

Brief Reports

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Diffractive angular distribution for $^{16}\text{O}(\pi^+, \pi^-)^{16}\text{Ne}(\text{g.s.})$

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The 164-MeV angular distribution for the non-analog double charge exchange reaction $^{16}\text{O}(\pi^+, \pi^-)^{16}\text{Ne}(\text{g.s.})$ has been measured and is found to be basically diffractive. This is contrasted with the previously measured analog distribution for $^{18}\text{O}(\pi^+, \pi^-)^{18}\text{Ne}$ (double isobaric analog) which is not diffractive.

NUCLEAR REACTIONS $^{16}\text{O}(\pi^+, \pi^-)^{16}\text{Ne}(\text{g.s.})$; $E_\pi = 164$ MeV, $\theta = 0^\circ - 50^\circ$, measured $\sigma(\theta)$, non-analog transition, eikonal fit to angular distribution.

An intriguing observation in pion double charge exchange (DCX) is the existence of large non-analog cross sections for (π^+, π^-) reactions on $T=0$ nuclei.^{1,2} As seen in the oxygen isotopes at 164 MeV, cross sections for ^{16}O ($T=0$) are nearly as large as those for ^{18}O ($T=1$). This result is in conflict with expectations that transitions to double-isobaric-analog states (DIAS) would dominate because of the high degree of overlap between initial- and final-state wave functions.

With rare exception,³ all previous treatments of DCX completely ignore excitation of non-analog final states. The non-analog process has important consequences for the interpretation of double-analog DCX on $T=1$ nuclei, particularly when $T=1$ targets can be viewed as $T=0$ cores plus two neutrons.

In addition to other interesting features,² the angular distribution for DCX on ^{18}O has a peculiar shape at 164 MeV, but not at 292 MeV. The anomalous shape is observed at the bombarding energy (164 MeV) at which the non-analog process from ^{16}O has

greatest strength. It thus became of interest to measure the angular distribution for DCX on ^{16}O at the same energy.

The measurements were performed at the Clinton P. Anderson Meson Physics Facility (LAMPF) using the EPICS pion spectrometer facility. The small-angular DCX apparatus² was installed and ice targets of 2-g/cm² areal density were used to maximize count rates consistent with acceptable resolution. Absolute normalizations were determined by comparing π^+p elastic scattering, under the same conditions, with phase-shift calculations. Target-thickness effects were checked by comparing yields with a 0.25-g/cm² ice target. The measured absolute cross section scale is believed accurate to $\pm 15\%$.

The measured angular distribution is shown in Fig. 1. There is a sharp falloff from forward angles with a distinct minimum near 32° , then a short rise to a second maximum near $45-50^\circ$. This shape is in striking contrast with that for ^{18}O , which is shown for comparison. The ^{16}O angular distribution is well

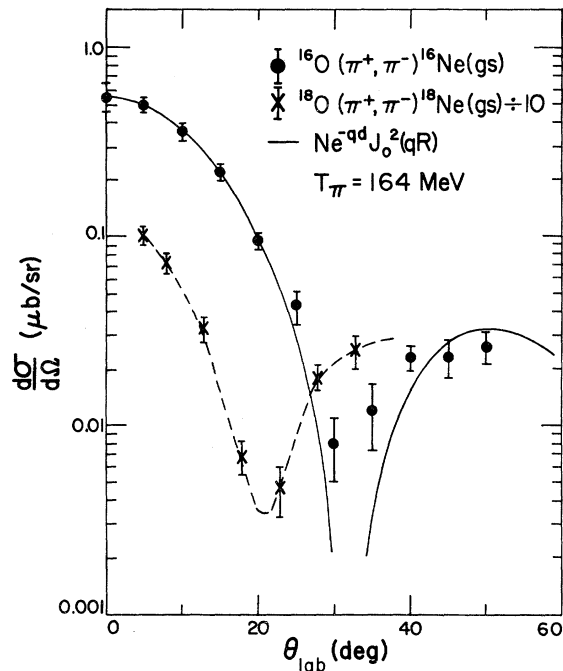


FIG. 1. Comparison of 164-MeV angular distributions for the DCX reactions $^{16,18}\text{O}(\pi^+, \pi^-)^{16,18}\text{Ne}(\text{g.s.})$. The solid curve is an eikonal calculation explained in the text. The dashed line merely outlines the ^{18}O data.

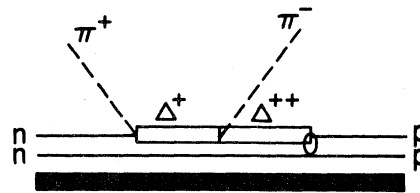


FIG. 2. A proposed (Ref. 5) single-step process for non-analog DCX.

described by an eikonal form,⁴ $\sigma(\theta) = Ne^{-qd}J_0^2(qR)$, shown by the solid line in Fig. 1. For this curve the parameters are $d = 1.2$ fm and $R = 3.3$ fm, which is a reasonable value for the strong absorption radius of ^{16}O .

The diffractive nature of this angular distribution suggests that DCX on ^{16}O is a single-step process involving only one basic amplitude. In Ref. 5, this amplitude is suggested to result from the process outlined in Fig. 2, viz., DCX on a neutron, turning it into a Δ^{++} in the final state. It is clear that any single-step DCX process must involve a Δ . Any theoretical treatment of DCX on $T = 1$ nuclei must be able to correctly account for non-analog DCX on the core ($T = 0$ nuclei) before it has any hope of success.

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