

**$\nu i_{13/2}$  and  $\nu h_{9/2}$  isomers in odd-*A* Pt nuclei\***M. Piiparinen,<sup>†</sup> S. K. Saha, and P. J. Daly*Chemistry Department, Purdue University, West Lafayette, Indiana 47907*

C. L. Dors, F. M. Bernthal, and T. L. Khoo

*Departments of Chemistry and Physics and Cyclotron Laboratory,  
Michigan State University, East Lansing, Michigan 48824*

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$\frac{13}{2}^+$  isomers with half-lives of  $311 \pm 15$ ,  $143 \pm 5$ , and  $95 \pm 5$   $\mu$ s in  $^{187}\text{Pt}$ ,  $^{189}\text{Pt}$ , and  $^{191}\text{Pt}$ , respectively, have been observed in  $(\alpha, 3n\gamma)$  reactions on isotopically enriched Os targets. Isomeric decay schemes based on the present measurements and on  $^{189}\text{Au}$  and  $^{191}\text{Au}$  decay data are presented. It is proposed that the  $^{187,189,191}\text{Pt}$   $\frac{13}{2}^+$  isomers deexcite by *M2* transitions to  $\frac{9}{2}^-$  intrinsic states of  $\nu h_{9/2}$  character. In  $A \geq 193$  Pt nuclei the corresponding  $\frac{9}{2}^-$  states must lie above the known  $\nu i_{13/2}$  isomers, which have much longer half-lives.

[ NUCLEAR REACTIONS  $^{186,188,190,192}\text{Os}(\alpha, 3n\gamma)$ ,  $E=30-50$  MeV, enriched targets; measured  $\sigma(E)$ ,  $E_\gamma$ ,  $I_\gamma$ ,  $\gamma-\gamma$  coin,  $\sigma(\theta)\gamma$ ,  $T_{1/2}$ , deduced  $^{187}\text{Pt}$ ,  $^{189}\text{Pt}$ ,  $^{191}\text{Pt}$  isomeric decay schemes,  $J$ ,  $\pi$ . ]

## I. INTRODUCTION

In a recent study<sup>1</sup> of the high-spin level systematics of  $^{186-194}\text{Pt}$  by  $(\alpha, xn\gamma)$  reactions on isotopically enriched Os targets, we identified decoupled  $\nu i_{13/2}$  bands in the four odd-*A* nuclei  $^{187}\text{Pt}$ ,  $^{189}\text{Pt}$ ,  $^{191}\text{Pt}$ , and  $^{193}\text{Pt}$ . The  $\frac{13}{2}^+$  bandheads which were observed to be very strongly populated in the  $(\alpha, 3n)$  reactions are low-lying isomers and not the ground states of these nuclei. The present paper describes the properties and modes of deexcitation of these  $\frac{13}{2}^+$  isomers. Very recently, the ground state spins of  $^{187}\text{Pt}$ ,  $^{189}\text{Pt}$ , and  $^{191}\text{Pt}$  have been determined to be  $\frac{3}{2}$  by the atomic beam magnetic resonance method.<sup>2</sup>

In  $^{193}\text{Pt}$  and  $^{195}\text{Pt}$ ,  $\frac{13}{2}^+$  *M4* isomers with half-lives of 4.3 and 4.2 days, respectively, have been known for several years.<sup>3</sup> No isomers with comparable half-lives have been found in the lighter odd-*A* Pt nuclei, but there have been reports of shorter lived isomers which are pertinent to the present investigation. In the  $\text{Ir}(p, 2n)$  reaction, Conlon<sup>4</sup> identified a 107- $\mu$ s  $^{191}\text{Pt}$  isomer deexciting by a 91-keV *E2* transition. In similar proton bombardments of natural iridium, Fraser and Moore<sup>5</sup> observed  $\gamma$  rays of 167 and 90 keV decaying with half-lives of 145 and 95  $\mu$ s, respectively, and tentatively assigned them to new  $^{190}\text{Pt}$  and  $^{192}\text{Pt}$  isomers. These assignments are shown on the most recent Chart of the Nuclides.<sup>6</sup> In studies of the radioactive decay of high-spin  $^{189}\text{Au}^m$ , the ISOLDE collaboration<sup>7</sup> identified a  $^{189}\text{Pt}$  isomer decaying by a 167-keV *E2* transition with a half-life of 464 ns.

## II. EXPERIMENTAL PROCEDURE AND ANALYSIS

The experimental methods used were very similar to those described in the preceding paper.<sup>8</sup> Targets of isotopically enriched  $^{186}\text{Os}$ (62%),  $^{188}\text{Os}$ (87%),  $^{190}\text{Os}$ (95%), and  $^{192}\text{Os}$ (98%) were bombarded with 30–50 MeV  $\alpha$  particles from the Michigan State University cyclotron and  $\gamma$ -ray singles,  $\gamma$ - $\gamma$  coincidence, angular distribution, and lifetime measurements were performed. Isotopic assignments of individual  $\gamma$  rays to  $^{187}\text{Pt}$ ,  $^{189}\text{Pt}$ ,  $^{191}\text{Pt}$ , and  $^{193}\text{Pt}$  were based primarily on excitation function determinations. When the  $\gamma$ - $\gamma$  coincidence data for all four nuclei had been completely analyzed, no strong  $\gamma$  rays of  $^{193}\text{Pt}$  remained unplaced in the level scheme. However, intense  $\gamma$  rays of 91.1, 166.7, and 117.3 keV which the excitation function measurements quite clearly showed to be due to  $^{191}\text{Pt}$ ,  $^{189}\text{Pt}$ , and  $^{187}\text{Pt}$ , respectively, remained unplaced. These  $\gamma$  rays did not appear in the prompt coincidence spectra, their angular distributions were isotropic and, in beam-sweeping lifetime measurements, no decay of the lines was detectable during a 500-ns period. Therefore longer lifetime measurements were undertaken using a beam pulsing system which greatly extended the range of measurable lifetimes. For the isomers under study here, it was found suitable to allow beam on target every millisecond for 200  $\mu$ s, followed by an 800- $\mu$ s beam-off period. In Fig. 1, the decay data obtained and the half-lives determined for the three isomers are shown.

It is clear from these results that the 145- and

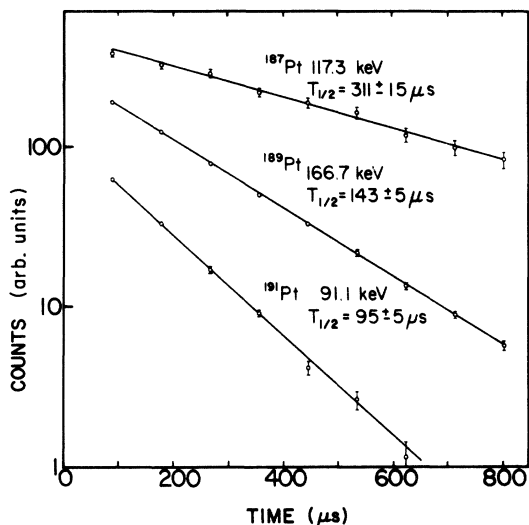


FIG. 1. The decay of the 117-, 167-, and 91-keV  $\gamma$  rays in  $^{187}\text{Pt}$ ,  $^{189}\text{Pt}$ , and  $^{191}\text{Pt}$ , respectively.

95- $\mu\text{s}$  activities reported by Fraser and Moore<sup>5</sup> should be reassigned to  $^{189}\text{Pt}$  and  $^{191}\text{Pt}$ , respectively. In addition, the 95- $\mu\text{s}$   $^{191}\text{Pt}$  91-keV  $\gamma$ -ray activity is obviously the same as the 107- $\mu\text{s}$   $^{191}\text{Pt}$  activity reported by Conlon.<sup>4</sup>

At first, the 143- $\mu\text{s}$  half-life determined for the 167-keV  $\gamma$ -ray activity in  $^{189}\text{Pt}$  was difficult to reconcile with the 0.5- $\mu\text{s}$  half-life observed in the  $^{189}\text{Au}^m$  decay study.<sup>7</sup> However, when the  $^{189}\text{Pt}$  coincidence tapes were re-sorted with a time gate set on a 200–400 ns delayed portion of the time-to-amplitude converter spectrum and an energy gate on the 167-keV peak, several  $\gamma$  rays including a moderately strong 320-keV line were found to be in delayed coincidence with the 167-keV  $\gamma$  ray. This result is compatible with the observation in the radioactivity study of a 321-keV transition populating the 0.5- $\mu\text{s}$  state. It is apparent from the combined results that the 167-keV  $\gamma$  ray occurs in the decay of the 0.5- $\mu\text{s}$  isomer, which in the  $(\alpha, 3n)$  reaction is fed predominantly in the deexcitation of the 143- $\mu\text{s}$  isomer and relatively weakly by other deexcitation pathways. There seems to be no indication that the 143- $\mu\text{s}$  isomer is populated in the  $^{189}\text{Au}^m$  decay.

In view of these findings, a search in the 200–400 ns range was also made for  $\gamma$  rays in delayed coincidence with the 117- and 91-keV  $\gamma$  rays of  $^{187}\text{Pt}$  and  $^{191}\text{Pt}$ , but none were found. Consequently, a lower limit of 1  $\mu\text{s}$  could be placed on the half-lives of the levels deexcited by these transitions.

### III. ISOMERIC DECAY SCHEMES

#### Nucleus $^{193}\text{Pt}$

The 4.3-day  $\frac{13}{2}^+$  isomer in  $^{193}\text{Pt}$  is well established.<sup>3</sup> Here the  $\frac{13}{2}^+$  state deexcites by a 135-keV  $M4$  transition to a  $\frac{5}{2}^-$  state at 14 keV. No additional isomeric states in  $^{193}\text{Pt}$  were observed in the present work. All the high-spin states which are strongly populated in the  $(\alpha, 3n)$  reaction deexcite by  $\gamma$ -ray cascades leading to the 4.3-day  $\frac{13}{2}^+$  isomeric state.

#### Nucleus $^{191}\text{Pt}$

In the delayed  $\gamma$ -ray spectra, the only radiations observed to decay with the 95- $\mu\text{s}$  half-life were Pt  $K$  x rays and 91.1-keV (strong) and 48.4-keV (weak)  $\gamma$  rays. The 91-keV transition is known to be of  $E2$  character from earlier work.<sup>3,4</sup> From the  $\gamma$ -ray intensities observed in the  $^{190}\text{Os}(\alpha, 3n)^{191}\text{Pt}$  singles spectra, the 91-keV transition intensity was determined to be slightly larger than the sum of the intensities of all the  $\gamma$  rays feeding the  $\frac{13}{2}^+$  isomeric state. By assuming that the 48- and 91-keV transitions occur in cascade in the deexcitation of the 95- $\mu\text{s}$  isomer, a total conversion coefficient of  $(6.0 \pm 0.8) \times 10^2$  was derived for the 48-keV transition from intensity balance requirements, indicating probable  $M2$  multipolarity for this transition.

Important complementary information about the  $^{191}\text{Pt}$  level structure was obtained in a detailed study<sup>9</sup> of the radioactive decay of 3.2-h  $^{191}\text{Au}$  ( $J^\pi = \frac{3}{2}^+$ ). This investigation located many low-spin  $^{191}\text{Pt}$  levels not detectably populated in the  $(\alpha, 3n)$  reaction. In addition, we found that the  $\frac{13}{2}^+$  isomer is populated quite strongly by cascades of  $\Delta J = +1$  and  $\Delta J = +2$  transitions from levels directly fed in the  $\beta$  decay. In earlier conversion electron studies of the  $^{191}\text{Au}$  decay, a 48.4-keV  $M2$  transition in  $^{191}\text{Pt}$  has been firmly established from subshell intensity ratios.<sup>10</sup> All our measurements indicate that this is the transition deexciting the 95- $\mu\text{s}$  isomer.

The proposed isomeric decay scheme is shown in Fig. 2. The placement of the  $\frac{5}{2}^-$  level at 10 keV is strongly supported by our  $^{191}\text{Au}$  decay data.  $J^\pi = \frac{9}{2}^-$  for the 101-keV level is established by the multipolarities of the transitions feeding and deexciting the level. As noted earlier, the delayed coincidence measurements in-beam indicate that the half-life of the  $\frac{9}{2}^-$  level is greater than 1  $\mu\text{s}$ .

#### Nucleus $^{189}\text{Pt}$

In this nucleus, the only radiations observed to decay with the 143- $\mu\text{s}$  half-life were Pt  $K$  x rays and 167-keV  $\gamma$  rays. The 167-keV transition is of

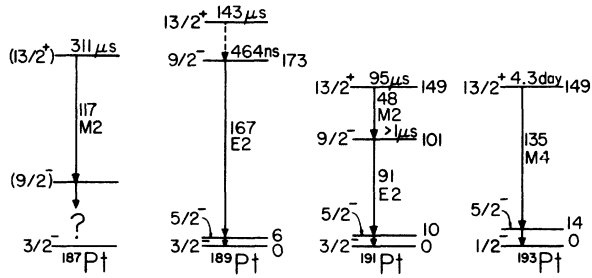


FIG. 2. The proposed isomeric decay schemes in odd- $A$  Pt nuclei. The known  $\frac{3}{2}^-$  level at 1.6 keV in  $^{193}\text{Pt}$  is not shown, although it is populated in the 4.3-day decay.

known  $E2$  character,<sup>3</sup> and its intensity observed in the  $\gamma$ -ray singles measurements was slightly larger than the summed intensities of the  $\gamma$  rays feeding the  $\frac{13}{2}^+$  isomer. The ISOLDE studies<sup>7</sup> of the 4.6-min  $^{189}\text{Au}$  ( $J^\pi = \frac{1}{2}^-$ ) decay have shown that the 167-keV transition occurs in the deexcitation of a 464-ns  $^{189}\text{Pt}$  isomer, with probable spin-parity of  $\frac{9}{2}^-$ . We propose a decay scheme for the 143- $\mu\text{s}$  isomer in  $^{189}\text{Pt}$  which is similar to that for  $^{191}\text{Pt}$  (Fig. 2). In this case the suggested  $M2$  transition has not been observed; in our measurements, a low energy  $M2$  transition ( $E_\gamma < 45$  keV), or one close in energy to the Pt  $K$  x rays, would almost certainly have escaped detection. The level at 6 keV is known from a study<sup>11</sup> of the decay of 28.3-min  $^{189}\text{Au}$  ( $J^\pi = \frac{3}{2}^+$ ), and the  $\frac{5}{2}^-$  spin-parity assignment is proposed because the pattern of its population following the  $\beta$  decay is very similar to that of the 10-keV level in  $^{191}\text{Pt}$ .

#### Nucleus $^{187}\text{Pt}$

Here the only radiations observed to decay with the 311- $\mu\text{s}$  half-life were Pt  $K$  x rays and 117-keV  $\gamma$  rays. By attributing all the delayed Pt  $K$  x rays to internal conversion of the 117-keV transition, a  $K$  conversion coefficient of  $18 \pm 4$  was estimated, indicating  $M2$  multipolarity for the 117-keV transition. The requirement of intensity balance between the  $\gamma$  rays feeding the  $\frac{13}{2}^+$  isomer and the 117-keV transition deexciting it also strongly favors the  $M2$  multipolarity assignment. We suggest that the 117-keV transition may be analogous to the  $\frac{13}{2}^+ \rightarrow \frac{9}{2}^-$  isomeric transitions in  $^{189}\text{Pt}$  and  $^{191}\text{Pt}$ , but in the case of  $^{187}\text{Pt}$  no decay scheme can be proposed because the low-lying level structure is poorly established.

#### IV. DISCUSSION

The decay schemes illustrated in Fig. 2 suggest an explanation for the striking difference between the long half-lives of the  $\nu i_{13/2}$  isomers in odd- $A$  Pt nuclei with  $A \geq 193$  and the submillisecond half-lives of the  $\frac{13}{2}^+$  isomers in  $^{191}\text{Pt}$ ,  $^{189}\text{Pt}$ , and  $^{187}\text{Pt}$ . Our detailed studies have established<sup>1,12</sup> that the  $\frac{13}{2}^+$  levels are definitely the lowest-lying members of the  $\nu i_{13/2}$  level families in  $^{193}\text{Pt}$ ,  $^{191}\text{Pt}$ , and  $^{189}\text{Pt}$ , and the  $\frac{13}{2}^+$  level is probably lowest-lying in  $^{187}\text{Pt}$  also. Therefore the most likely cause of the differences in the  $\frac{13}{2}^+$  half-lives is that a level of moderately high spin occurs below the  $\frac{13}{2}^+$  in the  $A \leq 191$  nuclei, but above it in the heavier nuclei. The results of the present study strongly suggest that this intermediate level in the  $A \leq 191$  nuclei has  $J^\pi = \frac{9}{2}^-$ .

The nature of the  $\frac{9}{2}^-$  levels is of some interest. The reduced transition probabilities for the transitions deexciting the  $\frac{9}{2}^-$  levels in  $^{189}\text{Pt}$  and  $^{191}\text{Pt}$  are  $B(E2; 167 \text{ keV}) = 0.08$  s.p.u. and  $B(E2; 91 \text{ keV}) < 0.15$  s.p.u., so they are obviously not collective excitations built on lower-lying levels. The most revealing evidence regarding these  $\frac{9}{2}^-$  levels is provided by the 4.6-min  $^{189}\text{Au}$  decay data. The 0.5- $\mu\text{s}$   $\frac{9}{2}^-$  state in  $^{189}\text{Pt}$  is populated in the  $\beta$  decay of the  $^{189}\text{Au}$   $\frac{1}{2}^-$  isomer with a  $\log ft$  of 4.7, indicating an allowed unhindered transition.<sup>7</sup> Since the  $\pi h_{11/2}$  character of the 4.6-min  $\frac{1}{2}^-$  state in  $^{189}\text{Au}$  is firmly established,<sup>13,14</sup> the  $^{189}\text{Pt}$   $\frac{9}{2}^-$  state fed in its  $\beta$  decay must be of  $\nu h_{9/2}$  parentage. In the prolate limit ( $\gamma = 0^\circ$ ) this  $\frac{9}{2}^-$  state would be designated  $\frac{9}{2}[505]\nu$ ; however, our results for the  $\nu i_{13/2}$  level family in  $^{189}\text{Pt}$  indicate<sup>12</sup> a triaxial nuclear shape ( $\gamma \sim 26^\circ$ ). Almost certainly the  $\frac{9}{2}^-$  levels in  $^{191}\text{Pt}$  and  $^{187}\text{Pt}$  have a similar nature. It is not surprising to find this  $\frac{9}{2}^-$  intrinsic state at low energies in the light Pt nuclei, since  $\frac{9}{2}[505]\nu$  is known to be the ground state configuration in  $^{191}\text{Os}$ , and the same state has been located<sup>3</sup> 30 keV above ground in  $^{189}\text{Os}$ . Presumably the  $\frac{9}{2}^-$  state lies above the  $\frac{13}{2}^+$  isomers in the  $A \geq 193$  Pt nuclei. In  $^{195}\text{Pt}$ , a  $\frac{9}{2}^-$  level approximately 300 keV above the  $\frac{13}{2}^+$  isomer is fairly strongly populated directly in the  $\beta$  decay<sup>15</sup> of 3.7-h  $^{195}\text{Ir}^m$  ( $J^\pi = \frac{1}{2}^-$ ).

As we have noted in the preceding paper,<sup>8</sup> the existence of low-lying  $\frac{9}{2}^-$  intrinsic states in  $^{189}\text{Pt}$  and  $^{191}\text{Pt}$  appears to be closely related to the occurrence of  $10^-$  isomers of  $\nu i_{13/2}$ ,  $\nu h_{9/2}$  character in the neighboring core nuclei  $^{190}\text{Pt}$  and  $^{192}\text{Pt}$ .

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- †Present address: Department of Physics, University of Jyväskylä, Finland.
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