Rapid Communications

The Rapid Communications section is intended for the accelerated publication of important new results. Manuscripts submitted to this section are given priority in handling in the editorial office and in production. A Rapid Communication may be no longer than 3½ printed pages and must be accompanied by an abstract and a keyword abstract. Page proofs are sent to authors, but, because of the rapid publication schedule, publication is not delayed for receipt of corrections unless requested by the author.

Width of the 6⁻, T = 1, $E_x = 14.36$ MeV state in ²⁸Si and its relationship to intermediate energy inelastic scattering

Dean Halderson, K. W. Kemper, and J. D. Fox Department of Physics, Florida State University, Tallahassee, Florida 32306

R. O. Nelson, E. G. Bilpuch, and C. R. Westerfeldt Duke University, Durham, North Carolina 27706 and Triangle Universities Nuclear Laboratory, Durham, North Carolina 27706

G. E. Mitchell

North Carolina State University, Raleigh, North Carolina 27650 and Triangle Universities Nuclear Laboratory, Durham, North Carolina 27706

(Received 6 May 1981)

The width of the 6⁻, T = 1, $E_x = 14.36$ MeV state in ²⁸Si has been measured via the ²⁷Al(p,p)²⁷Al reaction. The value of $\Gamma = 3.7 \pm 0.4$ keV yields a single particle fraction $\beta_1^2 = 0.68 \pm 0.07$ which is shown to be consistent with the combined results for the inelastic excitation of this state by π^{\pm} , p, and e^- and the ²⁷Al(³He, d)²⁸Si (g.s.) reaction.

NUCLEAR REACTIONS ²⁷Al($p, p_{0,1,2}$)²⁷Al; $E_p = 2.70 - 3.06$ MeV; determined width and deduced single particle fraction for ²⁸Si state at $E_x = 14.36$ MeV.

The last three years have provided nuclear physics with a variety of new intermediate energy inelastic scattering data.¹ Proton inelastic scattering in the 100-140 MeV range, pion scattering near the (3,3) resonance, and high resolution electron scattering appear to be powerful tools for investigating nuclear structure.² However, conclusions concerning nuclear structure will depend upon (1) understanding the reaction mechanism and (2) selecting reasonable models for comparison with data. It is necessary to perform numerous tests of these two points before the intermediate energy results can add to the knowledge of nuclear structure in a manner consistent with low energy measurements. The present paper provides one such test.

There are many nuclear levels for which excitation by p, π^{\pm} , and e^{-} yield consistent results under the assumption that the levels and corresponding ground states possess a certain structure. Perhaps the most useful of these levels have been the stretched state configurations in which a particle is promoted from the active hole orbit with largest j_h to the active particle orbit with largest j_p to form a state with maximum total J. Because of their high spins, these states are preferentially excited in high momentum transfer reactions and also remain relatively pure due to the low density of states with large J. Owing to their stretched character, the direct inelastic excitation of these states is completely described by a single spin transition density which allows for a direct comparison of (e,e'), (p,p'), and (π, π') scattering.²⁻⁵ Obtaining consistent results for the different probes for these stretched states provides evidence that the reaction mechanisms are understood and that a reasonable approximation for their structure has been assumed.

It is now desirable to select from those observed states which have yielded consistent inelastic results that subgroup of states which can be compared with low energy results. One such state is the 6⁻, T = 1, $E_x = 14.36$ MeV level in ²⁸Si. The purpose of this paper is to present a measurement of the width of this state via the resonance reaction, ²⁷A1(*p*,*p*), and to demonstrate that the measured value of 3.7 ± 0.4 keV

<u>24</u>

786

corroborates the structure assumed in the inelastic scattering results. In addition this measurement provides information complementary to inelastic scattering and should aid in understanding future excitations of this level.

The experiment was performed with the high resolution system on the 3 MV Van de Graaff accelerator of Triangle Universities Nuclear Laboratory. The ${}^{27}\text{Al}(p,p_0)$, (p,p_1) , and (p,p_2) reactions were measured from $E_p = 2.70$ to 3.06 MeV. The targets were $1-2 \ \mu\text{g/cm}^2$ of aluminum evaporated onto thin carbon backings; the overall resolution was 350 eV full width at half maximum (FWHM).

The elastic scattering cross section for a portion of these data is shown in Fig. 1. The solid line in Fig. 1 is a preliminary fit to these data utilizing a multilevel, multichannel *R*-matrix code.⁶ The elastic cross section is dominated by a series of broad l=0 resonances and one very strong resonance near $E_p = 2.87$ MeV. The strong resonance at $E_p = 2.873$ MeV definitely has l=3 with a strongly preferred J=6 assignment, and is thus the T=1 state observed in the inelastic scattering reactions on ²⁸Si.

Twenty-two resonances were observed in the preliminary analysis; some of these resonances have appreciable strength in our inelastic channels. The 6⁻ state has no observed strength in either of our inelastic channels, although inelastic decay from low spin states makes a precise determination of the lower limit difficult. α channels are open, but decay to the 0⁺ ground state is not allowed, and decay to the 2⁺ first excited state of ²⁴Mg was not observed by Neal and Chagnon.⁷ The capture reaction was studied by Miehe *et al.*⁸ and by Neal and Lam.⁹ We obtain an excellent fit to the 6⁻ state assuming $\Gamma_p = \Gamma_{\text{total}}$. Our value for the width is $\Gamma_p = 3.7 \pm 0.4$ keV.

If one proceeds as in Refs. 2-5, the ²⁸Si ground



FIG. 1. The ²⁷Al(p,p) differential cross section at $\theta_L = 160^\circ$ in the vacinity of the $J^{\pi} = 6^-$, T = 1 state at $E_p = 2.873$ MeV. The solid curve is the result of a multilevel, multichannel *R*-matrix calculation.

state, $|i\rangle$, is assumed to be a closed $0d_{5/2}$ shell plus a more complicated component,

$$|i\rangle = \alpha_1|0\rangle + \alpha_2|c_i\rangle$$
, $\alpha_1^2 + \alpha_2^2 = 1$. (1)

The 6^- , T = 1 state is similarly taken to be a pure particle hole excitation plus a more complicated component,

$$|f\rangle = \beta_1 [a_{f_{7/2}}^{\dagger} a_{d_{5/2}}^{\dagger}]^{6,1} |0\rangle + \beta_2 |c_f\rangle ,$$

$$\beta_1^2 + \beta_2^2 = 1 .$$
(2)

Since the direct inelastic excitation operator is a single particle operator, the transition, $|0\rangle \rightarrow |np-nh\rangle$, n > 1, will not proceed. Therefore it is not unreasonable to expect that the dominant contribution to the inelastic cross section comes from $|0\rangle \rightarrow [a^{\dagger}a]^{6,1}|0\rangle$. If the transition between pure states $|0\rangle \rightarrow [a^{\dagger}a]^{6,1}|0\rangle$ is declared a sum rule, then the fraction of the sum rule exhausted by the transition provides a measure of $\alpha_1^2\beta_1^2 \equiv S_{ph}^2$. For the 6⁻, T = 1 state, the (π, π') , $^{10}(p,p')$, 11 and $(e,e')^{12}$ data agree remarkably well with the average value of $S_{ph}^2 \sim 0.31.^{2-5}$ The results of these experiments are summarized in Ref. 4.

This consistency among experiments provides confidence in one's understanding of the reaction mechanisms. One may then proceed to make the conclusion that the ²⁸Si ground state and 6⁻, T = 1state deviate substantially from the single particle shell model, such that $S_{ph}^2 \sim 0.31$. A comparison with ²⁷Al(p,p)²⁷Al results can now be demonstrated consistent with this conclusion.

Experimental evidence is available for the structure of the ²⁸Si ground state. If one assumes that ²⁷A1 (ground state) = $d_{5/2}^{-1}(p)|0\rangle$ and that ²⁸Si (ground state) = $|0\rangle$, then (³He,d) data account for only 44% of calculated strength.¹³ This is a considerable departure from the single particle model and the 6⁻, T = 1 state would necessarily be a fairly pure single particle state if S_{ph}^{-2} does equal 0.31. Specifically, it would have to look like an $f_{7/2}$ particle coupled to the A = 27 core 0.31/0.44 × 100 = 70% of the time.

The present experiment provides this information on the 6⁻, T = 1 state. The measured width of 3.7 keV can be used to extract the single particle fraction as follows. The width of an $f_{7/2}$ proton resonance at $E_p^{c.m.} = 2.77$ MeV in a Woods-Saxon well with geometry, V = 58.62 MeV, $r_0 = r_{so} = r_c = 1.2$ fm, $a_0 = a_{so} = 0.569$ fm, and λ of the Thomas spin-orbit term =25, is 10.8 keV. For a pure T = 1 state, which spends one-half of its time as a proton coupled to 2^{7} Al, this gives $\Gamma_p = 5.4$ keV. The well geometry was determined by requiring that the resonance appear at 2.77 MeV and that the well best fitted in the important surface region that potential obtained by single folding the V_0 component of the g matrix of Bertsch *et al.*¹⁴ The single particle fraction is then Γ/Γ_p = $3.7/5.4 = 0.68 \pm 0.07$, which is in near agreement with the value of 0.70 from combining the inelastic scattering and (³He, d) results. The inelastic scattering conclusion concerning the single particle fraction of the $0^+ \rightarrow 6^-$, T = 1 transition is therefore consistent with resonance measurements.

This paper has described a measurement of the width of the 6⁻, T = 1, $E_x = 14.36$ MeV state in ²⁸Si. It was demonstrated that the measured value of $\Gamma = 3.7 \pm 0.4$ keV yields a single particle fraction of 0.68 ± 0.07 which is consistent with intermediate energy inelastic scattering values. This measurement has not only provided a test of the understanding of the intermediate energy inelastic scattering mechanism, but also can provide the experimentalist with valuable information on high energy resolution at

- ¹F. Petrovich and W. G. Love, in Proceedings of the International Conference on Nuclear Physics, Berkeley, California, August 24-30, 1980, edited by R. M. Diamond and J. O. Rasmussen (North-Holland, Amsterdam, 1981); also see various papers in Workshop on Nuclear Structure with Intermediate-Energy Probes, conference proceedings, LA-8303-C, Los Alamos, 1980 (unpublished).
- ²F. Petrovich, *The* (*p*,*n*) *Reaction and the Nucleon-Nucleon Force* (Plenum, New York, 1979), p. 115.
- ³R. A. Lindgren, W. J. Gerace, A. D. Bacher, W. G. Love, and F. Petrovich, Phys. Rev. Lett. <u>42</u>, 1524 (1979).
- ⁴F. Petrovich, W. A. Love, A. Picklesimer, G. E. Walker, and E. R. Siciliano, Phys. Lett. 95B, 166 (1980).
- ⁵F. Petrovich, D. Halderson, J. Carr, and H. McManus (unpublished).
- ⁶E. G. Bilpuch, A. M. Lane, G. E. Mitchell, and J. D. Moses, Phys. Rep. 28C, 145 (1976).
- ⁷G. F. Neal and P. R. Chagnon, Phys. Rev. <u>C11</u>, 1461 (1975).

this excitation energy. In addition, the present measurement shows that the 6^- , T = 1 state is a good single particle state and other reactions which proceed via a single particle mechanism should show large yields. Finally, this measurement represents only one demonstration of the consistency between low energy and intermediate energy results. Centainly an effort should be made to increase the amount of data where low and intermediate energy experiments overlap so that additional comparisons can be made.

The authors wish to acknowledge helpful discussions with R.A. Lindgren, F. Petrovich, and D. Robson. This work was supported in part by the National Science Foundation and the United States Department of Energy.

- ⁸C. Miehe, J. P. Gonidec, A. Huck, and G. Walter, Rev. Phys. Appl. 8, 22 (1973).
- ⁹G. F. Neal and S. T. Lam, Phys. Lett. 45B, 127 (1973).
- ¹⁰C. Olmer, B. Zeidman, D. J. Geesaman, T. -S. H. Lee, R. E. Segel, L. W. Swenson, R. L. Boudrie, G. S. Blanpied, H. A. Theissen, C. L. Morris, and R. E. Anderson, Phys. Rev. Lett. 43, 612 (1979).
- ¹¹G. S. Adams, A. D. Bacher, G. T. Emergy, W. P. Jones, R. T. Kouzes, D. W. Miller, A. Picklesimer, and G. E. Walker, Phys. Rev. Lett. 38, 1387 (1977).
- ¹²S. Yen, R. Sobie, H. Zarek, B. O. Pich, T. E. Drake, C. F. Williamson, S. Kowalski, and C. P. Sargent, Phys. Lett. 93B, 250 (1980).
- ¹³R. W. Barnard and G. D. Jones, Nucl. Phys. <u>A108</u>, 641 (1968).
- ¹⁴G. Bertsch, J. Borysowicz, H. McManus, and W. G. Love, Nucl. Phys. A284, 399 (1977).