

### Analog transitions in the $A = 30$ system\*

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The  $\gamma$ -ray decay of the  $2^+$ ,  $T=1$  state at  $E_x = 4.183$  MeV in  $^{30}\text{P}$  was investigated. A new branch ( $2 \pm 1$ )% to the  $0^+$ ,  $T=1$  state at  $E_x = 0.677$  MeV was found. The decay mode is discussed in terms of isospin and an estimate of the lifetime of this state is derived.

[ NUCLEAR REACTION  $^{29}\text{Si}(p, \gamma)$ ,  $E = 0.70\text{--}0.75$  MeV, measured  $\sigma(E; E_\gamma, \theta)$ , deduced  $^{30}\text{P}$  level branching. Enriched  $^{29}\text{Si}$  target. ]

A comparison of analog electromagnetic transitions is a useful test of the validity of selection rules based on the concept of isospin. According to these rules,<sup>1</sup> the reduced strength of a transition connecting two levels of the same isospin can be written as

$$|M_{M\lambda}|^2 = (S + T_x V)^2 \quad (1)$$

The symbols  $S$  and  $V$  denote matrix elements of the isoscalar and isovector interaction, respectively, which are reduced in space and isospace. From this basic expression it follows that within a  $T$  multiplet with  $T \geq 1$  the  $2T+1$  analog transition strengths are determined by only two independent parameters  $S$  and  $V$ .

Within the  $s, d$  shell fairly good agreement with these predictions has been found for the  $E2$  transitions between the two lowest  $T = T_> = 1$  states in the nuclear mass triplets  $A = 18$  (Ref. 2), 26 (Ref. 3), 30 (Refs. 4, 5) and 34 (Ref. 5).

In this note we report an extension to transitions from the second  $2^+$ ,  $T_> = 1$  multiplet in the mass  $A = 30$  system consisting of the nuclei  $^{30}\text{Si}(T_x = -1)$ ,  $^{30}\text{P}(0)$ , and  $^{30}\text{S}(+1)$  (Fig. 1). The  $E2$  transition strengths to the  $0^+$ ,  $T_>$  ground states are known to be  $|M_{E2}|^2 = 1.3 \pm 0.3$  Weisskopf units (W.u.) (Ref. 6) for  $^{30}\text{Si}$  and  $|M_{E2}|^2 = 0.6 \pm 0.2$  W.u. (Ref. 4) for  $^{30}\text{S}$ . Since the  $T_x = 0$  nucleus  $^{30}\text{P}$  has low-lying  $T_<$  as well as  $T_>$  states, the interesting  $E2$ ,  $T_> \rightarrow T_>$  analog transitions are obscured by very strong  $M1$ ,  $T_> \rightarrow T_<$  transitions (Fig. 1). The second  $2^+$ ,  $T_>$  state at  $E_x = 4.183$  MeV is known to decay with a very short lifetime ( $\tau < 15$  fs, Ref. 6) to the two lowest  $J^\pi = 1^+$ ,  $T_<$  states. No branch to the  $0^+$ ,  $T_>$  state at  $E_x = 0.677$  MeV has been reported.

A search for this branch was made at the  $E_p = 731$  keV resonance in the reaction  $^{29}\text{Si}(p, \gamma)^{30}\text{P}$  which has a 19% decay<sup>7</sup> to the  $2^+$ ,  $T_>$  state at 4.183 MeV. An enriched (95%)  $^{29}\text{SiO}_2$  target (thickness  $\approx 5$  keV) was bombarded with protons from the 3-MV Stanford Van de Graaff generator. Beam intensities

were typically between 1 and 2  $\mu\text{A}$ . The  $\gamma$  rays were detected with a Ge(Li) detector with an active volume of about 50  $\text{cm}^3$  and a resolution of 3.8 keV at 1333 keV. Spectra were taken on and off resonance with the Ge(Li) detector placed at angles of 35 and 90° at a distance of  $d = 2.7$  cm. Figure 2 shows the interesting parts of a spectrum taken at 35° for a total accumulated charge of 100 mC. A normalized spectrum taken off resonance at the same angle has been subtracted from these data. In addition to the known  $M1$ ,  $\Delta T = 1$  transitions of energies 4.183 and 3.474 MeV, a weak  $\gamma$  ray at  $E = 3.506$  MeV is observed which is identified as the missing analog  $E2$ ,  $\Delta T = 0$  transition on the basis of the following facts: (i) The  $\gamma$  ray is resonant; (ii) Its energy fits with the expected value and cannot be explained by any other possible

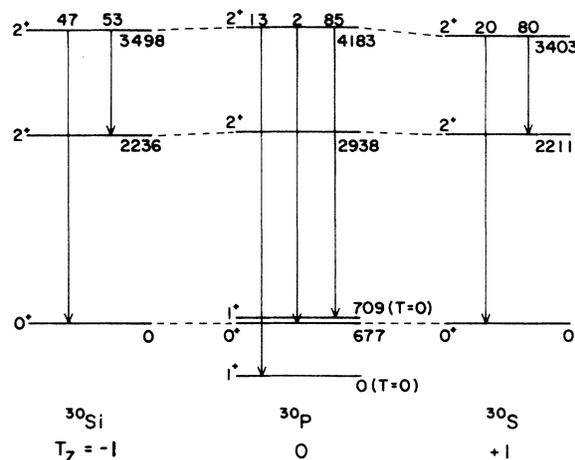


FIG. 1. Decay modes of the second  $2^+$ ,  $T=1$  states in the mass  $A = 30$  system. The branching ratios in  $^{30}\text{Si}$  and  $^{30}\text{S}$  are taken from Refs. 4 and 6, respectively, and in  $^{30}\text{P}$  from the present work. The energies are given in MeV. The isospin of all states is  $T = 1$  unless otherwise noted.

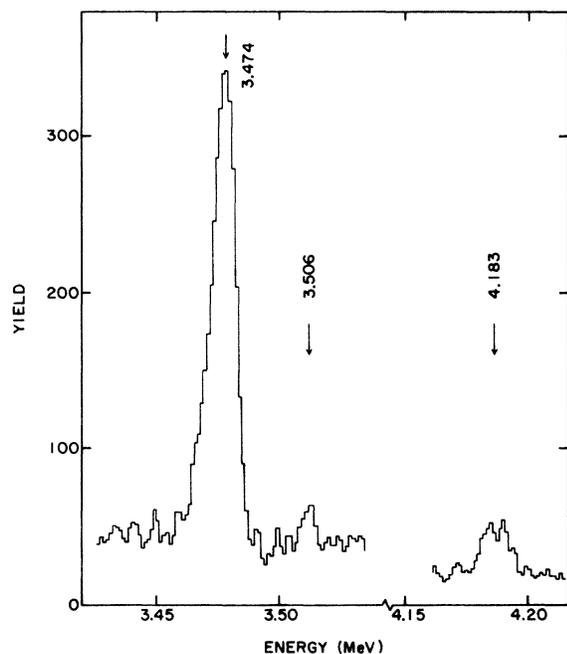


FIG. 2. Relevant parts of the  $\gamma$ -ray spectrum obtained with the Ge(Li) detector placed at  $35^\circ$ .

transition between the known states of  $^{30}\text{P}$  which could be produced in the reaction; (iii) Its observed Doppler shift of  $5 \pm 2$  keV is of the same magnitude as that of the strong 3.474 MeV transition; (iv) The population parameters of the  $E_x = 4.183$  MeV state (as determined from the angular distribution of the 3.474 MeV transition) indicate that the intensity of the  $2^+ \rightarrow 0^+$  transition should be larger at  $35^\circ$  than at  $90^\circ$ , as was indeed found in the spectra taken at these two angles. The branching ratio of this  $\gamma$  ray, corrected for angular correlation and the relative efficiency of the Ge(Li) detector, was estimated to be  $2 \pm 1\%$ . The strong branches to the ground state and the second excited state at  $E_x = 0.709$  MeV were found to be  $(13 \pm 2)$  and  $(85 \pm 2)\%$ , respectively, which are in reasonable agreement with the values quoted in the literature.<sup>6</sup>

The value observed for the new branch limits the  $E2$  strength to

$$|M_{E2}|^2 > 0.27 \text{ W.u.}$$

This value then sets a lower limit on the squared isoscalar reduced transition matrix element  $S^2$ . From the complete set of transition strengths for

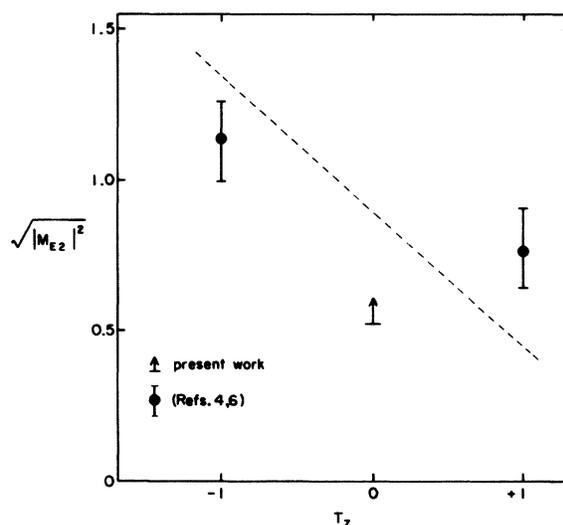


FIG. 3. The square roots of the  $E2$  transition strengths plotted as a function of  $T_z$ . The transitions are those from the second  $2^+$ ,  $T=1$  level to the lowest  $0^+$ ,  $T=1$  levels in  $^{30}\text{Si}$  ( $T_z = -1$ ),  $^{30}\text{P}(0)$ , and  $^{30}\text{S}(+1)$ . The dashed line is obtained from a shell model calculation (Ref. 8).

the  $A=30$  nuclei the following values for  $S^2$  and  $V^2$  are obtained:

$$S^2 = 0.9 \pm 0.3 \text{ W.u.}$$

$$V^2 = 0.03 \pm 0.05 \text{ W.u.}$$

The fact that the isoscalar part  $S^2$  is the dominant term in the transition matrix elements is in agreement with results for the other known sets of analog transitions in this mass region. In Fig. 3 the square roots of the  $E2$  strengths are plotted as a function of  $T_z$ ; according to Eq. (1), they should lie on a straight line. For comparison, results of a shell model calculation<sup>8</sup> are also given. The over-all agreement is quite satisfactory.

With the isoscalar and the isovector parts of the  $E2$  transition matrix elements now determined it is possible to give an approximate value for the lifetime of the second  $2^+$ ,  $T_1$  state in  $^{30}\text{P}$  which turns out to be

$$\tau = 6.0 \pm 3.5 \text{ fs.}$$

The  $M1$  strengths of transitions to the ground state and the second excited state at 0.709 MeV are then estimated to be  $0.93 \pm 0.54$  and  $10.5 \pm 6.0$  ( $10^{-2}$  W.u.), which can be compared with the shell model values<sup>8</sup> of 7.0 and 4.9 ( $10^{-2}$  W.u.), respectively.

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