## Comment on ${}^{12}C({}^{12}C, \alpha)^{20}$ Ne Excitation Functions\*

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The excitation function for the  ${}^{12}C({}^{12}C, \alpha)^{20}$ Ne reaction leading to the 2<sup>+</sup> state at 7.83 MeV possesses oscillations similar in width to those observed for the low-lying states. For the 7.83-MeV state, the oscillations are considerably damped in magnitude and are superimposed on a large underlying cross section that varies much more slowly with energy.

In a recent study<sup>1</sup> of the <sup>12</sup>C(<sup>12</sup>C,  $\alpha$ )<sup>20</sup>Ne reaction, certain states were observed to be selectively populated. These included the 0<sup>+</sup> state at 7.20 MeV and the 2<sup>+</sup> state at 7.83 MeV. Angular distributions for those states were similar to what one might expect for a direct mechanism. On the other hand, excitation functions showed strong oscillations. Exhaustive studies of angular distributions and excitation functions for the reaction  ${}^{12}C({}^{12}C, \alpha){}^{20}Ne$ leading to the low-lying states of <sup>20</sup>Ne had been performed previously.<sup>2</sup> Those data were successfully interpreted in terms of a compound-nucleus reaction mechanism. The question naturally arises: Are the excitation functions for the selectively populated states qualitatively different from those of the low-lying states?

Previous data<sup>1,3</sup> indicated that there were gualitative differences, e.g., the prominent structure in the excitation function measured for the 7.83-MeV state appeared to be much broader than that measured for the low-lying states. These results, however, were not conclusive, since data<sup>1,3</sup> were taken only in 500-keV (lab) steps. We have investigated the excitation function for this state in smaller steps [50 keV (lab)], in order to make a direct comparison with the low-lying states.

In Ref. 1, the most prominent structure in the excitation function for the 7.83-MeV state was a relatively broad peak near 26.5-MeV bombarding energy. This peak appeared to have a width of about 2 MeV, whereas the structure observed earlier<sup>2</sup> for the low-lying states was about 100-200keV in width. We have measured excitation functions in the bombarding energy range 25-27 MeV in 50-keV steps for the 7.83-MeV 2<sup>+</sup> state and for the low-lying states.

The data were obtained with the  ${}^{12}C^{4+}$  beam from the University of Pennsylvania tandem. Targets were  $7-\mu g/cm^2$  self-supporting foils of isotopically enriched (99.99%) <sup>12</sup>C. This thickness corresponds to an energy loss of  $\sim 30$  keV for 25-MeV <sup>12</sup>C ions. Outgoing  $\alpha$  particles were detected in a Si surfacebarrier detector of depletion depth 1000  $\mu$ . A Ta

foil placed in front of the detector stopped the <sup>12</sup>C ions without seriously affecting the energy resolution for the  $\alpha$  particles. Tests for carbon buildup during the course of the experiment were made by frequent measurements of repeat points. A small buildup was observed and has been corrected for.

An  $\alpha$ -particle spectrum obtained at a bombarding energy of 25.0 MeV and a lab angle of  $20^{\circ}$  is displayed in Fig. 1. Particle groups of interest are those leading to the  $0^+$  ground state (g.s.), the 2<sup>+</sup> 1.63-MeV, 4<sup>+</sup> 4.25-MeV, 2<sup>-</sup> 4.97-MeV, and 2<sup>+</sup> 7.83-MeV states.

Excitation functions measured at  $20^{\circ}$  (lab) for those states are displayed in Fig. 2. Also displayed are the earlier data<sup>3</sup> for the 7.83-MeV state measured at  $\theta(lab) = 18.75^\circ$ , in 500 keV steps. The general features of the present data are similar to those previously<sup>2</sup> measured for the ground and 1.63-MeV states. Rapid oscillations are observed, with peak-to-valley ratios of 10:1, and widths of ~250 keV. Structure of similar width is observed in the data for the 7.83-MeV state, but the peak-tovalley ratios are considerably smaller.

The average cross section measured in this energy interval for the 7.83-MeV state is 9.1 mb/sr, whereas for the 1.63-MeV 2<sup>+</sup> state, the average



FIG. 1.  $\alpha$  spectrum obtained at 20° for the <sup>12</sup>C- $({}^{12}C, \alpha)^{20}$ Ne reaction at 25.0 MeV.

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FIG. 2. Excitation functions measured at 20° (lab) for the  ${}^{12}C({}^{12}C,\alpha){}^{20}Ne$  reaction at low-lying states.

cross section is only 3.7 mb/sr. Calculations<sup>4</sup> based on a purely statistical process predict that the 7.83-MeV 2<sup>+</sup> state will be only  $\frac{1}{4}$  as strong as the 1.63-MeV 2<sup>+</sup> state.

It thus appears that the rapid oscillations observed for the low-lying states is also present in the 7.83-MeV data. However, for the 7.83-MeV state, these oscillations appear to be superimposed on a large background having a width much larger than that of the rapid oscillations. There is no indication of such an underlying background for the low-lying states – neither in the present data nor in the much more exhaustive data of Ref. 2.

The excitation function data for the 7.83-MeV state is thus qualitatively different from that of the other states. What, then, is the reaction mechanism responsible for population of the 7.83-MeV state? It was suggested in Ref. 1 that a large direct component was present, in which two  $\alpha$  particles were transferred to the <sup>12</sup>C target, and thus that the 7.83-MeV state had a large parentage of a <sup>12</sup>C core plus two  $\alpha$  particles.

Noble<sup>5</sup> has calculated the excitation functions expected from such a reaction. He finds that inclusion of the momentum transfer dependence at both vertices ( $^{12}C \rightarrow ^{**8}Be'' + \alpha$  and  $^{12}C + ^{**8}Be'' \rightarrow ^{20}Ne$ ) produces oscillations in the excitation functions with a width of a few MeV. It thus appears that, at least for the 7.83-MeV state, the data are consistent with the presence of such a reaction mechanism, but with the usual compound-nucleus oscillations superimposed on this direct background. Measurements to further elucidate the  $^{12}C(^{12}C, \alpha)$ - $^{20}Ne$  reaction mechanism are underway at this laboratory.

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<sup>3</sup>R. Middleton, J. D. Garrett, H. T. Fortune, and R. R. Betts (unpublished).

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<sup>5</sup>J. V. Noble, Phys. Rev. Letters 28, 111 (1972).