

TABLE I. Mean shifts, attenuation factors, and mean lifetimes for transitions deexciting the nuclear levels of ^{64}Zn . The uncertain weak transitions are indicated inside parentheses.

Level (keV)	E_γ (keV)	I_γ rel. ^b	ΔE_γ (keV)	E_α ^a (MeV)	$F(\tau)$	τ (fs)
991.5	991.5	3800	0.13 ± 0.03	6.4	0.025 ± 0.007	$2700 \pm \begin{smallmatrix} 800 \\ 500 \end{smallmatrix}$
1799.4	807.9	1000	0.12 ± 0.03	7.0	0.026 ± 0.007	$2600 \pm \begin{smallmatrix} 800 \\ 500 \end{smallmatrix}$
	1799.5	390	0.26 ± 0.07		0.026 ± 0.007	$2600 \pm \begin{smallmatrix} 800 \\ 500 \end{smallmatrix}$
1910.3	918.8	150	0.10 ± 0.03	7.0	0.020 ± 0.006	$3400 \pm \begin{smallmatrix} 1500 \\ 800 \end{smallmatrix}$
2306.7	1315.2	660	0.88 ± 0.10	7.6	0.114 ± 0.013	$630 \pm \begin{smallmatrix} 140 \\ 100 \end{smallmatrix}$
2609.3	1617.8	88	0.50 ± 0.18	7.6	0.053 ± 0.020	$1400 \pm \begin{smallmatrix} 800 \\ 500 \end{smallmatrix}$
2736.5	937.1	250	0.14 ± 0.04	7.6	0.026 ± 0.008	$3000 \pm \begin{smallmatrix} 1200 \\ 700 \end{smallmatrix}$
	1745	≈ 25				
2793.9	1802.4	200	9.9 ± 0.5	7.6	> 0.90	< 13
2980.2	1180.8	100	< 0.20	8.0	< 0.03	> 2500
	1988.5	63	< 0.20		< 0.02	> 3800
3005.7	1206.8	45	3.8 ± 0.5	8.0	0.53 ± 0.07	$81 \pm \begin{smallmatrix} 25 \\ 17 \end{smallmatrix}$
	2013.7	75	9.0 ± 2.0		0.75 ± 0.17	$32 \pm \begin{smallmatrix} 30 \\ 23 \end{smallmatrix}$
	3005.7	50	11.0 ± 2.0		0.60 ± 0.12	$65 \pm \begin{smallmatrix} 20 \\ 15 \end{smallmatrix}$
2998.3	2006.8	165	< 0.50	8.0	< 0.040	> 1500
3078.3	771.5	60	0.20 ± 0.10	8.0	0.040 ± 0.020	$2000 \pm \begin{smallmatrix} 1500 \\ 800 \end{smallmatrix}$
	2086.8	60	0.38 ± 0.15		0.031 ± 0.012	$2500 \pm \begin{smallmatrix} 1400 \\ 800 \end{smallmatrix}$
3094.3	2103.8	97	3.9 ± 0.5	8.0	0.31 ± 0.04	$190 \pm \begin{smallmatrix} 40 \\ 30 \end{smallmatrix}$
3187.0	(1276)	≈ 10		8.0		
	1387.6	60	1.0 ± 0.3		0.12 ± 0.04	$580 \pm \begin{smallmatrix} 300 \\ 180 \end{smallmatrix}$
	(2195)	≈ 15				
3206.5	1407.1	68	1.20 ± 0.30	8.0	0.15 ± 0.04	$470 \pm \begin{smallmatrix} 200 \\ 120 \end{smallmatrix}$
	(2214)	≈ 10				
3261.9	(1462)	≈ 20		8.0		
	2270.4	40	11.4 ± 1.0		0.85 ± 0.08	$20 \pm \begin{smallmatrix} 11 \\ 12 \end{smallmatrix}$
3296.8	2305.3	65	2.1 ± 0.3	8.0	0.15 ± 0.02	$450 \pm \begin{smallmatrix} 100 \\ 80 \end{smallmatrix}$
(3365)	(1566)	≈ 15		8.0		
	(2374)	≈ 10				
	3364	20	15 ± 3		0.75 ± 0.15	$37 \pm \begin{smallmatrix} 23 \\ 22 \end{smallmatrix}$
(3425)	(3425)	≈ 15	20 ± 3	8.0	> 0.90	< 15

^a Energy of the incident $^4\text{He}^{++}$ beam.

^b Relative intensities at $E_\alpha = 8$ MeV.

^{64}Zn and the corresponding γ -ray transitions. The proposed level scheme is shown in Fig. 2. Up to the level at 2793 keV the spin-parity assignments are obtained from previous distorted wave Born approximation analysis^{1,2} or angular correlation measurements.³ It should be noted that the levels at 2609 and 2793 keV are possibly connected to the 2^+ 1799 keV level. As a matter of fact, those tran-

sitions (indicated by dashed lines) may have been masked by the strong transitions of 807.9 and 991.5 keV, respectively. However, a careful analysis of those peaks indicated that the possible 809.9 and 994.5 keV transitions must have intensities lower than 40 and 140 (units of Table I), respectively. States with $J^\pi = 3^-$ and 4^+ have been predicted at about 3.02 and 3.11 MeV, respectively, using the

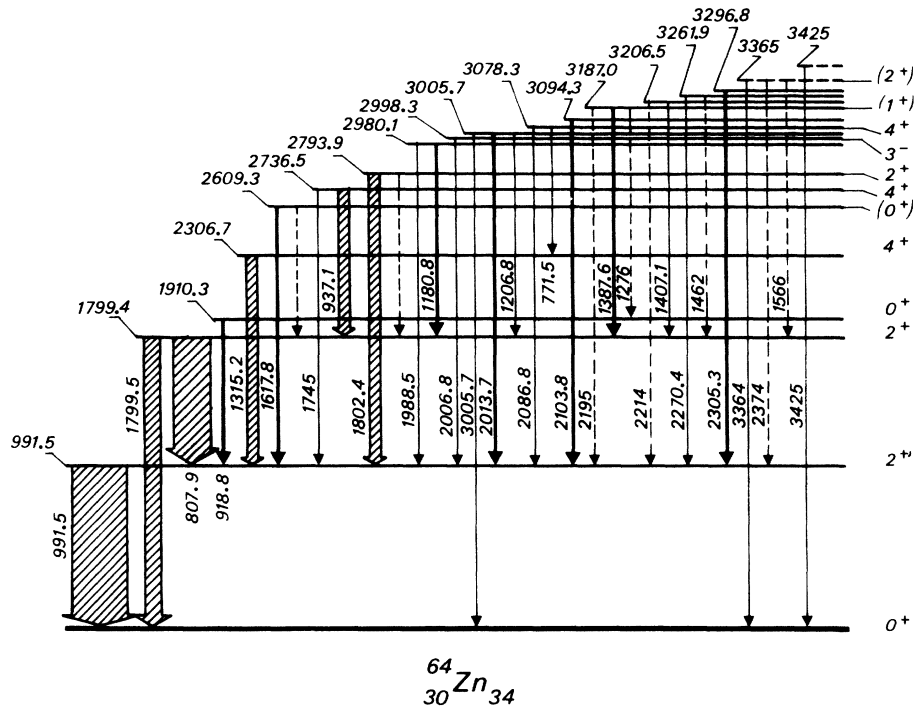


FIG. 2. Proposed level scheme of ^{64}Zn resulting from the present investigation.

($\alpha\alpha'$) and (t, p) reaction data.^{1,2} Those states were finally associated with the 2998 and 3078 keV levels on the evidence of γ - γ angular correlation experiments.³ Unfortunately, no evidence is available for the assignments of the 2980 and 3006 keV levels which were also observed via the (p, p') reaction. The high lying levels of our scheme are also excited with the (p, p') reaction. The very weak or uncertain transitions are indicated by dashed lines.

In the following columns of Table I are given the

TABLE II. Electromagnetic transition rates for some transitions in ^{64}Zn . The reduced transition probabilities are given in Weisskopf units.

Level	I^π	Transition	Mult.	T_{exp}/T_W
991.5	2^+	991.5	$E2$	20.4
1799.4	2^+	807.9	$E2$	42.3
		1799.4	$E2$	0.30
1910.3	0^+	918.8	$E2$	23.6
2306.7	4^+	1315.2	$E2$	21.3
2609.3	0^+	1617.8	$E2$	3.4
2736.5	4^+	937.1	$E2$	21.8
		1745	$E2$	0.10
2793.9	2^+	1802.4	$E2$	>213
		1802.4	$M1$	>0.45
2998.3	3^-	2006.8	$E1$	$<5.2 \times 10^{-5}$
3078.3	4^+	771.5	$E2$	42.5
		2086.8	$E2$	0.31

measured $F(\tau)$ and the deduced mean lifetimes for each level. It should be noted that the lifetime of the first excited state is in good agreement with previous measurements^{4,7} ($T_{1/2} = 1.75$ ps, $\tau = 2.52$ ps). Those results are discussed in the next section.

IV. DISCUSSION

From the measured branching ratios, the level lifetimes presented in Table I, and the expected multiplicities, we can deduce the reduced transition probabilities. Their values in Weisskopf units (W.u.) are given in Table II for some ^{64}Zn states. The $B(E2)$ value of the one-phonon state is about the same as in ^{66}Zn and ^{68}Zn .⁷⁻⁹ The $B(E2)$ values of the levels at 1799 (2^+), 1910 (0^+), and 2306 (4^+) are in the range expected for the two-phonon state. Those results can be compared to the values deduced from the lifetimes measured in ^{66}Zn .⁹ The $B(E2)$ corresponding to the 4^+ state is of the same order of magnitude (43 W.u.) but there is a strong deviation for the 2^+ state due to the surprising large $B(E2)$ value in ^{66}Zn (468 W.u.). A three-phonon state interpretation can be suggested for the 2736 keV state (4^+) which decays primarily to the 2^+ two-phonon level. The very short lifetime of the 2793 keV level should be noted. The reduced transition probability values

suggest that the 1802 keV transition is not pure $E2$ but rather ($E2 + M1$). The $B(E1)$ value of the 2998 keV level is lower by at least a factor of 10 than those deduced from measured lifetimes of the first 3^- states in ^{66}Zn and ^{68}Zn .⁹

In conclusion, it seems difficult at present to achieve a theoretical understanding of this experiment. The calculation of sufficiently realistic wave functions for Zn isotopes is now required to make full use of the experimental data.

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