

Mass and excited states of <sup>100</sup>Nb

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The <sup>100</sup>Mo(*t*, <sup>3</sup>He)<sup>100</sup>Nb reaction has been used to determine the mass and the energy levels of <sup>100</sup>Nb. A mass excess of  $-79480 \pm 30$  keV is obtained for <sup>100</sup>Nb, in substantial disagreement with  $\beta$ -decay results. Twenty-two excited states of <sup>100</sup>Nb with  $E_x < 1.3$  MeV have been observed.

[NUCLEAR REACTION <sup>100</sup>Mo(*t*, <sup>3</sup>He),  $E=23.0$  MeV; measured  $\sigma(E_{\text{He}}, \theta)$ . <sup>100</sup>Nb] deduced levels. New mass of <sup>100</sup>Nb.

<sup>100</sup>Nb has previously<sup>1</sup> been observed in studies of the <sup>100</sup>Mo(*n*, *p*)<sup>100</sup>Nb reaction, in the  $\beta^-$  decay of <sup>100</sup>Zr and as a fission by-product. Two isomers of <sup>100</sup>Nb have been identified: Their half-lives<sup>2</sup> are  $1.5 \pm 0.2$  s and  $3.1 \pm 0.3$  s. Studies<sup>2-4</sup> of the  $\gamma$  rays in <sup>100</sup>Mo, formed in the  $\beta^-$  decay of these two isomers, show that the 1.5-s isomer should have a low spin since it decays to 0<sup>+</sup> and 2<sup>+</sup> states of <sup>100</sup>Mo. This is confirmed by the fact that it is the only one of the two isomeric states populated in the decay of <sup>100</sup>Zr ( $J^\pi=0^+$ ). The decay of the 3.5-s isomeric state is poorly known but its population in fission, the fact that it is not populated in the <sup>100</sup>Zr  $\beta^-$  decay, and its probable<sup>2</sup>  $\beta^-$  decay to a 4<sup>+</sup> state in <sup>100</sup>Mo at 1135 keV suggest a large spin.  $Q_{\beta^-} = 6240 \pm 100$  keV from the decay<sup>3</sup> of the 1.5-s isomeric state. Wapstra and Bos<sup>5</sup> adopt  $6230 \pm 130$  keV. It is not known which of the two isomeric states is the ground state or what their separation is. No other states of <sup>100</sup>Nb have previously been observed.

The masses and the low-lying states of <sup>100</sup>Nb have been studied by the (*t*, <sup>3</sup>He) reaction on <sup>100</sup>Mo. A 23-MeV triton beam from the LASL three-stage Van de Graaff facility and a magnetic spectrometer of the quadrupole-dipole-dipole-dipole (Q3D) type which has a focal plane detector consisting of a 1-m long helix detector with 0.8-mm spatial resolution<sup>6</sup> were used.

A self-supporting target<sup>7</sup> of molybdenum enriched<sup>8</sup> to 95.9% <sup>100</sup>Mo (also containing 1.7% <sup>98</sup>Mo) was oriented at 20° to the incident triton beam. The target was 338  $\mu\text{g}/\text{cm}^2$  thick. Data were taken with the <sup>100</sup>Mo target at  $\theta = 25^\circ, 30^\circ,$  and  $35^\circ$  and total integrated beam currents of 3.16 to 4.32 mC. Runs were also made with an enriched <sup>24</sup>Mg target under identical conditions preceding and following each run to calibrate<sup>9</sup> the channel number versus

the energy of the outgoing <sup>3</sup>He ions.

Figure 1 shows spectra obtained at  $\theta_{\text{lab}} = 25^\circ$  and  $35^\circ$ , and a partial spectrum of the results at  $30^\circ$ . The numbered groups correspond to states in <sup>100</sup>Nb: See Table I. The resolution of single groups (full width at half maximum  $\approx 25$  keV) observed under similar conditions in other experiments<sup>9,10</sup> is clearly not sufficient to resolve the many states of <sup>100</sup>Nb with  $E_x < 1.4$  MeV. The 23 groups displayed in Table I are a lower limit to the states of <sup>100</sup>Nb in that range. States  $\leq 20$  keV apart would not be resolved, nor could we resolve

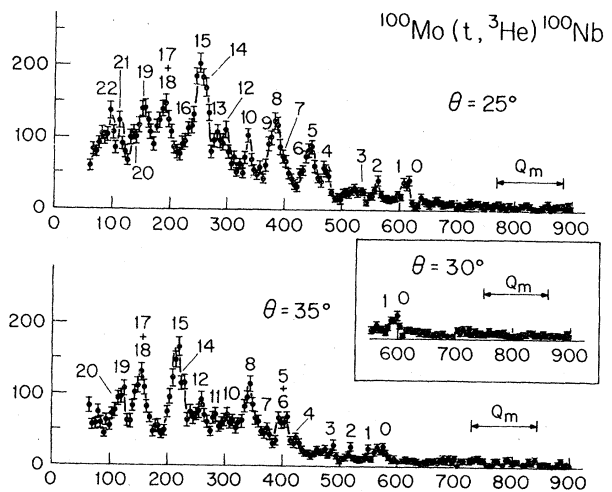


FIG. 1. Spectra of the <sup>3</sup>He ions from the <sup>100</sup>Mo(*t*, <sup>3</sup>He)<sup>100</sup>Nb reaction at  $E_t=23.0$  MeV,  $\theta_{\text{lab}}=25^\circ$  and  $35^\circ$ ,  $B=5.4532$  kG. The ordinate shows the total number of counts recorded in a 5-channel bin. The abscissa shows the channel number. The inset shows the region corresponding to  $-6.0 < E_x < -6.8$  MeV at  $30^\circ$ . For a discussion of the regions labeled  $Q_m$  see the text. The numbered groups are due to states in <sup>100</sup>Nb: See Table I.

TABLE I. Energy levels of  $^{100}\text{Nb}$ .

Group No. <sup>a</sup>	$E_x$ in $^{100}\text{Nb}$ (keV)	$d\sigma/d\Omega^b$ ( $\mu\text{b}/\text{sr}$ )
0	0 <sup>c</sup>	0.090
1	25 ± 10	0.093
2	131 ± 10	0.080
3	(210 ± 15) <sup>d</sup>	
4	348 ± 15	0.19
5	410 ± 15	
6	450 ± 20 <sup>e</sup>	
7	520 ± 20 <sup>e</sup>	
8	565 ± 10	
9	(595 ± 20) <sup>e</sup>	
10	680 ± 20	0.28
11	(720 ± 20) <sup>f</sup>	
12	784 ± 20	
13	820 ± 20 <sup>f</sup>	
14	865 ± 20 <sup>e</sup>	
15	893 ± 20	
16	945 ± 20 <sup>e</sup>	
17	1040 ± 20 <sup>e</sup>	
18	1075 ± 20 <sup>e</sup>	
19	1136 ± 20 <sup>f</sup>	
20	1180 ± 25 <sup>f</sup>	
21	1260 ± 30 <sup>g</sup>	
22	1300 ± 30 <sup>g</sup>	

<sup>a</sup> See Fig. 1.<sup>b</sup>  $\theta_{\text{lab}} = 25^\circ; \pm 40\%$ .<sup>c</sup>  $Q_0$  measured in this experiment is  $-6.690 \pm 0.030$  MeV.<sup>d</sup> Observed at only one angle. There is weak structure at all three angles suggesting unresolved states with  $200 < E_x < 300$  keV.<sup>e</sup> Not resolved.<sup>f</sup> Resolved at one angle.<sup>g</sup> Kinematically observable at only one angle.

states which are weakly populated if in close proximity to strong groups.

The mass of the ground state of  $^{100}\text{Nb}$  derived from our results is  $99.914675$  (32) u, using  $Q_0$

$= -6690 \pm 30$  keV as measured from our data, the masses of  $^{100}\text{Mo}$ ,  $t$  and  $^3\text{He}$  from Wapstra and Bos<sup>5</sup> as well as the conversion factor  $931.5016$  (26) MeV/u used by them. The atomic mass excess of  $^{100}\text{Nb}$  is then  $-79480$  (30) keV, yielding  $Q_{\beta^-}(\text{max}) = 6709 \pm 30$  keV. This is puzzling in view of the fact that Stippler *et al.*<sup>3</sup> found  $Q_{\beta^-}$  for the 1.5-s, low spin isomeric state of  $^{100}\text{Nb}$  to be  $6240 \pm 100$  keV. We show in Fig. 1 the region at all three angles where a group corresponding to  $Q_{\beta^-} = 6230 \pm 130$  keV (the Wapstra-Bos value<sup>5</sup>) should be observed. No structure is evident at any of these angles, which is particularly surprising in view of the fact that the 1.5-s isomer is a low spin state. The relatively constant background for the channels above group 0 can be easily understood in several ways and may be in part due to the presence of counts from the  $^{98}\text{Mo}(t, ^3\text{He})^{98}\text{Nb}$  reaction ( $^{98}\text{Mo}$  was a 1.7% contaminant) whose ground state  $Q$  value is<sup>5</sup>  $-4.566$  MeV. However, the more likely explanation is leak-through of  $\alpha$  particles into the  $^3\text{He}$  spectrum since the  $(t, \alpha)$  reaction is three orders of magnitude more prolific than is the  $(t, ^3\text{He})$  reaction. These counts, from  $^{100}\text{Mo}(t, \alpha)^{99}\text{Nb}$ , would appear as a continuum due to the high excitation energy in the residual nucleus  $^{99}\text{Nb}$ . It would be interesting to see if additional  $\beta^-$  decay studies would confirm  $Q_{\beta^-}$ . Studies of the branching ratios of the decay of both isomeric states would permit limits to be placed on the  $J^\pi$  of the two isomers, and might suggest their ordering in  $^{100}\text{Nb}$  in conjunction with the results presented here.

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