

### Coulomb excitation studies in antimony isotopes

K. C. Jain,\* K. P. Singh,\* G. Singh, S. S. Datta, and I. M. Govil

Department of Physics, Panjab University, Chandigarh 160 014, India

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Various low-lying energy levels up to 1.145 MeV in  $^{121,123}\text{Sb}$  have been Coulomb excited using 3.0–4.0 MeV proton-beam energies. The reduced transition probabilities of the excited levels have been newly deduced from the measured yields of the deexcitation gamma rays. The values of mixing ratios for a few transitions and the removal of ambiguities in the spin values of some levels have been presented from the angular distribution analysis.

The odd- $A$  transitional nuclei  $^{121,123}\text{Sb}$  have been subjected to several theoretical<sup>1,2</sup> and experimental<sup>3,4</sup> investigations. The properties of the levels up to 600-keV excitation energies are well established. However, the information about many levels beyond 600 keV are incomplete. In several cases the experimental  $B(E2)$  values reported by various groups<sup>1,3,4</sup> are quite discrepant among themselves and also differ from the theoretical predictions.<sup>1</sup> The assigned spin values for many of the levels are quite ambiguous.<sup>3,4</sup>

Most of the reported Coulomb excitation studies<sup>3–8</sup> on both the antimony isotopes were carried out with heavy ions and alpha particles using relatively coarse resolution and poor efficiency of the detectors. In the present investigations the gamma-ray yields have been measured with high resolution ( $\sim 1.9$  keV at 1.332 MeV) and better efficiency HPGe detector ( $57\text{ cm}^3$ ). The reduced quadrupole transition probabilities of six levels in  $^{121}\text{Sb}$  and four levels in  $^{123}\text{Sb}$  up to 1.145 MeV have been extracted from the gamma-ray yields observed at  $55^\circ$  with respect to

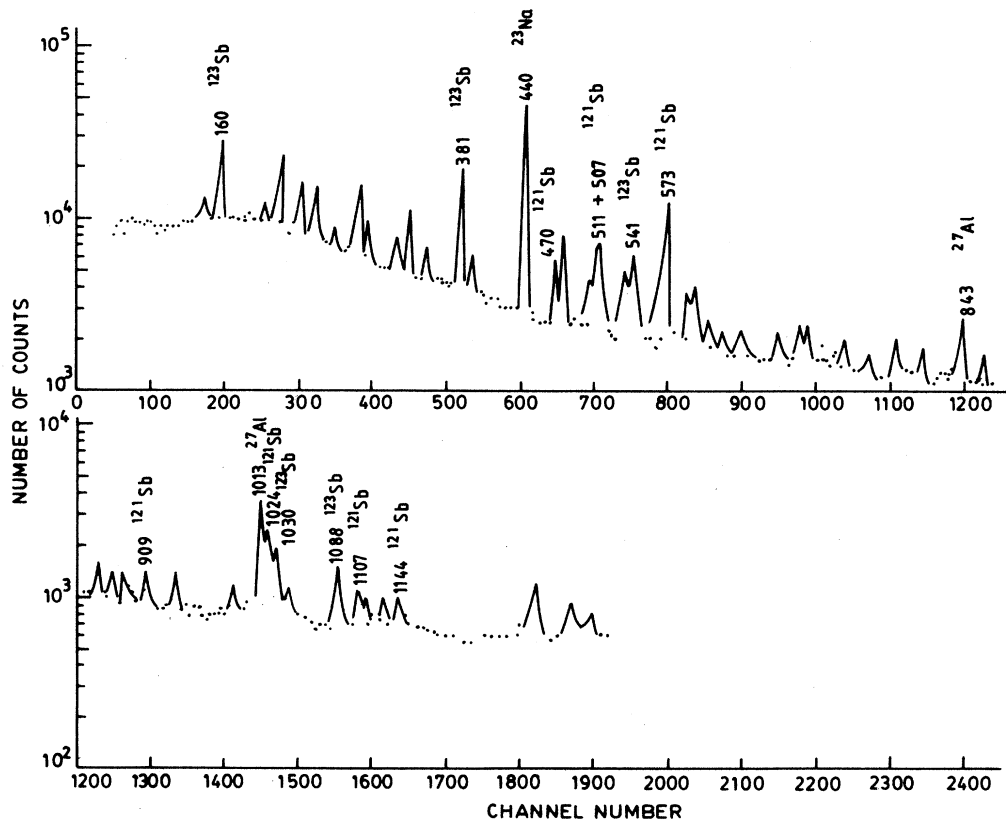


FIG. 1. Gamma-ray spectrum with 4.0-MeV protons incident on a thick natural antimony target. The peaks left unmarked are from  $(p, n\gamma)$  reaction and background.

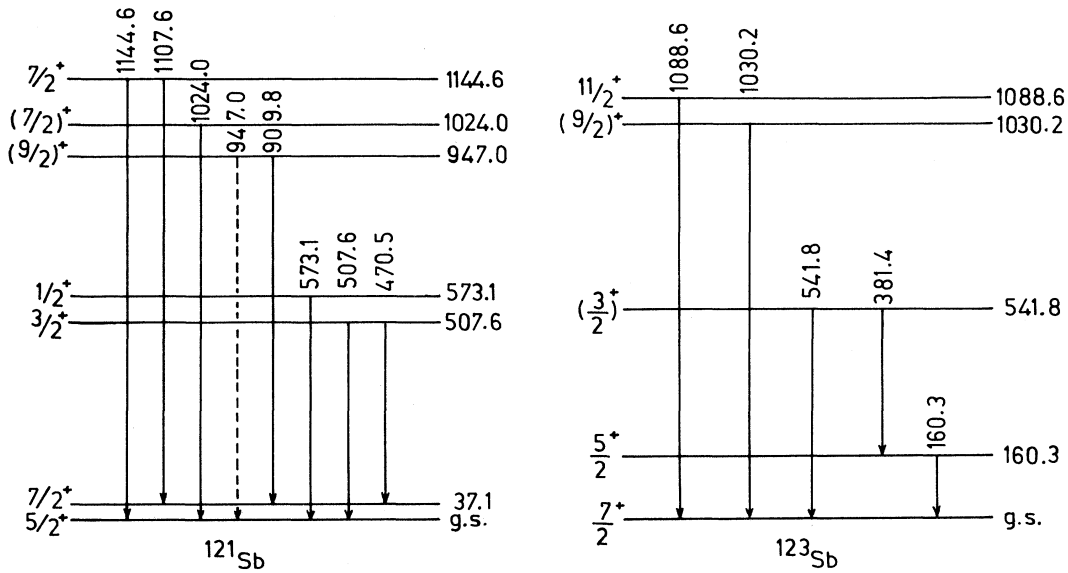


FIG. 2. Level schemes proposed for the Coulomb excited levels in  $^{121}\text{Sb}$  and  $^{123}\text{Sb}$ .

beam direction. The results have been discussed in the light of earlier experimental data<sup>5-8</sup> and the predictions of the intermediate coupling model.<sup>2</sup>

The low-lying levels in  $^{121,123}\text{Sb}$  were excited with 3.0–4.0-MeV proton beams using a thick target of spectroscopically pure antimony. The details of the experiments are described elsewhere.<sup>9</sup> A typical spectrum at  $90^\circ$  with a 4.0-MeV proton beam is shown in Fig. 1. The angular distributions were measured at  $0^\circ$ ,  $30^\circ$ ,  $45^\circ$ ,  $75^\circ$ , and  $90^\circ$  for spin assignments. The photopeaks corresponding to the gamma ray from energy levels at 507.6,

573.1, 947.0, 1024.0, and 1144.6 keV in  $^{121}\text{Sb}$  and from energy levels at 150.3, 541.8, 1030.2, and 1088.6 keV in  $^{123}\text{Sb}$  are clearly resolved in the spectrum. The unmarked peaks in the spectrum belong to  $(p,n\gamma)$  reactions from these two nuclei. The compound contribution to the total  $(p,p'\gamma)$  cross sections at this proton energy were calculated with the computer code CINDY (Ref. 10) and found to be less than 5% of the Coulomb contribution. The proposed level schemes for the Coulomb excited levels in  $^{121}\text{Sb}$  and  $^{123}\text{Sb}$  are shown in Fig. 2.

The present experimental  $B(E2)$  values and the

TABLE I. Present branching ratios and  $B(E2)$  values along with the earlier data on Coulomb excited levels in  $^{121,123}\text{Sb}$ .  $B(E2)$  ( $e^2\text{cm}^4 \times 10^{-50}$ ).

Energy level (keV)	$E_\gamma$ (keV)	Present branching ratio (%)	Present	Barnes <i>et al.</i> (Ref. 5)	Galperin <i>et al.</i> (Ref. 6)	Andreev <i>et al.</i> (Ref. 7)	Kulkarni and Patrawale (Ref. 8)	Hooper <i>et al.</i> <sup>a</sup> (Ref. 2)
$^{121}\text{Sb}$								
507.59	507.59	80						
	470.47	20	$1.32 \pm 0.41$	$0.7 \pm 0.2$	$1.3 \pm 0.4$	$1.1 \pm 0.2$	b	0.61
573.14	573.14	100	$2.70 \pm 0.40$	$2.7 \pm 0.3$	$2.0 \pm 0.4$	$2.8 \pm 0.2$		2.30
946.98	909.84	100	$0.06 \pm 0.02$			$0.07 \pm 0.02$		0.01
1024.00	1024.00	100	$1.80 \pm 0.35$	$10.0 \pm 1.6$		$7.0 \pm 0.5$		2.20
1144.65	1144.65	74	$2.24 \pm 0.33$		$3.0 \pm 0.1$	$8.1 \pm 0.5$		6.40
	1107.60	26						
$^{123}\text{Sb}$								
160.3	160.33	100	$0.23 \pm 0.04$	$0.43 \pm 0.07$	$0.23 \pm 0.08$		$0.46 \pm 0.04$	0.40
541.8	541.8	26						
	381.4	74	$3.00 \pm 0.30$	$2.8 \pm 0.6$	$2.8 \pm 0.6$	$4.0 \pm 0.3$		3.10
1030.23	1030.23	100	$1.96 \pm 0.30$	$8.0 \pm 1.0$	$9.0 \pm 2.0$	$9.0 \pm 2.0$	$7.0 \pm 4.5$	4.40
1088.64	1088.64	100	$2.06 \pm 0.3$	$7.0 \pm 1.0$	$4.2 \pm 0.9$	$7.6 \pm 0.8$	$3.0 \pm 0.28$	$6.5, 0.2^c$

<sup>a</sup>Results from intermediate coupling model.

<sup>b</sup> $B(E2)$  values by Kulkarni *et al.* (Ref. 8) for  $^{121}\text{Sb}$  contain large systematic errors.

<sup>c</sup>The 6.5 and 0.2  $e^2\text{cm}^4 \times 10^{-50}$  values have been calculated for  $J^\pi = \frac{11}{2}^+$  and  $\frac{9}{2}^+$ , respectively.

TABLE II. Present  $B(E2)$  values compared with the single-particle estimates.

Isotope	Energy level	Spin $J^{\pi a}$	$B(E2)\downarrow$ ( $e^2 \text{ cm}^4 \times 10^{-50}$ )	$B(E2)\downarrow / B(E2)_{s.p.}^b$
$^{121}\text{Sb}$	507.6	$\frac{3}{2}^+$	$1.98 \pm 0.62$	$4.73 \pm 1.5$
	573.1	$\frac{1}{2}^+$	$8.10 \pm 1.20$	$19.4 \pm 2.7$
	947.0	$(\frac{9}{2})^+$	$0.036 \pm 0.012$	$0.09 \pm 0.03$
	1024.0	$(\frac{7}{2})^+$	$1.35 \pm 0.26$	$3.2 \pm 0.6$
	1144.6	$\frac{7}{2}^+ a$	$1.68 \pm 0.25$	$4.0 \pm 0.6$
$^{123}\text{Sb}$	160.3	$\frac{5}{2}^+$	$0.31 \pm 0.05$	$0.72 \pm 0.12$
	541.8	$(\frac{3}{2})^+$	$6.0 \pm 0.6$	$14.0 \pm 1.4$
	1030.2	$(\frac{9}{2})^+$	$1.57 \pm 0.24$	$3.7 \pm 0.6$
	1088.6	$\frac{11}{2}^+ a$	$1.37 \pm 0.20$	$3.2 \pm 0.5$

<sup>a</sup> $J^{\pi}$  values for 1144.6 and 1088.6 keV were assigned in the present work, other values were taken from literature. Those enclosed in the parentheses are tentative assignments.

<sup>b</sup>The single-particle estimates  $B(E2)_{s.p.}$  have been obtained from  $e^2/4\pi | \frac{3}{5} R_0^2 |^2$ , with  $R_0 = 1.25 \times 10^{-13} A^{1/3}$  cm.

branching ratios along with the  $B(E2)$  values from the earlier measurements<sup>5-8</sup> are summarized in Table I. Our values in  $^{121}\text{Sb}$  for the 507.6-, 573.1-, and 947.0-keV levels agree with the earlier data.<sup>3,4</sup> However, for the levels at 1024 keV, our results are lower by a factor of  $\sim 4$  compared to others. This may be partly due to the fact that the earlier authors, due to poor resolution, could not resolve this state from the adjacent peaks due to  $^{27}\text{Al}$  and  $^{123}\text{Sb}$ . In the case of  $^{123}\text{Sb}$ , our results for 160.3 and 541.8 keV agree with earlier data, except at 1030.2-keV levels for which our  $B(E2)$  value is  $\sim 4$  times smaller than others. This is also due to the same reason of poor resolution and interference of adjacent peaks due to  $^{27}\text{Al}$  and  $^{121}\text{Sb}$  in the earlier data.

Table II compares the experimental  $B(E2)$  values with single-particle estimates. Our data confirm that the 947-

keV level in  $^{121}\text{Sb}$  and 160.3-keV level in  $^{123}\text{Sb}$  are essentially of single-particle character, while the 573.1- and 541.9-keV levels have collective structures. The other levels at 507.6, 1024.0, and 1144.6 keV in  $^{121}\text{Sb}$  and 1030.2 and 1088.6 keV in  $^{123}\text{Sb}$  seem to have mixed nature intermediate between these two extremes. These observations are consistent with the predictions of Conjeaud *et al.*<sup>11</sup> based on  $^{120}\text{Sn}(^3\text{He},d)$  and  $^{122}\text{Te}(t,\alpha)$  reactions.

The angular distribution for the 1088.6-keV level in  $^{123}\text{Sb}$  predicts the  $A_2$  value equal to  $0.175 \pm 0.021$ , which suggests the spin of this level as  $\frac{11}{2}$  rather than  $\frac{9}{2}$ . Similarly, for the 1144.6-keV level in  $^{121}\text{Sb}$  our experimental value of  $A_2 = -0.048 \pm 0.007$  suggests the spin of this level as  $\frac{7}{2}$  and mixing ratio ( $\delta$ ) as  $-0.162$  for the 1144.6-keV transition from this state.

\*Present address: D.A.V. College, Chandigarh, India.

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