

## Population of $^{11}\text{O}^*$ in two-neutron removal from $^{13}\text{O}$

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I have used a simple model to estimate the relative populations expected for the ground state (g.s.) and three excited states of  $^{11}\text{O}$  in  $2n$  removal from  $^{13}\text{O}$ . Results are ratios exc./g.s. of at most a few percent for each excited state, compared to ratios near unity suggested in a recent experiment involving  $^{11}\text{O}$ .

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### I. INTRODUCTION

Recently, Webb *et al.* [1] produced  $^{11}\text{O}$  with  $2n$  removal from  $^{13}\text{O}$  and observed it by detecting  $^9\text{C}$  and  $2p$  in coincidence. To reproduce the decay energy spectrum, they found that they needed four states—the ground state (g.s.) and three excited states, whose  $J^\pi$  they took to be  $5/2^+$ ,  $3/2^-$ , and  $5/2^+$ . Here, I attempt to estimate the relative population of these states in a simple model.

### II. CALCULATIONS AND ANALYSIS

Throughout, I treat  $^{11}\text{O}$  and  $^{11}\text{Li}$  as mirrors, and likewise for  $^{13}\text{O}/^{13}\text{B}$  and  $^{11}\text{B}/^{11}\text{C}$ . The main components of the g.s. of  $^{13}\text{B}$  are  $^{13}\text{B}_{1p}$  and  $^{11}\text{B}_{1p} \times \nu(\text{sd})^2$ , so that one may write

$$^{13}\text{O}(\text{g.s.}) = A \ ^{11}\text{C}_{1p} \times \pi(\text{sd})^2 + B \ ^{13}\text{O}_{1p},$$

where the subscripts  $1p$  denote structures totally within the  $1p$  shell. The best evidence for core excitation of the type  $^{11}\text{B} \times \nu(\text{sd})^2$  in  $^{13}\text{B}(\text{g.s.})$  involves the observation [2] of the Gamow-Teller  $\beta$  decay of  $^{14}\text{Be}(\text{g.s.})$  to a  $1^+$  level at 1.28 MeV in  $^{14}\text{B}$ , followed by emission of a neutron to the  $3/2^-$   $^{13}\text{B}(\text{g.s.})$ . The measured neutron width of the  $^{14}\text{B}$  level then allows an estimate of the  $^{13}\text{B}(\text{g.s.})$  configuration mixing. Details of the procedure are given elsewhere [2,3]. Results are listed in Table I, along with one other estimate [4]. Thus, for  $^{13}\text{B}$ , estimates of  $A^2$  range from 0.21 to 0.30 [2–4], with a “best” value of 0.21(2) [3].

The  $^{11}\text{Li}(\text{g.s.})$  has long been treated as a combination of  $^{11}\text{Li}_{1p}$  and  $^9\text{Li}_{1p} \times \nu(\text{sd})^2$ , so that one may use

$$^{11}\text{O}(\text{g.s.}) = a \ ^9\text{C}_{1p} \times \pi(\text{sd})^2 + b \ ^{11}\text{O}_{1p}.$$

In a two-state model, the excited  $3/2^-$  state is just the orthonormal linear combination

$$^{11}\text{O}(3/2^- \text{exc}) = -b \ ^9\text{C}_{1p} \times \pi(\text{sd})^2 + a \ ^{11}\text{O}_{1p}.$$

The latest estimate [5] of  $a^2$  is  $0.33^{+0.03}_{-0.05}$ . Other values are similar [6,7]. Amplitudes for producing these two  $3/2^-$  states in two-neutron removal from  $^{13}\text{O}$  are thus

$$\mathcal{A}(\text{g.s.}) = Aa \ \mathcal{A}(^{11}\text{C}_{1p} \rightarrow ^9\text{C}_{1p}) + Bb \ \mathcal{A}(^{13}\text{O}_{1p} \rightarrow ^{11}\text{O}_{1p});$$

$$\mathcal{A}(3/2^- \text{exc}) = -bA \ \mathcal{A}(^{11}\text{C}_{1p} \rightarrow ^9\text{C}_{1p}) + aB \ \mathcal{A}(^{13}\text{O}_{1p} \rightarrow ^{11}\text{O}_{1p}).$$

Note that all transitions involve removal of a  $p$ -shell neutron pair, while the  $sd$ -shell protons act as spectators. Thus, the expected amplitude ratio is approximately

$$\mathcal{A}(3/2^- \text{exc.})/\mathcal{A}(\text{g.s.}) \sim (-bA + aB)/(Aa + Bb) = 0.135,$$

with an estimated uncertainty of about 24%, or a predicted cross section ratio of about 0.02(1).

The lowest positive-parity state in  $^{11}\text{Li}$  has the structure  $^9\text{Li}_{1p} \times \nu((1p_{1/2})(2s_{1/2}))_{1-}$ . This configuration actually contains three states, with  $J^\pi = 5/2^+$ ,  $3/2^+$ , and  $1/2^+$ . Thus, the configuration of this state in  $^{11}\text{O}$  is  $^9\text{C}_{1p} \times \pi((1p_{1/2})(2s_{1/2}))_{1-}$ . This configuration in  $^{11}\text{O}$  cannot be produced from the  $^{13}\text{O}(\text{g.s.})$  considered above. Rather, to make this configuration from  $^{13}\text{O}$  by two-neutron removal requires an  $sd$ -shell neutron in its g.s. By correspondence with the mirror  $^{13}\text{B}$ , the lowest such admixtures in  $^{13}\text{O}(\text{g.s.})$  would be  $^{11}\text{N}(1/2^-) \times (\text{sd})^2_{10 \text{ or } 21}$ , where the subscripts denote  $JT$  of the  $sd$ -shell pair, which must be one proton and one neutron. In the mirror  $^{13}\text{B}$ , I have argued elsewhere [8] that these admixtures are extremely small, 0.026 and  $<0.036$ , respectively. To make the  $5/2^+$  state from either of these requires removal of two neutrons from two different major shells. Thus, I estimate that, even if  $(s_{1/2})(p_{3/2})$  removal is as strong as  $(p_{3/2})^2$  removal, the  $5/2^+$  strength should be at most a few percent of  $^{11}\text{O}(\text{g.s.})$ .

As mentioned by Webb *et al.*, the second  $5/2^+$  state involves a  $p_{3/2}$  to  $s_{1/2}$  excitation of  $^{11}\text{O}(\text{g.s.})$ . Hence, its population in two-neutron removal from  $^{13}\text{O}(\text{g.s.})$  should be even weaker.

These predictions are listed in Table II, where they are compared with relative yields used in the spectrum fit of Webb *et al.*

I was able to reproduce the experimental  $^{11}\text{O}$  to  $^9\text{C} + 2p$  decay spectrum using only the g.s. of  $^{11}\text{O}$  [9]. The energy and width needed were quite close to those predicted previously [10,11]. Thus, I consider it very unlikely that the  $^{11}\text{O}$  decay spectrum contains excited states at anything near the magnitudes used by Webb *et al.* If the  $\ell$  value(s) of the  $2p$  decay pair can be determined, then the situation could be clarified.

TABLE I. Estimates of  $^{11}\text{B} \times v(\text{sd})^2$  neutron intensity in  $^{13}\text{B}(\text{g.s.})$ .

Intensity	Reference
0.33 <sup>a</sup>	2
0.25(3) <sup>b</sup>	3
0.30(3) <sup>c</sup>	3
0.21(2) <sup>d</sup>	3
0.25(5)	4

<sup>a</sup> $R$  matrix using a  $^{14}\text{B}(1^+)$  width of 49 keV and no reduction for  $S < 1$ .

<sup>b</sup>As above, but with  $S(^{11}\text{B} \rightarrow ^{12}\text{B}(1^+)) = 0.75(5)$ .

<sup>c</sup>Potential model analysis, using width of 49 keV and  $S = 0.75(5)$ .

<sup>d</sup>As above, but using width of 34(3) keV extracted in Ref. [3] from data of Ref. [2].

### III. SUMMARY

In summary, I have estimated the relative yields expected for the g.s. of  $^{11}\text{O}$  and three excited states in  $2n$  removal from

TABLE II. Current predictions compared with relative populations of states included in fit to  $^{11}\text{O} \rightarrow ^9\text{C} + 2p$  spectrum (energies and widths in MeV) [1].

$J^\pi$	Predicted contribution <sup>a</sup>	$E_r^b$	$\Gamma^b$	Contribution by Webb <i>et al.</i> [1]
$3/2^-$	1.0	4.16	1.30	39%
$5/2^+$	0.02–0.04	4.65	1.06	29%
$3/2^-$	0.02(1)	4.85	1.33	
$5/2^+$	<0.02–0.04	6.28	1.96	32%

<sup>a</sup>Present.

<sup>b</sup>Reference [1].

$^{13}\text{O}$ . Ratios exc./g.s. are predicted to be at most a few percent for each excited state, compared to ratios near unity suggested in a recent treatment of  $^{11}\text{O}$ .

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