Energy dependence of cross sections for pion double charge exchange on ^{60,62,64}Ni

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Double charge exchange cross sections have been measured at $\theta_L = 5^\circ$ and $T_{\pi} = 230$, 180, and 140 MeV for targets of ^{60,62,64}Ni. Results are compared with previous data at 292 MeV and with a generalized seniority model. [S0556-2813(97)02005-0]

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Pion-induced double charge exchange (DCX) has been investigated at several incident energies for ⁵⁸Ni, for which the double isobaric analog state (DIAS) is the ground state (g.s.). References [1,2] present 5° cross sections at a single incident energy of 292 MeV, while Ref. [3] measured an excitation function from 120 to 292 MeV, again at 5°.

A separate experiment [4] presented 292-MeV results for 60,62,64 Ni. Combined with 58 Ni data, the DIAS and g.s. cross sections were compared to predictions of a generalized seniority model [5,6], which gives

$$\sigma_{\text{DIAS}} = \frac{n(n-1)}{2} \left| \alpha + \frac{\beta}{n-1} \right|^2,$$

$$\sigma_{\text{g.s.}} = \frac{n(n-2)(2\Omega + 2 - n)}{2(2\Omega + 1)(n-1)} |\beta|^2.$$

Apart from a factor $(A_0/A)^C$, the amplitudes α, β are independent of N, Z, A of the target, but they do depend on incident energy and scattering angle. The quantity Ω is six for Ni (2 Ω represents a full shell), and *n* is the neutron excess,

TABLE I. Cross sections (in nb/sr) for the DIAS as a function of energy at 5° (laboratory) for natural and isotopic Ni targets.

T_{π} (MeV)	⁶⁰ Ni	⁶² Ni	⁶⁴ Ni	^{nat} Ni
292 ^a	295±55	471±72	974±172	198±26
260				196 ± 46
230	172 ± 29	360 ± 71	584 ± 114	168 ± 37
180	89 ± 38	325 ± 89	463 ± 142	
164				100 ± 33
140	52±45	93±82	212 ± 159	136±117

^aRef. [4].

n=N-Z. In the amplitude, the exponent *C* is about 1.5. As we are dealing with a very narrow range of *A*, we ignore this factor. At 292 MeV, the results were consistent with the generalized seniority model. The present experiment deals with results for ^{60,62,64}Ni at other energies.

The experimental setup was as in Ref. [4], except that the strip target contained two pieces of 60 Ni, and no nat Ni. Rather, nat Ni data points were obtained with a separate full target. Peak areas were extracted in the usual manner. Normalization runs using a CH₂ target were performed at each bombarding energy. Resulting DIAS cross sections are listed in Table I and plotted in Fig. 1 for each isotope.

Because DIAS cross sections and pion beam fluxes are larger at 292 MeV, statistics are limited at these lower energies. Nevertheless, we can compare results with senioritymodel predictions. Within the limited statistical accuracy, the ^{60,62,64}Ni DIAS data allow extraction of α,β , and their relative phase φ . The quantities of α,β are related to the earlier [4,5] A,B by $A = \alpha + \beta/\Omega$, $B = [(\Omega - 1)/\Omega]\beta$. At 292 MeV, fitting only DIAS cross sections resulted in $\alpha \approx \beta$

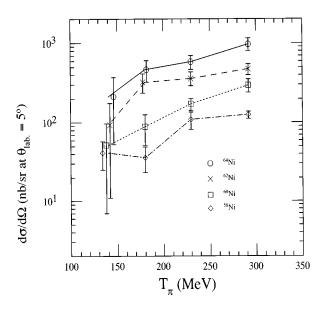


FIG. 1. Cross sections at 5° vs T_{π} for the DIAS in ^{58,60,62,64}Ni(π^+, π^-). Data for ⁵⁸Ni are from Ref. [3], and other 292-MeV data are from Ref. [4].

2468

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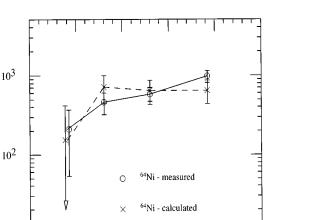
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10³

 10^{1}

100

 $d\sigma/d\Omega$ (nb/sr at $\theta_{lab.} = 5^{\circ}$)



250

300

350

FIG. 2. Measured cross sections for ⁶⁴Ni (circles) and those calculated (crosses) from 58,60,62Ni using generalized seniority: $\sigma_{64} = (4\sigma_{58} - 18\sigma_{60} + 20\sigma_{62})/7.$

200

 T_{π} (MeV)

150

and $\varphi \sim 90^\circ$, whereas a simultaneous fit to g.s. and DIAS gave $\varphi \sim 0^\circ$ and a smaller β . Or we can combine with ⁵⁸Ni (Table II) and arrive at a test of the model — which allows one of the four DIAS cross sections to be written in terms of the other three. Figure 2 plots the values of σ_{64} thus obtained, along with its measured values. The agreement is fair.

Seniority (or generalized-seniority) models provide a convenient prescription for seeing how T=1 nuclei differ from all others — it is only for T=1 that the β term contributes fully; for T > 1, it is attenuated by a factor 1/(n-1). Seldom are statistics good enough to separately distinguish the β

TABLE II. Measured 5° cross sections (nb/sr) for ^{58,64}Ni and calculated ⁶⁴Ni cross section.

		64	⁴ Ni	
T_{π} (MeV)	⁵⁸ Ni ^a	measured	calculated ^b	
292	125 ± 13	974±172	648±211	
230	109 ± 27	584 ± 114	649 ± 217	
180	36±13	463 ± 142	720 ± 272	
140	41 ± 16	$212{\pm}159$	155 ± 261	

^aRef. [3].

^bCalculated from other three isotopes using generalized seniority (Refs. [4,5]).

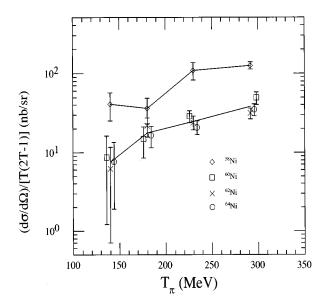


FIG. 3. Plot of $\sigma/[T(2T-1)]$ for ^{58,60,62,64}Ni.

coefficient in these higher T nuclei, but it is frequently the case that T=1 differs appreciably from the others. In the present example, we can note that difference by plotting $\sigma/[T(2T-1)]$ as is done in Fig. 3. There we see that the reduced cross section for T=1 is significantly greater than for the other three nuclei - all of which are essentially identical.

One motivation of the present study was to try to determine if the energy dependence of the DIAS cross sections were different for different isospin. A hint of a difference is apparent in Fig. 1 where the gap between data for ^{58,60}Ni and ^{62,64}Ni appears larger at 180 MeV than at other energies. This difference is scarcely statistically significant, but we can average cross sections at the two lowest and two highest energies separately (Table III) to decrease the uncertainties. When we do that, we find

$$\frac{\langle \sigma(140,180 \text{ MeV}) \rangle}{\langle \sigma(230,292 \text{ MeV}) \rangle}$$

for 58,60 Ni is 0.33 ± 0.08 , while for 62,64 Ni it is 0.50 ± 0.11 . The ratio of the two ratios is 0.65 ± 0.21 —not quite twostandard deviations from unity. Thus, despite the appearance in Fig. 1, we observe no statistically significant differences in the excitation functions of the four nuclei.

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TABLE III. Cross sections averaged for two low and two high energies.

$T_{\pi}(\text{MeV})$	⁵⁸ Ni	⁶⁰ Ni	⁶² Ni	⁶⁴ Ni	^{nat} Ni
230,292	122±12	199±26	415±51	703±95	188±21
140,180	38 ± 10	74 ± 29	209 ± 60	352 ± 106	103 ± 32
Ratio	0.31 ± 0.09	0.37 ± 0.15	0.50 ± 0.16	0.50 ± 0.16	0.55 ± 0.18
Ave. all T_{π}	80 ± 8	137 ± 19	312 ± 39	528 ± 71	162 ± 18
$\sigma_{\rm ave}/[T(2T-1)]$	80 ± 8	22.8 ± 3.2	20.8 ± 2.6	18.9 ± 2.5	

- [1] J.D. Zumbro et al., Phys. Rev. C 36, 1479 (1987).
- [2] K.K. Seth, S. Iversen, M. Kaletka, D. Barlow, A. Saha, and R. Soundranayagam, Phys. Lett. B 173, 397 (1986).
- [3] H.T. Fortune, G.B. Liu, J.M. O'Donnell, J.D. Silk, and S. Mordechai, Phys. Rev. C 50, 306 (1994).
- [4] D.R. Benton, H.T. Fortune, J.M. O'Donnell, R. Crittenden, M.

McKinzie, E. Insko, R. Ivie, D.A. Smith, and J.D. Silk, Phys. Rev. C 47, 140 (1993).

- [5] N. Auerbach, W.R. Gibbs, J.N. Ginocchio, and W.B. Kaufmann, Phys. Rev. C 36, 1277 (1988).
- [6] J.N. Ginocchio, Phys. Rev. C 48, 1460 (1993).