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Quadrupole moment of ${}^{107}Ag^m$ determined by level mixing resonance on oriented nuclei

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The quadrupole moment ratio of ${}^{107}\text{Ag}^m$ to ${}^{109}\text{Ag}^m$ has been determined as $Q({}^{107}\text{Ag}^m)/Q({}^{109}\text{Ag}^m) = 0.960(10)$ with the level mixing resonance on oriented nuclei technique. Using the previously established electric field gradient of silver in zinc, $Q({}^{107}\text{Ag}^m) = 0.98(11)$ b has been deduced. The results are in agreement with the theoretical predictions.

The precision of the electric quadrupole moments, determined through hyperfine interactions in solid matrices, is limited on account of the experimental error on the electric field gradients. Nevertheless, the recently developed technique of level mixing resonances on oriented nuclei (LMR/ON) allows the determination of the quadrupole resonance frequency to a high precision,¹⁻³ and thus to establish an accurate ratio of quadrupole moments. This feature is particularly interesting, when the nuclear quadrupole moment varies smoothly from one isotope to another, as in the case of silver isotopes for which a smooth increase of the prolate deformation versus mass number is predicted by a macroscopic-microscopic calculation of Möller and Nix.⁴ The predicted quadrupole moments of some silver isotopes have been confirmed in a previous paper,⁵ but the experi-mental errors are rather high: $Q(^{108}Ag^m)$ and $Q(^{110}Ag^m)$ are established by optical spectroscopy⁶ and corrected for the are established by optical spectroscopy and corrected for the Sternheimer effect, $Q(^{106}Ag^m)$ was obtained from low temperature nuclear orientation of ^{106}Cd in zinc, and the isotope $^{109}Ag^m$ was measured by LMR/ON.⁵ For the evaluation of this latter measurement we used the electric field gradient determined from the low temperature nuclear orientation of ¹¹⁰Ag^m in zinc.⁵ The convolution of the errors of each measurement leads thus to a limited absolute precision on Q.

As ${}^{107}\text{Ag}^m$ is also open to LMR/ON, we measure this resonance in a single crystal sample ${}^{107}\text{Cd}Zn$.

Natural silver foils were irradiated with the 28 MeV deuteron beam of the synchrocylotron of the University of Lyon. The $^{107}Cd(6.5 h)$ activity was separated chemically from the silver foil, electrodeposited onto a Zn single crystal cut perpendicularly to the *c* axis. The quality of each crystal has been checked previously by Laue backscattering. The carrier-free ^{107}Cd activity was diffused into the crystal at 300 °C during 3 h under H₂ atmosphere, then the remaining activity was cleaned from the surface.

In order to assure that 109 Ag^m and 107 Ag^m experience the same electric field gradient, we tried to introduce 109 Cd and 107 Cd into the same crystal. On account of the high inactive carrier concentration of 109 Cd, it was not possible to diffuse 109 Cd and 107 Cd together, during only a few hours. (In our previous investigation 2 109 Cd has been diffused during 24 h.) The evaporation and the diffusion of 109 Cd, followed by that of the 107 Cd activity, introduced 107 Cd and 109 Cd to very different depths, and LMR/ON showed an appreciable line broadening, which can be due to imperfections in the crystal sites. So, it was preferable to introduce 107 Cd alone into the Zn single crystal.

The $\Delta m = 2$ and 3 resonances on ${}^{107}\text{Ag}^m$ have been ob-



FIG. 1. The $\Delta m = 3$ LMR/ON resonance on 107 Ag^mZn observed at a temperature of 6.7 mK. The tilting angle between the magnetic field and the electric field gradient was 7°.

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served in a dilution refrigerator on several samples (Fig. 1). The positions of the resonances between various samples showed scattering up to about $\pm 1\%$. This scattering is higher than that observed in our previous investigation on ¹⁰⁹Ag^m, where the long lifetime of ¹⁰⁹Cd allowed for a long diffusion and more careful annealing of the samples. We thus find for ¹⁰⁷Ag^m in zinc

$$eQV_{zz}/\mu = 2.96(3) \text{ T}$$

and with $\mu = 4.398(5)$ from Ref. 7

$$v_O(^{107} \text{Ag}^m Zn) = eQV_{zz}/h = 99.2(10) \text{ MHz}$$
.

Reevaluating with the new magnetic moment value of Ref. 7

$$v_O(^{109}\text{Ag}Zn) = 103.3(7) \text{ MHz}$$
.

Using the electric field gradient of silver in zinc of Ref. 5 we find

$$Q(^{107}\text{Ag}, \frac{7}{2}) = 0.98(11) \text{ b}, \quad Q(^{109}\text{Ag}, \frac{7}{2}) = 1.02(12) \text{ b}.$$

- ¹R. Coussement, P. Put, L. Hermans, M. Rots, I. Berkes, R. Brenier, and G. Marest, Phys. Lett. **97A**, 301 (1983).
- ²P. Put, R. Coussement, G. Scheveneels, F. Hardeman, I. Berkes, B. Hlimi, G. Marest, J. Sau, and E. H. Sayouty, Phys. Lett. 103A, 151 (1984).
- ³P. Put, R. Coussement, G. Scheveneels, F. Hardeman, I. Berkes, B. Hlimi, G. Marest, and J. Sau, Hyperfine Interact. 22, 131 (1985).

The errors on Q are mainly due to the error on the electric field gradient.

The ratio of the electric quadrupole moments, exempt of this error

$$Q(^{107}\text{Ag},\frac{7}{2}^+)/Q(^{109}\text{Ag},\frac{7}{2}^+) = 0.960(10)$$
,

is almost one order of magnitude more precise than the numerical value of Q, and it is in excellent agreement with the theoretical estimation of the ratio of the intrinsic quadrupole moments 0.95 of Möller and Nix.⁴ The absolute value of the intrinsic quadrupole moment $Q_0(^{107}\text{Agm}) = 2.1(2)$ b also agrees fairly well with the theoretical value of 1.8 b.

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- ⁴P. Möller and J. R. Nix, At. Data Nucl. Data Tables **26**, 165 (1981).
- ⁵I. Berkes, B. Hlimi, G. Marest, E. H. Sayouty, R. Coussement, F. Hardeman, P. Put, and G. Scheveneels, Phys. Rev. C **30**, 2026 (1984).
- ⁶W. Fischer, H. Hühneman, and Th. Meier, Z. Phys. A 274, 79 (1975).
- ⁷R. Eder, H. Hagn, and E. Zech, Phys. Rev. C 31, 190 (1985).