

Low-energy ${}^7\text{Li}(n, \gamma_0){}^8\text{Li}$ and ${}^7\text{Li}(p, \gamma_0){}^8\text{Be}$ cross sections

F. C. Barker

Department of Theoretical Physics, Research School of Physical Sciences and Engineering, The Australian National University, Canberra, ACT 0200, Australia

(Received 3 September 1996)

The lack of p -wave strength in a recent measurement of the low-energy ${}^7\text{Li}(n, \gamma_0){}^8\text{Li}$ cross section, as compared with significant strength found in the ${}^7\text{Li}(p, \gamma_0){}^8\text{Be}$ reaction, may be attributed to the qualitatively different low-energy behaviors of penetration factors for neutron and charged-particle channels. [S0556-2813(97)05001-2]

PACS number(s): 25.40.Lw, 27.20.+n

Blackmon *et al.* [1] recently measured the ${}^7\text{Li}(n, \gamma_0){}^8\text{Li}$ cross section for $E_n = 1.5 - 1340$ eV, in part to look for evidence of p -wave contributions. They found none. Earlier a ${}^7\text{Li}(p, \gamma_0){}^8\text{Be}$ measurement by Chasteler *et al.* [2] at $E_p = 70$ keV had indicated substantial p -wave strength, of uncertain origin. Chasteler *et al.* pointed out that such strength could significantly reduce the zero-energy astrophysical S factor obtained by assuming pure s waves, and that a similar phenomenon might be present in the ${}^7\text{Be}(p, \gamma){}^8\text{B}$ reaction, which is of importance in the solar neutrino problem and to which ${}^7\text{Li}(n, \gamma_0){}^8\text{Li}$ is the mirror reaction. Here we show that the difference in the observed low-energy p -wave strengths may be attributed to the qualitatively different low-energy behaviors of penetration factors for neutron and charged-particle channels.

The ℓ -wave contributions to the low-energy cross sections for the three reactions ${}^7\text{Li}(n, \gamma_0){}^8\text{Li}$, ${}^7\text{Li}(p, \gamma_0){}^8\text{Be}$, and ${}^7\text{Be}(p, \gamma){}^8\text{B}$ are proportional to the penetration factors P_ℓ in the initial channels ${}^7\text{Li} + n$, ${}^7\text{Li} + p$ and ${}^7\text{Be} + p$, respectively. We are particularly interested in the ratio P_1/P_0 . For the charged-particle channels, the low-energy penetration factors are determined mainly by the Coulomb potential, which extends to larger distances than either the centrifugal potential or the nuclear interaction. Since the Coulomb potential is ℓ independent, the ratio of the P_ℓ for different ℓ values tends to a constant as the energy tends to zero. With the penetration factor defined as in R -matrix theory [3],

$$P_\ell = \rho / (F_\ell^2 + G_\ell^2),$$

where F_ℓ and G_ℓ are the regular and irregular Coulomb functions evaluated at the channel radius a , and $\rho = ka$, the formulas in Ref. [4] lead to

$$[P_\ell / P_{\ell-1}]_{E=0} = [K_{2\ell-1}(x) / K_{2\ell+1}(x)]^2,$$

where $K_n(x)$ is a Bessel function [5] and $x = (8\rho\eta)^{1/2}$, with η the Sommerfeld parameter. For neutron channels, however, the ℓ -dependent centrifugal potential dominates at large distances, and at low energies one has

$$P_\ell / P_{\ell-1} = \rho^2 / (2\ell - 1)^2,$$

which tends to zero as E tends to zero.

For the conventional value of the channel radius, $a = 1.45(7^{1/3} + 1)$ fm = 4.22 fm, the ratio P_1/P_0 has the zero-energy values 0.034 and 0.048 for ${}^7\text{Li} + p$ and ${}^7\text{Be} + p$, respectively, and increases with energy. For ${}^7\text{Li} + n$, P_1/P_0 has the values 1.2×10^{-6} and 4.7×10^{-4} at $E_n = 1.78$ eV and 721 eV, the average energies for the extreme bins in the measurement of Blackmon *et al.* [1].

Chasteler *et al.* [2] obtained fits to their ${}^7\text{Li}(p, \gamma_0){}^8\text{Be}$ data with (18–95)% p -wave strength and found that these could reduce the zero-energy S factor by (7–38)%. By attributing the p -wave strength to the tails of known 1^+ levels of ${}^8\text{Be}$, Barker [6] found acceptable fits to the data of Chasteler *et al.* with $\lesssim 9\%$ p -wave strength. Suppose that there were appreciable p -wave strength in ${}^7\text{Li}(p, \gamma_0){}^8\text{Be}$, say 10%. If all relevant quantities except penetration factors are assumed the same for ${}^7\text{Li}(n, \gamma_0){}^8\text{Li}$ as for ${}^7\text{Li}(p, \gamma_0){}^8\text{Be}$, then the different values of the ratio P_1/P_0 would imply p -wave strengths in ${}^7\text{Li}(n, \gamma_0){}^8\text{Li}$ of $3.8 \times 10^{-4}\%$ and 0.15% at 1.78 and 721 eV, respectively. The uncertainties in the measurement of Blackmon *et al.* [1] are, however, only sufficient to limit the p -wave strength to less than about 30%. Similarly, this measurement would not impose any significant restriction on possible low-energy p -wave strength in the ${}^7\text{Be}(p, \gamma){}^8\text{B}$ reaction.

- [1] J.C. Blackmon, A.E. Champagne, J.K. Dickens, J.A. Harvey, M.A. Hofstee, S. Kopecky, D.C. Larson, D.C. Powell, S. Raman, and M.S. Smith, *Phys. Rev. C* **54**, 383 (1996).
 [2] R.M. Chasteler, H.R. Weller, D.R. Tilley, and R.M. Prior, *Phys. Rev. Lett.* **72**, 3949 (1994).
 [3] A.M. Lane and R.G. Thomas, *Rev. Mod. Phys.* **30**, 257 (1958).

- [4] F.L. Yost, J.A. Wheeler, and G. Breit, *Phys. Rev.* **49**, 174 (1936).
 [5] G.N. Watson, *A Treatise on the Theory of Bessel Functions* (Cambridge University Press, Cambridge, England, 1958), p. 78 and Table IV.
 [6] F.C. Barker, *Aust. J. Phys.* **48**, 813 (1995).