Brief Reports

Brief Reports are short papers which report on completed research or are addenda to papers previously published in the Physical Review. A Brief Report may be no longer than 3½ printed pages and must be accompanied by an abstract and a keyword abstract.

Alpha-decay properties of ^{205,206,207,208}Fr: Identification of ²⁰⁶Fr^m

B. G. Ritchie

University of South Carolina, Columbia, South Carolina 20208

K. S. Toth

Oak Ridge National Laboratory, Oak Ridge, Tennessee 37830

H. K. Carter, R. L. Mlekodaj, and E. H. Spejewski UNISOR, Oak Ridge Associated Universities, Oak Ridge, Tennessee 37830 (Received 12 January 1981)

Alpha-particle and γ -ray spectral measurements were made for ²⁰⁵⁻²⁰⁸Fr. A new α emitter ($T_{1/2} = 0.7 \pm 0.1$ sec and $E_{-} = 6.930 \pm 0.005$ MeV) was observed and identified with the decay of a previously unknown isomer in ²⁰⁶Fr. From the α particle and γ ray intensities, α decay branching ratios were deduced for ²⁰⁵⁻²⁰⁸Fr utilizing available information concerning the nuclides' (electron capture + positron) decay properties. Reduced widths were calculated and compared with those of neighboring nuclei.

RADIOACTIVITY 205,206,207,208 Fr, measured E_{γ} , I_{γ} , E_{α} , I_{α} ; deduced α -branching
ratios; identified 206 Fr^m. Mass separation.

The α decays of francium nuclei with $N \le 126$ have been studied by several groups and their findings have been summarized by Gauvin *et al.*¹ For various reasons, none of these groups incorporate observations of the electron capture + positron decay (EC + β^{+}) properties of these nuclei. We have reexamined the decays of $^{205-208}$ Fr and, by observing both the intensities of the α and (EC + β^{+}) decay modes, have deduced their α branching ratios in a manner which is more straightforward than the methods used by previous investigators.

In the present investigation, evidence was also obtained which indicates the existence of a previously unobserved α -emitting isomer in ²⁰⁶Fr. Gamma-ray data taken simultaneously with the α data provide clues to levels within ²⁰⁶Fr and its daughter ²⁰²At.

The francium nuclei of interest were produced by bombarding natural iridium with 120-MeV ²⁰Ne ions accelerated in the Oak Ridge isochronous cyclotron. The reaction products were mass separated on line at the University Isotope Separator at Oak Ridge (UNISOR) facility. The separator ion sources were operated in a surfaceionization mode to release exclusively the francium product nuclei. A high temperature ion source² was used for most of the experiments.

The mass-separated francium nuclei were assayed with surface barrier Si(Au) α particle detectors and large volume Ge(Li) γ -ray detectors positioned in absolutely calibrated geometries. Energy and absolute efficiency calibrations were made with standard sources of known strength. The α and γ data were taken simultaneously in a spectrum multiscale mode using the UNISOR data acquisition system described previously.³ A precision pulser was used to provide time normalization for the multiscale data.

The francium α -branching ratios were based on α intensities and the intensities of the firstexcited to ground state γ transitions in the (EC + β^+) daughters. In parallel investigations, ^{4, 5} no transitions to the ground states were observed to bypass the first-excited states. Additionally, direct (EC + β^+) feeding to the radon ground states is unlikely since the spin difference between the parent and daughter is in all cases $\geq 2\hbar$. Thus the first-excited to ground state transition intensities represent the (EC + β^+) decay strengths.

23

2342

© 1981 The American Physical Society



FIG. 1. Alpha spectrum measured for A = 206 nuclides. The spectrum shows a previously unreported α group at 6.930 ± 0.005 MeV. The remaining peaks are due to the α decay of ²⁰⁶Fr and its daughter ²⁰²At.

The $A = 206 \alpha$ multiscale spectra revealed the presence of a new α group at 6.930 ± 0.005 MeV (see Fig. 1) which decayed with a half-life of 0.7 ± 0.1 s. Based on the following arguments, we identify it as the α decay of an isomer in ²⁰⁶Fr:

(a) no element with Z greater than 87 (francium) can be produced in the heavy-ion reaction used;

(b) the decay curve of the group was that characteristic of a first order (parent) decay equation rather than that of a higher order (daughter) curve; and

(c) radon, astatine, and polonium atoms do not surface ionize to an appreciable extent, requiring any of their isotopes deposited as sources in these experiments to be daughters of francium decay.

Analysis of the γ -ray data revealed two transitions of energy 391 and 531 keV with half-lives $(1.1\pm0.4 \text{ and } 0.7\pm0.1 \text{ s}, \text{ respectively})$ nearly equal to that of the 6.930 MeV α group. The correspondence of the half-lives and energy differences between these γ rays and the two α groups of ²⁰⁶Fr indicate that they are transitions within ²⁰⁶Fr and ²⁰²At. Their suggested placements in the two nuclei are indicated by the level schemes shown in Fig. 2. Gamma-ray coincidence data taken in these experiments possessed insufficient



FIG. 2. Proposed α -decay schemes for ²⁰⁶Fr and ²⁰⁶Fr^m. Level and transition energies are given in MeV.

statistical significance either to support or detract from such a proposal. Therefore, since no excited levels within ²⁰²At or ²⁰⁶Fr are known at this time, further experiments are needed to confirm the proposed level structures.

Table I summarizes our determinations of the ^{205,206,207,208}Fr branching ratios, half-lives, and α -decay energies. The interested reader is referred to the compilation of Gauvin $et \ al.^1$ for a detailed comparison with currently available data^{6,7,8} for the same four nuclides. We will only mention that within error limits our α -decay energies agree with the published values. The same is true for half-lives with the exception of those given in Ref. 6 for 208 Fr and 207 Fr, i.e., 37.5 ± 2.0 and 18.7 ± 0.8 s, respectively. Only Hornshöj et al.⁸ have measured α branches. Their values, $74 \pm 3\%$, $93 \pm 3\%$, $85 \pm 2\%$, and $\ge 97\%$ for ²⁰⁸Fr, ²⁰⁷Fr, ²⁰⁶Fr, and ²⁰⁵Fr, respectively, are generally lower than our ratios (see Table I). Perhaps these discrepancies are due to differences in the experimental techniques. Their method⁸ is based on the observed intensities of the francium and daughter (astatine) α groups. A correction must be made for recoil losses following α decay and the daughter α branches have to be known. Neither the correction nor the daughter branching ratios are required when using our method.

TABLE I. Summary of results obtained for ²⁰⁵⁻²⁰⁸Fr.

| Nuclide | E_{α} (MeV) | lpha-decay branching ratio (%) | Half-life (s) |
|---------------------|--------------------|-----------------------------------|-----------------|
| ²⁰⁸ Fr | 6.636 ± 0.005 | 90±4 | 59.1 ± 0.3 |
| ²⁰⁷ Fr | 6.766 ± 0.005 | $97\pm^{2}_{3}$ | 14.9 ±0.1 |
| 206 Fr ^m | 6.930 ± 0.005 | 0.3 ± 0.1 | 0.7 ±0.1 |
| ²⁰⁶ Fr | 6.790 ± 0.005 | 93±4 | 15.9 ±0.3 |
| ²⁰⁵ Fr | 6.917 ± 0.005 | ~100 ^a | 3.96 ± 0.04 |

^aNo evidence of (EC+ β^*) decay for ²⁰⁵Fr was observed—upper limit ~1%.

The information obtained for 206 Fr^m is also included in Table I. Its α -branching ratio of (0.3 \pm 0.1)% was deduced by comparing the intensities of the 206 Fr and 206 Fr^m α groups and by using the experimentally determined 206 Fr α branch.

Alpha-decay rates are customarily considered within a theoretical framework in which relative probabilities can be obtained after removing the dependence on Q value as well as on the atomic and mass numbers. One convenient formalism has been developed by Rasmussen⁹ wherein a reduced width δ^2 is defined as $\delta^2 = \lambda h/P$. In the equation, λ is the decay constant, h is Planck's constant, and P is the penetrability factor. The barrier for the α particle to tunnel through includes an optical-model potential; a centrifugal barrier is also added so that l dependence in α decay can be considered. In α decay, transitions between ground states of doubly-even nuclei are taken to represent unhindered decays. The reduced widths of these s-wave transitions are taken to be standard. Hindrance factors for other α decays are then deduced by comparing their widths with those of neighboring s-wave transitions.

The δ^2 values for neutron-deficient francium isotopes with $N \le 130$ were calculated by using half-lives, branching ratios, and decay energies from Table I for ${}^{205-208}$ Fr, and Ref. 1 for the remaining nuclides. Angular momentum transfers were deduced either from experimental spins¹⁰ or from inferences based on systematics. The calculations showed that the δ^2 values, when considered as a function of neutron number, have a sharp minimum at N=126. This effect has long been observed not only for francium nuclei (Ref. 8), but for other heavy elements as well.^{8,11} The pronounced dip has been explained¹² as being due to the influence of the major N=126 closed shell. In addition, the widths for both odd-N and even-Nfranciums were found to be similar in magnitude to those of neighboring s-wave transitions for polonium, radon, and radium nuclei (Ref. 11). The francium nuclides therefore have unhindered α -decay rates. This lack of hindrance reflects¹² a large overlap between the wave functions of the initial and final states connected by the α decay. It would appear that the levels in the francium parents and astatine daughters connected by each of the α transitions considered have the same spin and parity.

This research was sponsored in part by the U. S. Department of Energy. Oak Ridge National Laboratory is operated by Union Carbide Corporation for the Department of Energy under Contract No. W-7405-eng-26. UNISOR is a consortium of thirteen institutions and is partly supported by them and by the Department of Energy under Contract No. DE-AC05-760R00033 with the Oak Ridge Associated Universities. One of us (BGR) was supported by an appointment to the Laboratory Graduate Participant program administered by Oak Ridge Associated Universities for the Department of Energy.

- ¹H. Gauvin, Y. Le Beyec, J. Livet, and J. L. Reyss, Ann. Phys. (Paris) <u>9</u>, 241 (1975).
- ²R. L. Mlekodaj, E. H. Spejewski, and B. G. Ritchie, Nucl. Instrum. Methods 171, 451 (1980).
- ³H. K. Carter *et al.*, Nucl. Instrum. Methods <u>139</u>, 349 (1976).
- ⁴B. G. Ritchie, F. T. Avignone, H. K. Carter, R. L. Mlekodaj, and E. H. Spejewski, Phys. Rev. C <u>23</u>, 1717 (1981).
- ⁵B. G. Ritchie, Ph.D. dissertation, University of South Carolina, 1979 (unpublished).
- ⁶R. D. Griffioen and R. D. Macfarlane, Phys. Rev. <u>133</u>, 1373 (1964).

- ⁷K. Valli, E. K. Hyde, and W. Treytl, J. Inorg. Nucl. Chem. 29, 2503 (1967).
- ⁸P. Hornshöj, P. G. Hansen, and B. Jonson, Nucl. Phys. A230, 380 (1974).
- ⁹J. O. Rasmussen, Phys. Rev. <u>113</u>, 1593 (1959).
- ¹⁰Table of Isotopes, 7th ed., edited by C. Michael Lederer and Virginia S. Shirley (Wiley, New York, 1978), pp. 1279-1370.
- ¹¹K. S. Toth, M. A. Ijaz, C. R. Bingham, L. L. Riedinger, H. K. Carter, and D. C. Sousa, Phys. Rev. C <u>19</u>, 2399 (1979).
- ¹²H. J. Mang, Annu. Rev. Nucl. Sci. <u>14</u>, 1 (1964).