## Comment on Rotational Bands in <sup>19</sup>F<sup>†</sup>

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A recent suggestion that the 4.385-MeV  $\frac{7}{2}^+$  state in <sup>19</sup>F is in the ground-state rotational band and that the 5.473-MeV  $\frac{7}{2}^+$  state is in an excited band appears to be contradicted by  $\gamma$ -decay data and by data from the <sup>16</sup>O(<sup>6</sup>Li, <sup>3</sup>He)<sup>19</sup>F reaction – both of which assign the 5.473-MeV state to the ground-state band.

Recently, Dixon *et al.*,<sup>1</sup> assigned  $J^{\pi} = \frac{9^{+}}{2}$  to a state at  $E_x = 6592$  keV in <sup>19</sup>F, and  $J = \frac{11}{2}$  (parity uncertain) to a state at  $E_x = 7937$  keV. They then used these assignments, together with the known<sup>2</sup> assignment of  $\frac{7}{2}$  for the state at 5464 keV to suggest that states at 3907 keV  $(\frac{3}{2}^+)$ , 4549 keV  $(\frac{5}{2}^+)$ , 5464 keV  $(\frac{7}{2}^+)$ , 6592 keV  $(\frac{9}{2}^+)$ , and 7937 keV  $(\frac{11}{2})$ form a  $K^{\pi} = \frac{3}{2}^{+}$  rotational band. There are no observed *E*2 transitions connecting these states,<sup>1</sup> but very strong E2's could be present without being detected. (A 20 W.u. in-band E2 would correspond<sup>1</sup> to less than a 1% decay branch for these states.) The principal evidence<sup>1</sup> that these states form a band is that their excitation energies lie on a straight line when plotted as a function of J(J+1). The identification of these states as members of an excited  $\frac{3^+}{2}$  rotational band appears to contradict existing  $\gamma$ -decay data,<sup>3</sup> as well as information on three-particle transfer.<sup>4</sup> We review the situation below and then present evidence which suggests a different band assignment.

The available information is as follows: (1) The 5464-keV  $\frac{7}{2}$  state decays<sup>3</sup> to the  $\frac{3}{2}$  state at 1554 keV with a strongly enhanced E2 [B(E2)= 41 ± 21  $e^2$  fm<sup>4</sup>]. This large B(E2) is suggestive of an in-band transition. The fact that the 1554keV state is in the ground-state band is then strong evidence that the 5464-keV state is also a member of the ground-state band.

(2) The configuration of the states in the groundstate band is largely<sup>5, 6</sup> (sd)<sup>3</sup>. The configuration of the  $\frac{3^+}{2}$  state at 3907 keV is most likely<sup>7, 8</sup> (sd)<sup>7</sup>(1p)<sup>-4</sup> or (sd)<sup>5</sup>(1p)<sup>-2</sup>. Recent shell-model calculations<sup>6</sup> of <sup>19</sup>F which include a full sd shell basis and an inert <sup>16</sup>O core predict two low-lying  $\frac{7^+}{2}$  states, but only one low-lying  $\frac{3^+}{2}$  state. Thus, both the lowlying  $\frac{7^+}{2}$  states (at 4378 and 5464 keV) can be accounted for as (sd)<sup>3</sup> states, whereas the  $\frac{3^+}{2}$  state at 3907 keV cannot be. Calculations<sup>7</sup> that assume a <sup>12</sup>C core, however, do reproduce this second  $\frac{3^+}{2}$ state, and suggest a 7p-4h or 5p-2h configuration for it. This would put the  $\frac{3^+}{2}$  3907-keV state in a band different from both of the low-lying  $\frac{7^+}{2}$  states. (3) The large B(E2) observed for the 5464-keV  $\frac{7^+}{2}$  decay to the first  $\frac{3^+}{2}$  state is accounted for by the  $(sd)^3$  calculations.<sup>6</sup> The calculated B(E2) value is 18.4  $e^2$  fm<sup>4</sup>. The B(E2) for the decay of the other low-lying  $\frac{7^+}{2}$  state to the first  $\frac{3^+}{2}$  state is calculated<sup>6</sup> to be small.

(4) In a recent study<sup>4</sup> of the <sup>16</sup>O(<sup>6</sup>Li, <sup>3</sup>He)<sup>19</sup>F reaction, the  $\frac{1^+}{2}$ ,  $\frac{3^+}{2}$ ,  $\frac{5^+}{2}$ ,  $\frac{9^+}{2}$ , and  $\frac{13^+}{2}$  members of the ground-state band (all of whose identifications now appear firm) were strongly populated relative to other states of the same  $J^{\pi}$ . Furthermore, the  $\frac{7^+}{2}$  state at 5464 keV was a factor of 10 stronger than the  $\frac{7}{2}$  state at 4378 keV – again suggesting that the 5464-keV state is the  $\frac{7^+}{2}$  member of the ground-state band.

Distorted-wave calculations have now been carried out for the <sup>16</sup>O(<sup>6</sup>Li, <sup>3</sup>He)<sup>19</sup>F reaction to all these low-lying positive-parity states. The details of the distorted-wave analysis and the results for the ground-state band have already been presented.9 Those results are summarized in Table I. Results are also presented for the lowest state of each  $J^{\pi} = \frac{1}{2}^+, \frac{3}{2}^+, \frac{5}{2}^+, \frac{7}{2}^+$  that does not belong to the ground-state band. It can be seen that the ratio of experimental to theoretical cross section is roughly constant for the  $\frac{1}{2}^+$ ,  $\frac{3}{2}^+$ ,  $\frac{5}{2}^+$ , and  $\frac{9}{2}^+$ members of the ground-state band. This is to be expected, since these states can be reasonably well represented as different J projections from a single intrinsic state of  $(sd)^3$  configuration.<sup>5, 10</sup> In a simple picture of the <sup>16</sup>O(<sup>6</sup>Li, <sup>3</sup>He)<sup>19</sup>F reaction on an inert <sup>16</sup>O core, these are the only states in <sup>19</sup>F that should be populated. In that simple picture, other states would be populated only through 2p-2h and 4p-4h admixtures in the <sup>16</sup>O groundstate wave function. The ratio  $\sigma_{exp}/\sigma_{DW}$  for states not within the ground-state band is significantly smaller than for states within the ground-state band, usually by an order of magnitude. Thus, the observed spectroscopic strengths for the  $\frac{7}{2}$ states, 0.27 for the 5.47-MeV state and 0.030 for

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$J^{\pi}$	Ground-state band members			Ground-state band nonmembers			
	E <sub>x.</sub> (MeV)	$\sigma_{max}$ (mb/sr)	$\frac{\sigma_{exp}}{\sigma_{DW}}$	$E_x$ (MeV)	$\sigma_{ m max}$ (mb/sr)	$\frac{\sigma_{exp}}{\sigma_{DW}}$	$\frac{\sigma_{\exp}/\sigma_{\rm DWBA}}{\sigma_{\exp}/\sigma_{\rm DWBA}}$ (for ground-state band nonmember)
$\frac{1}{2}^{+}$	0.00	0.106	0.35	5.34 <sup>a,b</sup>	≲0.005 <sup>b</sup>	≲0.030	≤0.085
$\frac{3+}{2}$	1,56	0.279	0.38	3.92	0.012	0.025	0.067
<u>5+</u> 2	0.20	0.682	0.33	4.55 <sup>c, d</sup>	≲0.059 <sup>d</sup>	≲0.063	≲0.19
$\frac{7+}{2}$	5.47	0.318 <sup>e</sup>	0.27	4.39	0.037	0.030	0.11
<u>9+</u> 2	2.78	1.16	0.32				

TABLE I. Comparison of cross sections in the <sup>16</sup>O(<sup>6</sup>Li, <sup>3</sup>He)<sup>19</sup>F reaction leading to states within and outside of the ground-state rotational band.

<sup>a</sup> $J^{\pi}$  has not yet unambiguously been assigned for 5.34-MeV state.

 $^{\rm b}$  The 5.34-MeV state is a member of a probable doublet. The combined cross section for the two members of the doublet is given.

<sup>c</sup> The parity of 4.55-MeV state is not yet completely unambiguous.

 $^{d}$  The 4.55- and 4.56-MeV states are not completely resolved. The cross section given is for both states.

<sup>e</sup> Further work on the <sup>16</sup>O(<sup>6</sup>Li, <sup>3</sup>He)<sup>19</sup>F reaction reveals that of the 0.347 mb/sr listed for the 5.47-MeV state in Ref. 4, 0.029 mb/sr arises from the 5.49-MeV state.

the 4.39-MeV state, appear to present a strong argument for placing the 5.47-MeV state in the ground-state band. It is observed, however, that the strength of the 5.47-MeV state is somewhat smaller than the strengths of the other members of the ground-state band. This result would indicate more mixing between the  $\frac{7^+}{2}$  states than for states of other  $J^{\pi}$ . This is consistent with theoretical expectations.<sup>5, 8, 10</sup>

These results, coupled with the lifetime mea-

surements,<sup>3</sup> then lead to the conclusion that the 5.47-MeV state (and not the 4.39-MeV state) is to be identified with the  $\frac{7^+}{2}$  member of the ground-state band.

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