

Ultrahigh-resolution neutron spectroscopy of low-energy spin dynamics in UGe₂F. Haslbeck,^{1,2} S. Säubert,^{1,3} M. Seifert,^{1,3} C. Franz,³ M. Schulz,³ A. Heinemann,⁴ T. Keller,^{5,6} Pinaki Das,^{7,*} J. D. Thompson,⁷ E. D. Bauer,⁷ C. Pfeleiderer,¹ and M. Janoschek^{2,7,8,†}¹Physik-Department, Technische Universität München, D-85748 Garching, Germany²Institute for Advanced Study, Technische Universität München, D-85748 Garching, Germany³Heinz Maier-Leibnitz Zentrum Garching, Technische Universität München, D-85748 Garching, Germany⁴German Engineering Materials Science Centre (GEMS) at Heinz Maier-Leibnitz Zentrum (MLZ), Helmholtz-Zentrum Geesthacht GmbH, D-85747 Garching, Germany⁵Max-Planck-Institut für Festkörperforschung, D-70569 Stuttgart, Germany⁶Max Planck Society Outstation at the Forschungsneutronenquelle Heinz Maier-Leibnitz (MLZ), D-85747 Garching, Germany⁷Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA⁸Laboratory for Scientific Developments and Novel Materials, Paul Scherrer Institut, Villigen PSI, Switzerland

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Studying the prototypical ferromagnetic superconductor UGe₂ we demonstrate the potential of the modulated intensity by zero effort (MIEZE) technique—a novel neutron spectroscopy method with ultrahigh energy resolution of at least 1 μeV —for the study of quantum matter. We reveal purely longitudinal spin fluctuations in UGe₂ with a dual nature arising from 5 f electrons that are hybridized with the conduction electrons. Local spin fluctuations are perfectly described by the Ising universality class in three dimensions, whereas itinerant spin fluctuations occur over length scales comparable to the superconducting coherence length, showing that MIEZE is able to spectroscopically disentangle the complex low-energy behavior characteristic of quantum materials.

DOI: [10.1103/PhysRevB.99.014429](https://doi.org/10.1103/PhysRevB.99.014429)**I. INTRODUCTION**

Ultralow spin dynamics represent a key characteristic of quantum matter such as quantum spin liquids [1], electronic nematic phases [2], topological spin textures [3,4], non-Fermi liquid behavior [5–8], and unconventional superconductivity [9]. For the clarification of these phenomena spectroscopic methods with excellent momentum and energy resolution are required, as key characteristics emerge typically in the low milli-Kelvin range. In principle, neutron scattering is ideally suited for studies of the relevant spin excitations. However, the typical energy resolution of conventional neutron spectroscopy corresponds to Kelvin temperatures. Although techniques such as neutron spin-echo spectroscopy offer ultrahigh resolution of sub- μeV , they are incompatible with conditions that depolarize neutron beams such as ferromagnetism (FM), superconductivity, or large magnetic fields.

The discovery of superconductivity in the FM state of UGe₂ highlights the combination of scientific and experimental challenges that arise in the study of the complex low-energy behavior of quantum matter that characteristically emerges due to the competition of high-energy atomic energy scales [10]. Namely, actinide-based compounds such as UGe₂ are formidable model systems, where the hybridization of itinerant d and localized f electrons drives low-energy

excitations that mediate a multitude of novel states [11–13]. Traditionally the concomitant subtle reconstruction of the electronic structure has been studied via the charge channel, which fails to provide the required high resolution [14,15]. Exploiting in contrast the spin channel recent advances in neutron spectroscopy provided new insights [16–18].

Using an implementation of neutron resonance spin echo (NRSE) spectroscopy that is insensitive to depolarizing conditions, namely the so-called modulated intensity by zero effort (MIEZE) [19], we report a study of UGe₂ in which we identify the enigmatic low-energy excitations as an unusual combination of fluctuations attributed normally either to itinerant or localized electrons in an energy and momentum range comparable to the superconducting coherence length and ordering temperature. As the superconductivity in UGe₂ represents a prototypical form of quantum matter, our study underscores also the great potential of the MIEZE technique in studies of quantum matter on a more general note.

At ambient pressure UGe₂ displays ferromagnetism with a large Curie temperature, $T_c = 53$ K, and a large ordered moment, $\mu_{\text{FM1}} = 1.2 \mu_B$ (FM1) [20,21]. Under increasing pressure FM order destabilizes, accompanied by the emergence of a second FM phase below $T_x < T_c$ where $\mu_{\text{FM2}} = 1.5 \mu_B$ (FM2). The FM2 and FM1 phases vanish discontinuously at $p_x \approx 12.2$ kbar and $p_c \approx 15.8$ kbar, respectively [21], while superconductivity emerges between ≈ 9 kbar $< p_x$ and p_c . Evidence for a microscopic coexistence of FM order and superconductivity makes UGe₂ a candidate for p-wave pairing, where the Cooper pairs form spin triplets [20]. This p-wave superconductivity is believed to be mediated by an abundance of low-lying longitudinal spin fluctuations associated

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with a FM quantum phase transition (QPT), where transverse spin fluctuations are theoretically known to break spin-triplet pairing [22]. Neutron triple-axis spectroscopy (TAS) at ambient pressure indeed identified predominantly longitudinal spin fluctuations in UGe₂ [23] but failed to provide insights into the character of the fluctuations in the momentum and energy range comparable to the superconducting coherence length and transition temperature, respectively.

Prior to our study the interplay of seemingly conflicting ingredients of the spectrum of spin fluctuations were unresolved. On one hand, the strong Ising anisotropy promotes longitudinal spin fluctuations as typically attributed to localized electrons in the presence of strong to spin-orbit coupling. This is contrasted, on the other hand, by the notion of Cooper pairs and a well developed, strongly exchange-split Fermi surface [22,24]. Consistent with this dichotomy characteristic of p-wave superconductivity, our ultrahigh resolution data reveals that the low-energy spin fluctuations of UGe₂ reflect a subtle interplay of itinerant and local electronic degrees of freedom on scales comparable to the superconductivity.

II. EXPERIMENTAL METHODS

NRSE achieves extreme energy resolution by encoding the energy transfer $\hbar\omega$ of the neutrons in their polarization as opposed to a change of wavelength. However, FM domains, Meissner flux expulsion, or applied magnetic fields typically depolarize the beam. We used therefore a novel NRSE technique, so-called MIEZE, implemented at the instrument RESEDA at the Heinz Maier-Leibnitz Zentrum (MLZ) [25–27]. Generating an intensity modulated beam by means of resonant spin flippers and a spin analyzer in front of the sample the amplitude of the intensity modulation assumes the role of the NRSE polarization. Because all spin manipulations are performed before the sample, beam depolarizing effects are no longer important. Using incident neutrons with a wavelength $\lambda = 6 \text{ \AA}$ and $\Delta\lambda/\lambda \approx 10\%$ provided by a velocity selector, we achieved an energy resolution of $\Delta E \approx 1 \mu\text{eV}$. MIEZE in small angle neutron scattering (SANS) configuration also provides high momentum \mathbf{q} resolution of approximately 0.015 \AA^{-1} . The MIEZE setup is described in the Supplemental Material [28].

A high-quality single crystal of UGe₂ was grown by the Czochralski technique followed by an annealing similar to Ref. [29]. A cylindrical piece with nearly constant diameter of 7 mm and 16 mm length ($m = 6 \text{ g}$) with the crystallographic c axis approximately parallel to the cylinder axis was cut for the MIEZE experiments. The sample was oriented using neutron Laue diffraction so that c was perpendicular to the scattering plane. The Laue images also confirm a high-quality single-grain sample [28]. Neutron depolarization imaging measurements [28] of the same sample reveal that the magnetic properties of the crystal are completely homogeneous with a Curie temperature $T_C = 52.68(3) \text{ K}$ demonstrating that this sample is optimal for the investigation of critical spin fluctuations. Magnetic susceptibility measurements were performed on a small piece ($m = 36 \text{ mg}$) of the same sample in a Quantum Design magnetic property measurement system (MPMS).

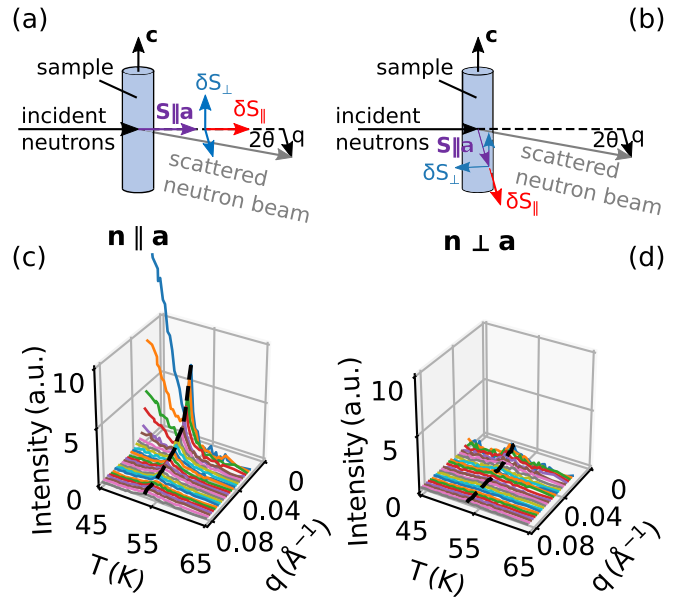


FIG. 1. Magnetic intensity in UGe₂ near the Curie temperature $T_C = 52.7 \text{ K}$. (a) and (b) show two experimental configurations with the a axis parallel or perpendicular to the incident neutron beam \mathbf{n} , respectively, used to differentiate longitudinal from transverse spin fluctuations (see text). (c) and (d) show the observed energy-integrated intensities for $\mathbf{n} \parallel \mathbf{a}$ and $\mathbf{n} \perp \mathbf{a}$ as a function of temperature T and momentum transfer \mathbf{q} . The black dashed line marks T_C .

III. RESULTS AND DISCUSSION

The magnetic cross section is related to the imaginary part of the dynamical magnetic susceptibility $\chi''_{ij}(\mathbf{q}, \omega)$ via

$$\frac{d^2\sigma}{d\Omega d\omega} \propto \frac{k_f}{k_0} (\delta_{ij} - \hat{q}_i \hat{q}_j) |F_{\mathbf{q}}|^2 [n(\omega) + 1] \chi''_{ij}(\mathbf{q}, \omega), \quad (1)$$

where k_0 and k_f are the wave vector of the incident and scattered neutrons, respectively. \hat{q} is a unit vector parallel to the scattering vector \mathbf{q} and $n(\omega)$ is the Bose function. $F_{\mathbf{q}}$ is the uranium magnetic form factor.

In Fig. 1 we show the temperature and \mathbf{q} dependence of the energy-integrated intensity of the spin fluctuations in UGe₂ that was obtained by switching the MIEZE setup off. Nonmagnetic background scattering obtained well above T_C was subtracted from all data sets shown. The temperature scan was carried out with the crystallographic a axis, which is the magnetic easy axis for UGe₂ oriented parallel ($\mathbf{n} \parallel \mathbf{a}$) and perpendicular ($\mathbf{n} \perp \mathbf{a}$) to the incident neutron beam, respectively. Due to the term $\delta_{ij} - \hat{q}_i \hat{q}_j$ in Eq. (1) neutron scattering is only sensitive to spin fluctuations that are perpendicular to \mathbf{q} . Because in SANS configuration \mathbf{q} is approximately perpendicular to the incident neutron beam, this allows us to separate longitudinal (δS_{\parallel}) from transverse spin fluctuations (δS_{\perp}) as illustrated in Figs. 1(a) and 1(b). For $\mathbf{n} \parallel \mathbf{a}$ both δS_{\parallel} and δS_{\perp} are perpendicular to \mathbf{q} . As shown in Fig. 1(c) substantial magnetic intensity is observed for this configuration. In contrast, for $\mathbf{n} \perp \mathbf{a}$ only δS_{\perp} is perpendicular to \mathbf{q} and the vanishingly small signal observed in this case [see Fig. 1(d)] can only come from transverse spin fluctuations. Because of the cylindrical shape of the sample differences in

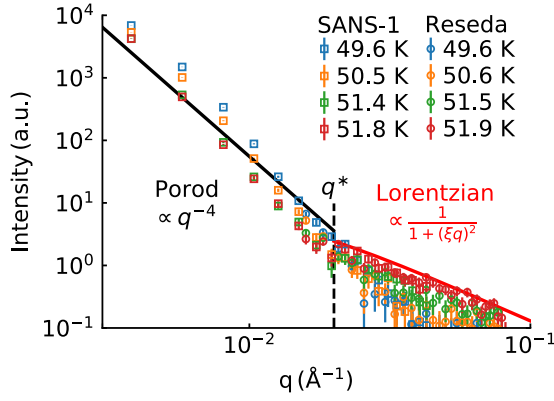


FIG. 2. \mathbf{q} dependence of the intensity for selected T below T_C for $\mathbf{n} \parallel a$. Below $q^* \approx 0.02 \text{ \AA}^{-1}$ the intensity is well described by Porod scattering due to ferromagnetic (FM) domains (black solid line), whereas above q^* a Lorentzian shape due to critical spin fluctuations is observed (red solid line).

neutron transmission between the two orientations are negligible. As shown in the Supplemental Material [28], the small intensity observed for $\mathbf{n} \perp a$ arises from finite \mathbf{q} resolution, demonstrating that the critical spin fluctuations in UGe_2 are solely longitudinal.

Inspecting the temperature dependence of the integrated intensity for $\mathbf{n} \parallel a$ [see Fig. 1(c)], a pronounced peak is centered at $T_C = 52.7 \text{ K}$ due to the divergence of critical spin fluctuations. For low q and for $T < T_C$ additional intensity is observed that increases like a magnetic order parameter. Figure 2 shows the q dependence of the intensity for a few temperatures below T_C . Below $q^* \approx 0.02 \text{ \AA}^{-1}$ the intensity is well described by a q^{-4} dependence that is characteristic for scattering from FM domains that form below T_C [30,31]. To follow this so-called Porod scattering towards lower q , we have performed a supporting SANS experiment on the instrument SANS-1 at MLZ (details are described in Ref. [28]) denoted with square symbols in Fig. 2. Observation of Porod scattering down to $q_{\min} = 0.004 \text{ \AA}^{-1}$ implies the onset of long-range order over length scales $\gg 2\pi/q_{\min} \approx 1600 \text{ \AA}$. In Fig. 3, we show the temperature dependence of the intensity for selected q below q^* . Near to T_C it evolves as $M^2(T) \propto (1 - T/T_C)^{2\beta}$. We find that $\beta = 0.32(1)$ describes our data perfectly in agreement with $\beta^{\text{theo}} = 0.32$ for a three-dimensional (3D) Ising system [32]. This is also in good agreement with $\beta = 0.36(1)$ from neutron diffraction [33].

For $q \geq q^*$ and for $T \approx T_C$ the q dependence is described by a Lorentzian line shape characteristic of critical spin fluctuations with a correlation length ξ . The corresponding dynamical magnetic susceptibility is

$$\frac{\chi''(\mathbf{q}, \omega)}{\omega} = \chi(\mathbf{q}) \frac{\Gamma_q}{\Gamma_q^2 + \omega^2} \quad (2)$$

$$\chi(\mathbf{q}) = \frac{\chi_0}{1 + (\xi q)^2}, \quad (3)$$

where Γ_q and χ_0 are the momentum dependent relaxation frequency and the static magnetic susceptibility, respectively. Because of the longitudinal character of the spin fluctuations only χ''_{aa} is nonzero, and we have thus dropped the indices i, j .

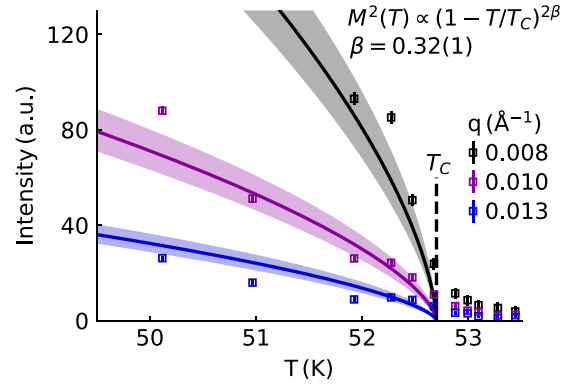


FIG. 3. The T dependence of the Porod scattering for $q < q^*$ follows the FM order parameter M via $M^2(T) \propto (1 - T/T_C)^{2\beta}$ with $\beta = 0.32(1)$ (solid lines). The shaded regions denote the uncertainty of the fit of β .

To investigate the critical scattering quantitatively, we subtract the Porod scattering (Fig. 2) from the observed intensities [Fig. 1(c)]. For our experimental conditions the quasistatic approximation [34,35] is valid and thus integrating Eq. (1) with respect to $\hbar\omega$, we obtain $\frac{d\sigma}{d\Omega} \propto T\chi(\mathbf{q})$ (see Supplemental Material [28]). We show $\chi(\mathbf{q})$ obtained by dividing the observed intensity by T for various temperatures in Fig. 4(b). The solid lines are fits to Eq. (3) to extract the T dependence of χ_0 and ξ shown in Figs. 5(a) and 5(b), respectively. For comparison we show the static magnetic susceptibility H/M determined by bulk magnetization measurements in Fig. 5(a) (blue squares) that scales perfectly with χ_0 .

We find that $1/\chi_0 \propto (1 - T/T_C)^\gamma$ with $\gamma = 1.23(3)$ and $1/\xi = \kappa = \kappa_0(1 - T/T_C)^\nu$ with $\kappa_0 = 0.31(2) \text{ \AA}^{-1}$ and $\nu = 0.63(2)$ [solid lines Figs. 5(a) and 5(b)]. The critical exponents are in excellent agreement with a 3D Ising FM, for which $\gamma^{\text{theo}} = 1.24$ and $\nu^{\text{theo}} = 0.63$ [32]. Huxley *et al.* found $\kappa_0 = 0.29 \text{ \AA}^{-1}$ in good agreement with our result. In contrast, they determined $\nu = 1/2$, consistent with a mean-field transition [23]. However, their study was limited to $q > 0.03 \text{ \AA}^{-1}$ and underestimates the divergence of the critical fluctuations.

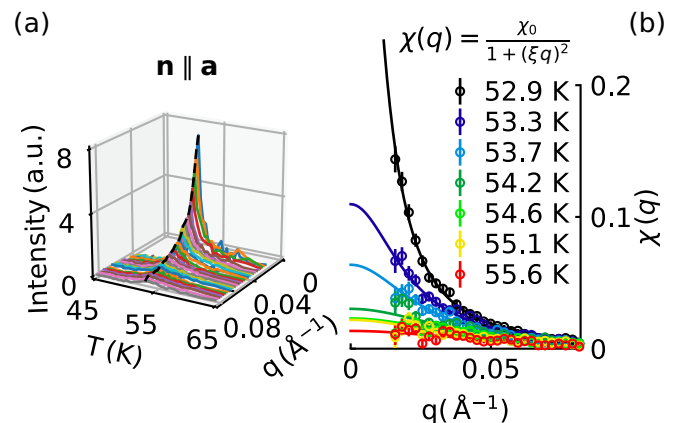


FIG. 4. (a) Temperature T and momentum transfer \mathbf{q} dependence of the critical Ising spin fluctuations in UGe_2 . (b) q dependence of the magnetic susceptibility $\chi(q)$. Solid lines are fits to Eq. (3).

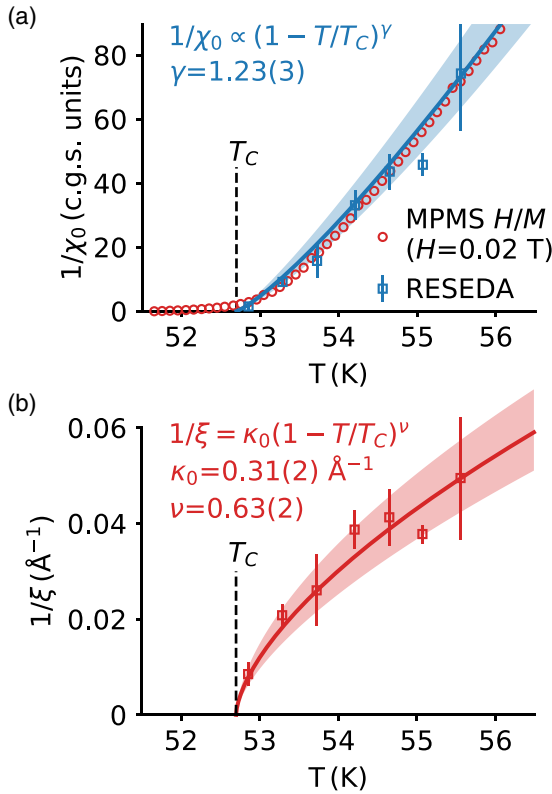


FIG. 5. The inverse susceptibility $1/\chi_0$ and inverse correlation length $1/\xi$ as a function of temperature T (near the Curie temperature $T_C = 52.7$ K), respectively, resulting from fits in Fig. 4. The blue squares in (a) denote the static easy-axis magnetic susceptibility H/M determined with a magnetic field $H = 0.1$ T. The solid black and red lines are fits to determine the critical exponents for χ_0 and ξ (see text), and the shaded region denotes the uncertainty of the fit.

We now discuss the results of our MIEZE measurements. MIEZE measures the intermediate scattering function $S(q, \tau)$ that is the time Fourier transform of the scattering function $S(q, \omega) = 1/\pi[n(\omega) + 1]\chi''_{ij}(\mathbf{q}, \omega)$ [cf. Eq. (1)]. In Fig. 6 we show $S(q, \tau)$ for various q at T_C . $S(q, \tau)$ for all other

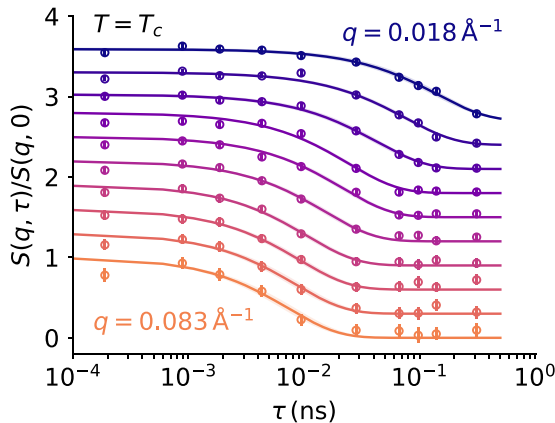


FIG. 6. Spin fluctuation spectrum of UGe_2 obtained by MIEZE. The intermediate scattering function $S(q, \tau)$ normalized to $S(q, 0)$ (static signal) is shown at $T_C = 52.7$ K for a range of momentum transfers q . Solid lines are fits to Eq. (4).

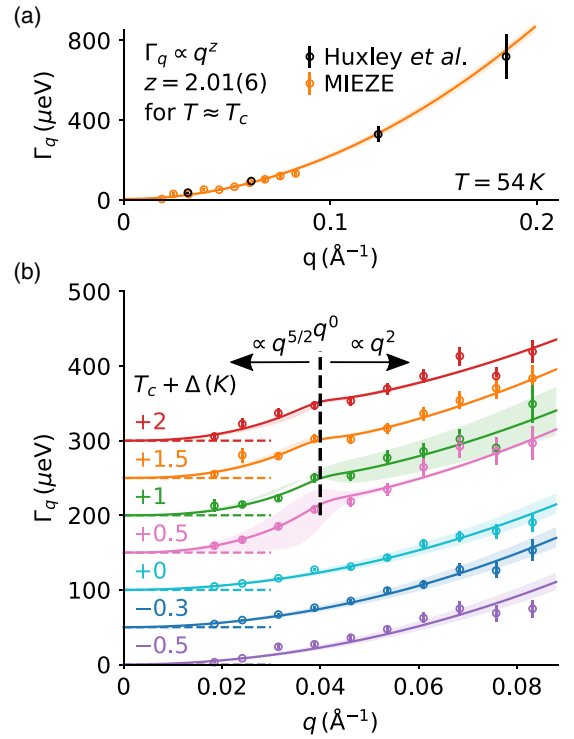


FIG. 7. The fluctuation frequency Γ_q of UGe_2 at various temperatures T as determined by fits of 4 to the data shown in Fig. 6. Solid lines are fits to $\Gamma_q \propto q^z$, where z is the dynamical critical exponent. (a) Comparison of our data to the high- q data by Huxley *et al.* [23] is shown. (b) We find two distinct regimes with $z = 2.5$ and 2 below and above $q^0 = 0.038 \text{ \AA}^{-1}$, respectively (see text). Data sets are shifted by $50 \mu\text{eV}$ for better readability as indicated by the horizontal dashed lines.

measured temperatures are shown in Ref. [28]. Because the spin fluctuations have Lorentzian lineshape [see Eq. (2)] we fit $S(q, \tau)$ with an exponential decay [solid lines in Fig. 6]:

$$S(q, \tau) = \exp(-|\Gamma_q| \cdot \tau). \quad (4)$$

The resulting fluctuation frequency Γ_q is shown in Fig. 7.

The momentum dependence of Γ_q is described by the dynamical exponent z via $\Gamma_q \propto q^z$. For $T \leq T_C$, we find that Γ_q is fitted perfectly by $z = 2.0(1)$ [Fig. 7(a)]. This is in excellent agreement with predictions for a 3D Ising FM, for which $z^{\text{theo}} = 2$ [32]. For $T > T_C$, Γ_q is also well described by $z = 2$, however, only above a crossover value of $q^0 = 0.04 \text{ \AA}^{-1}$. Below q^0 , our data is best fit by $\Gamma_q = Aq^z$ with $z = 2.53(4)$ (see Fig. 7). This is consistent with $z = 5/2$ calculated for itinerant FMs within critical renormalization group theory [36] and confirmed for various d -electron FMs such as Fe [37], Ni [38], and Co [39]. Notably, typical values reported for A are $3\text{--}350 \text{ meV \AA}^{5/2}$ [37–40] consistent with $A = 200(2) \text{ meV \AA}^{5/2}$ that we find for UGe_2 . As demonstrated in Fig. 7(a) for $T = 54$ K, the fit of Γ_q with $z = 2.0(1)$ also describes the data of Huxley *et al.* [23] (black empty circles) perfectly. However, they conclude that Γ_q remains finite for $q \rightarrow 0$ in contrast to our findings. This discrepancy is easily explained by considering that their experiment was

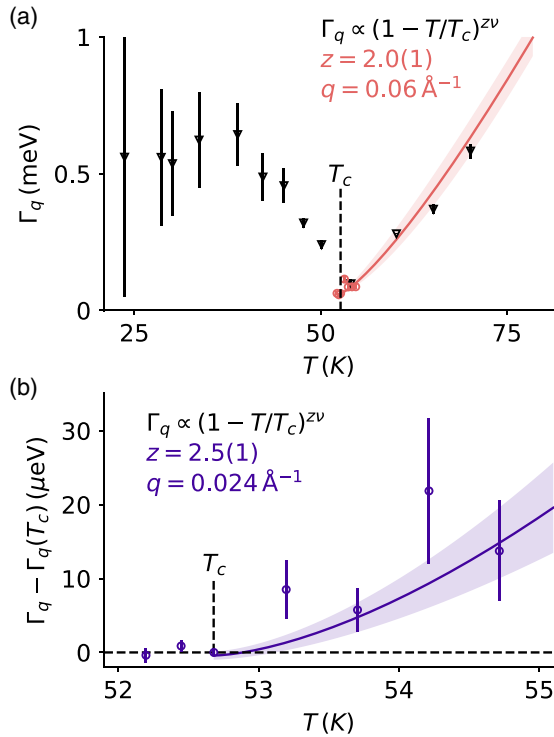


FIG. 8. Temperature (T) dependence of the fluctuation frequency Γ_q for q above (a) and below (b) q^0 . (a) Comparison to data of Ref. [23]. Solid lines denote $\Gamma_q \propto (1 - T/T_c)^{z\nu}$ (see text).

limited to $q \geq 0.03 \text{ \AA}^{-1}$, which is only slightly below q^0 where we observe the crossover to $z = 5/2$.

Figure 8 shows the T dependence of Γ_q . For finite q , it follows the T dependence of ξ via $\Gamma_q \propto (1/\xi)^z = (1 - T/T_c)^{z\nu}$ in agreement with the dynamical scaling prediction [41]. In Fig. 8(a) we show that for $q = 0.06 \text{ \AA}^{-1}$ both the results from Ref. [23] and our own are consistent with $z = 2$. Below q^0 , $z = 5/2$ agrees well with our data (solid line) consistent with the fits of Γ_q shown in Fig. 7.

For clean itinerant FMs the fluctuation spectrum is characterized by Landau damping as has been demonstrated for 3d transition metal materials [42,43]. Here the product of the magnetic susceptibility with the fluctuation frequency, $\chi(q)\Gamma_q$, is given by the Lindhard dependence $(2/\pi)v_F\chi_Pq$ for $T > T_c$, where v_F and χ_P are the Fermi velocity and the noninteracting Pauli susceptibility, respectively [44,45]. We show $\chi(q)\Gamma_q$ for UGe_2 in Fig. 9. Huxley *et al.* [23] who carried out measurements for $q \geq 0.03 \text{ \AA}^{-1}$ found that $\chi(q)\Gamma_q$ only weakly depends on q and concluded that it remains finite for $q \rightarrow 0$ (solid black line in Fig. 9). This difference with respect to prototypical 3d electron itinerant FMs is likely due to strong spin-orbit coupling that modifies the spin fluctuation spectrum. Our data agrees with the weak q dependence above q^0 but clearly shows that $\chi(q)\Gamma_q \rightarrow 0$ for $q \rightarrow 0$, implying that the uniform magnetization is a conserved quantity in UGe_2 . Our data is consistent with $\chi(q)\Gamma_q \propto q^{5/2}$ (solid blue line in Fig. 9). This more pronounced q dependence is expected by theory near T_c [45] and agrees with $\Gamma_q \propto q^{5/2}$. Here, we highlight that although the q

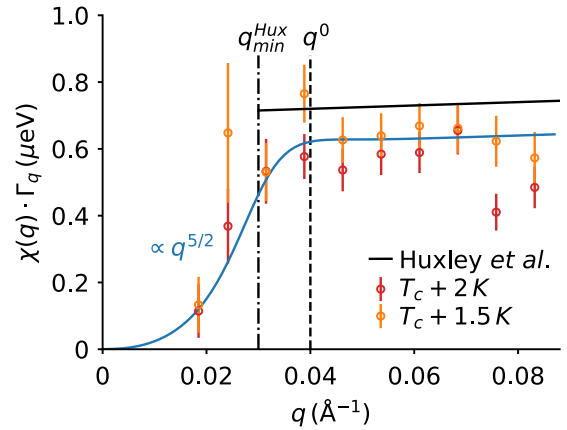


FIG. 9. Temperature (T) dependence of the product of the magnetic susceptibility with the fluctuation frequency, $\chi(q)\Gamma_q$. The black line is $\chi(q)\Gamma_q$ as determined in Ref. [23] that reports measurements down to $q_{\min}^{\text{Hux}} = 0.03 \text{ \AA}^{-1}$ denoted by the dashed-dotted line. The blue line is a guide to the eye.

range over which q^z with $z = 5/2$ is observed is limited, this behavior is corroborated via three independent methods that are illustrated in Figs. 7–9.

IV. SUMMARY

Our results demonstrate that the spin fluctuations in UGe_2 exhibit a dual character associated with localized 5f electrons that are hybridized with itinerant d electrons. Notably, as expected for a local moment FM with substantial uniaxial magnetic anisotropy all critical exponents determined from our results are in perfect agreement with the 3D Ising universality class [32]. Further, $\chi(q)\Gamma_q$ is approximately constant as a function of q down to q^0 highlighting that the underlying spin fluctuations are localized in real space. In contrast, the dynamical exponent $z = 5/2$ and $\chi(q)\Gamma_q \rightarrow 0$ for $q \rightarrow 0$ observed below the crossover value q^0 are characteristic of itinerant spin fluctuations. Because the contribution of the conduction electrons to the total ordered moment is less than 3% [33], below T_c fluctuations of localized f magnetic moments are dominant. Spin fluctuations with a dual character are consistent with the moderately enhanced Sommerfeld coefficient $\gamma = 34 \text{ mJ/K}^2 \text{ mol}$ of UGe_2 [46,47] and a next-nearest-neighbor uranium distance $d_{\text{U-U}} = 3.85 \text{ \AA}$ [48] near to the Hill value of 3.5 \AA [49] that both suggest that the 5f electrons in UGe_2 are hybridized with the conduction electrons.

In conclusion, the dual nature of spin fluctuations revealed by our MIEZE measurements strongly supports the scenario of p-wave superconductivity in UGe_2 . First, to promote strong longitudinal fluctuations requires strong Ising anisotropy that typically is a result of localized f electrons with substantial spin-orbit coupling and is consistent with critical Ising exponents that we observe above q^0 . Second, the theory for p-wave pairing assumes that it is the same itinerant electrons that are responsible for the coexisting FM and superconducting states [22], highlighting that the low-energy itinerant spin fluctuations below q^0 discovered here are crucial to mediate p-wave superconductivity. The maximum superconducting

critical temperature T_s occurs at the QPT at p_x [20,21]. Here a substantial increase of the Sommerfeld coefficient [50] and changes in the electronic structure observed near p_x [51,52] suggest that the hybridization of $5f$ electrons and conduction electrons increases at p_x and corroborates that spin fluctuations with a dual nature are relevant for p-wave superconductivity. This is supported by a theory based on competition of FM exchange and the Kondo interaction that results in a localized to itinerant transition at p_x [53,54].

Further, we note that our findings of longitudinal critical fluctuations in UGe₂ are also consistent with the findings for UCoGe [55], which is another material that is a candidate for p-wave superconductivity. However, the results on UCoGe by Hattori *et al.* [55] were obtained by NMR measurements that are unable to probe spin fluctuations at finite q and, in turn, are unable to observe an itinerant-to-localized crossover as we report it here. Similarly, TAS measurements of UCoGe by Stock *et al.* [56] lack the required momentum and energy transfer resolution.

Finally, we note that the crossover value q^0 corresponds to a length scale of approximately 160 Å. The superconducting coherence length of UGe₂ was estimated as $\xi^{SC} = 200$ Å [20], which shows that the spin fluctuations relevant to the p-wave pairing are present at $q < q^0$. This may explain why triple-axis measurements of the spin fluctuation near p_x with limited resolution were inconclusive [57]. Although, the pressure dependence of the crossover length scale q^0 remains to be determined to unambiguously associate it with the unconventional superconducting state in UGe₂, our results highlight that recent developments in ultrahigh resolution neutron spectroscopy are critical for the study of low-energy spin fluctuations that are believed to drive the emergence of quantum matter states. Here the fluctuations that

appear at zero q such as for ferromagnetic and electronic-nematic quantum states can immediately be investigated via the MIEZE SANS configuration used here. In addition, MIEZE can be extended in straightforward fashion to study quantum fluctuations arising at large q [58], allowing for insights in antiferromagnetic QPTs and topological forms of order.

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