeV run. Expected uncertainties of 20, 35, and 50% are assigned to the absolute cross-section scales of the 39-, 30-, and 35-MeV runs, respectively.

The spectra shown in Figs. 1 and 2 are typical of those observed in α -particle-producing, heavyion reactions, where ¹²C is the target or incident particle.⁹ Moderately strong transitions are observed to selected low-lying levels; at higher excitation energies, very intense groups are superimposed on a background caused by breakup and by formation or unresolved levels. particle all have zero isospin, the ${}^{12}C + {}^{14}N \rightarrow \alpha$ + ${}^{22}Na$ reaction can populate T = 1 final states only through isospin mixing. Indeed, those known or suspected T = 1 states that could be resolved from nearby states are very weakly populated. These include states at excitation energies of 0.657, 4.069, 5.165, 6.715, and 6.834 MeV (level Nos. 6, 15, 28, 50, and 52, respectively).

In both spectra (Figs. 1 and 2), the low-lying members of the ground-state rotational band (ground state, 3^+ ; 0.891 MeV, 4^+ ; and 1.528 MeV, 5^+) are among the strongest transitions observed below 4 MeV. The 3.708-MeV state (level 13), which has been suggested² to be the 6^+ member of the ground-state band, is as strongly populated in the 35-MeV ${}^{12}C({}^{14}N, \alpha)$ spectrum (Fig. 2) as the lower-spin members of this rotational band. The trend is apparent by inspection of the integrated

Since the ground states of ^{12}C , ^{14}N , and the α

TABLE I. Cross sections for the ${}^{12}C + {}^{14}N \rightarrow \alpha + {}^{22}Na$ Reaction.

Level No.	${E_x}^a$ (MeV)	J ^{π a}	K^{π} a	¹⁴ N(¹² C, α) ^b 39 MeV, 7.5°	$d\sigma/d\Omega \text{ (mb/sr}^{12}\mathrm{C(}^{14}\mathrm{N}, \alpha)$ 35 MeV, 7.5°) ¹² C(¹⁴ N, α) 30 MeV, 3.75°	$\sigma_{0-90^{\circ}}^{c}$ (mb) $^{12}C(^{14}N, \alpha)$ 30 MeV
0	0.0	3+	3^+	0.111	0.439	0.507	0.806
1	0.583	1+	0+	0.029	0.181	0.091	0.134
2	0.657	$0^+, T = 1$	0+	0.002	0.013	0.057	0.084
3	0.891	4^{+}	3^{+}	0.170	0.424	0.205	0.949
4	1.528	5^{+}	3^{+}	0.198	0.388	1.14	2.079
5	1.937	1+	1+	10044	(((
6	1.952	$2^+, T = 1$	0+	0.044	{0.328	1.03	21.51
7	1.984	3^{+}	0+	0.099	l	l	
8	2.211	1-	1-	0.014	0.348	0.156	0.323
9	2.572	2-	1-	0.078	0.250	0.250	0.919
10	2.969	3^{+}	d	0.182	0.059	0.495	0.913
11	3.059	2+	(1+)	0.057	0.252	0.152	0.628
1 2	3.521	3-	1-	0.109	0.430	0.100	0.800
13	3.708	(6+)	(3+)	0.068	0.353	0.211	0.835
14	3.944	1+	(1 ⁺) ^d	0.027	0.033	0.228	0.269
15	4.069	$(4)^+$, $(T=1)$	(0+)	0.014	0.002 ^f	0.023	0.125
16	4.294			0.011	(((
17	4.319	1+		0.068	$\{0.022$	$\{0.402$	{0.913
18	4.360	$2^{+}(1^{+})$	$2^{+}(1^{+})$	0.015	0.027	0.010 ^f	0.063 ^f
19	4.466	(4 ⁻) ^e	(1 ⁻) e	0.136	(((
20	4.522	(7 ⁺) ^e	(3 ⁺) e	0.431	${2.38}$	0.575	${3.240}$
21	4.583	$2^{-}(0^{-}, 1^{-}, 3^{-})$	2-	(0.009 f	(,	(c
22	4.622	1		0.035	0.016 ^f	0.010^{f}	0.063 f
23	4.708	$(5^+)^{e}$	(0 ⁺) ^e	0.307	0.165	0.552	2.01
24	4.770	0+-4+	d	0.128	/	0.008 f	0.050 ^f
25	5.061	$1^+, 2^+$		0.067	0.013^{f}	0.008 f	0.050 f
26	5.099	•		(Ì	((
27	5.117			0.148	0.413	0.992	${2.14}$
28	5.165	$(2)^+$, $(T=1)$	2^{+}	0.035	0.005 f	0.008 f	0.050 ^f

^a Excitation energies, spin, parity, and band assignments (except those noted by a superscript) are from the literature as summarized in Ref. 5.

^b Reference 13.

^c Cross section integrated from 0-90° in the center-of-mass system.

^dSee Ref. 5.

^e Suggested from the present study.

 $^{\rm f}$ Upper limit for the cross section.

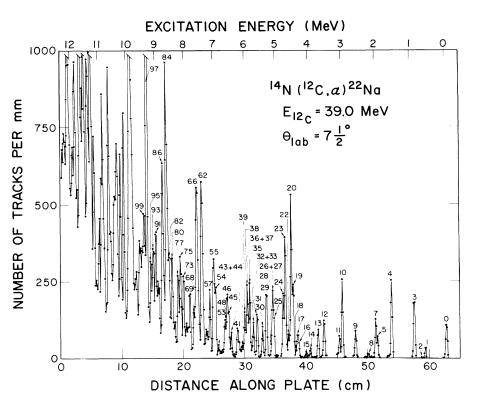


FIG. 1. α -particle spectrum from the ${}^{14}N({}^{12}C, \alpha){}^{22}Na$ reaction at a bombarding energy of 39 MeV. Excitation energies and differential cross sections corresponding to specific levels are listed in Table I. The level numbering scheme is the same as used in Refs. 3-5.

cross sections displayed in the last column of Table I. Even though there are a few exceptions, the general trend is that higher-spin states possess larger integrated cross sections. The cross section integrated from 0 to 90° for the 3.708-MeV level in the 30-MeV ¹²C(¹⁴N, α) study is larger than that for the other levels in this excitation region (Table I), even though this level is not selectively populated at the extreme forward angles.

By far the strongest transition observed below 7 MeV in both Figs. 1 and 2 is that to the 4.522-MeV level (level 20) of 22 Na. This state was only weakly populated by the single-^{4, 5} and two-nucleon³ transfer reactions (with no hint of a direct-like process) and was suggested⁵ as a possible high-spin state. It is an excellent candidate for the 7⁺ member of the ground-state rotational band.

For nuclei with A = 18-22, recent shell-model calculations¹⁴ in a full $1d_{5/2}-2s_{1/2}-1d_{3/2}$ basis give good agreement with most of the observed experimental properties of low-lying levels. In particular, predicted excitation energies of high-spin states appear¹⁴ to be quite reliable. Figure 3 contains the predicted excitation energies from these shell-model calculations for the lowest positiveparity state of each spin 3–9. Also plotted are the experimental excitation energies. The agreement between theory and experiment is seen to be excellent for the 3^+-6^+ members of the ground-state band. If the 4.522-MeV state is indeed the 7^+ member of this band, the agreement for that state is also quite good.

The next strongest state below 7 MeV in Fig. 1 and in the integrated cross sections is the state at $E_x = 4.708$ MeV (level 23). This level has been suggested from γ -decay studies² to be the 5⁺ member of the $K^{\pi}=0^+$, T=0 rotational band. This state was only weakly excited in the single-nucleon transfer reactions^{4, 5} and was populated by a probable L = 4 transfer in the ²⁰Ne(³He, *p*) reaction.³ All the evidence is consistent with the 4.708-MeV level's being the 5⁺ member of the $K^{\pi}=0^+$, T=0rotational band.

The 4.466-MeV state (level 19) also was weakly excited in one- and two-nucleon transfer and was suggested⁵ as a possible high-spin state. It is not resolved from the stronger 4.522-MeV level in the ${}^{12}C({}^{14}N, \alpha)$ runs, but it is populated with moderate strength in the ${}^{14}N({}^{12}C, \alpha)$ reaction. The 4⁻ member of the 1⁻ rotational band is expected² in this region of excitation. We tentatively identify this 4⁻ state as the experimental level at 4.466 MeV.

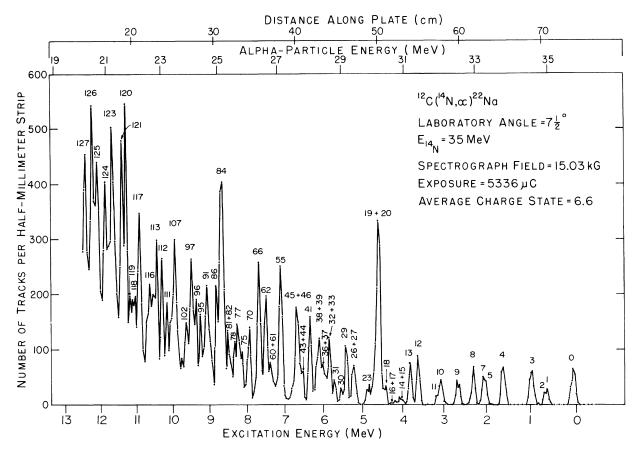


FIG. 2. α -particle spectrum from the ${}^{12}C({}^{14}N, \alpha){}^{22}Na$ reaction at a bombarding energy of 35 MeV. Excitation energies and differential cross sections corresponding to specific levels are listed in Table I. The level numbering scheme is the same as used in Refs. 3-5.

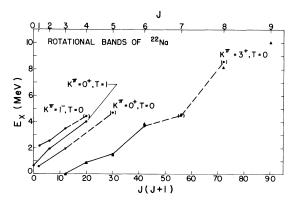


FIG. 3. Plot of excitation energies for states of spin J versus J(J+1) for the low-lying rotational bands of ²²Na. The identifications suggested in the present work are included in parentheses. The triangles indicate the predicted excitation (Ref. 14) of the lowest level of spin J. Other known rotational bands in ²²Na are summarized in Refs. 4 and 5.

The lowest 8^+ state of ²²Na is predicted¹⁴ at an excitation energy of 8.15 MeV, and a second 7⁺ level is predicted at 8.5 MeV. In the 35- and 40-MeV ¹²C(¹⁴N, α) exposures, as well as in the 39-MeV ¹⁴N(¹²C, α) exposure, the 3⁺, 4⁺, 5⁺, 6⁺, and suggested 7⁺ members of the ground-state rotational band are all strongly populated, and the transition to the 8.60-MeV level is the strongest observed below 9 MeV. (See Figs. 1 and 2.) Thus, it appears the 8.60-MeV state also has high spin and is a good candidate for the 8⁺ member of the ground-state band.

These tentative assignments should, of course, be confirmed by further study. More specific spectroscopic information may be obtained by studying the decay modes of these states. The ${}^{12}C + {}^{14}N \rightarrow \alpha + {}^{22}Na$ reaction should serve as a convenient means of selectively populating them for such a study.

Part of this work was performed at the BNL's tandem Van de Graaff facility, and we gratefully acknowledge the hospitality and cooperative spirit that were shown to us by Dr. Harvey Wegner, Dr. David Alburger, and the entire staff of the BNL tandem laboratory. The manuscript was edited and typed by Mrs. Mary E. White.

*Work supported in part through Atomic Energy Commission Contract No. AT(11-1)-3069.

†Present address: Los Alamos Scientific Laboratory, Los Alamos, New Mexico.

‡Work supported in part by the National Science Foundation.

¹E. K. Warburton, A. R. Poletti, and J. W. Olness, Phys. Rev. 168, 1232 (1960).

²J. W. Olness, W. R. Harris, P. Paul, and E. K. Warburton, Phys. Rev. C 1, 958 (1970)

³J. D. Garrett, R. Middleton, D. J. Pullen, S. A.

Andersen, O. Nathan, and O. Hansen, Nucl. Phys. A164, 449 (1971).

⁴J. D. Garrett, R. Middleton, and H. T. Fortune, Phys. Rev. C 4, 165 (1971).

⁵J. D. Garrett, H. T. Fortune, and R. Middleton, Phys. Rev. C 4, 1138 (1971).

⁶References 1 and 2, and references therein.

⁷K. W. Jones, A. Z. Schwarzschild, E. K. Warburton, and D. B. Fossan, Phys. Rev. 178, 1773 (1969), and references therein.

⁸R. Middleton, J. D. Garrett, and H. T. Fortune, Phys. Rev. Letters 24, 1436 (1970).

⁹R. Middleton, J. D. Garrett, H. T. Fortune, and R. R. Betts, J. Phys. (Paris) 32, C6 (1971).

¹⁰D. P. Balamuth, J. E. Holden, J. W. Noé, and R. W. Zurmühle, Phys. Rev. Letters 26, 1271 (1971); A. Gobbi, P. R. Maurenzig, L. Chua, R. Hadsell, P. D. Parker,

M. W. Sachs, D. Shapira, R. Stokstad, R. Wieband, and

D. A. Bromley, Phys. Rev. Letters 26, 396 (1971).

¹¹R. Middleton, in Proceedings of the International Conference on Nuclear Reactions Induced by Heavy Ions, Heidelberg, Germany, 15-18 July 1969 edited by R. Bock and R. Herring (North-Holland, Amsterdam, 1970), p. 263.

¹²L. C. Northeliffe and R. F. Schilling, Nucl. Data A7, 233 (1970).

¹³J. D. Garrett, Ph.D. thesis, University of Pennsylvania, 1970 (unpublished); J. N. Hallock, H. A. Enge, R. Middleton, J. D. Garrett, and H. T. Fortune, to be

published. ¹⁴E. C. Halbert, J. B. McGrory, B. H. Wildenthal, and S. P. Pandya, in Advances in Nuclear Physics, edited by M. Baranger and E. Vogt (Plenum, New York, 1971), Vol. IV; B. H. Wildenthal, private communication; and B. M. Preedom and W. H. Wildenthal, to be published.