

Role of band filling in tuning the high-field phases of URu₂Si₂

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We present a detailed study of the low-temperature and high-magnetic-field phases in the chemical substitution series URu₂Si_{2-x}P_x using electrical transport and magnetization in pulsed magnetic fields up to 65 T. Within the hidden-order x regime ($0 < x \lesssim 0.035$) the field-induced ordering that was earlier seen for $x = 0$ is robust, even as the hidden-order temperature is suppressed. Earlier work shows that for $0.035 \lesssim x \lesssim 0.26$ there is a Kondo lattice with a no-ordered state that is replaced by antiferromagnetism for $0.26 \lesssim x \lesssim 0.5$. We observe a simplified continuation of the field-induced order in the no-order x regime and an enhancement of the field-induced order upon the destruction of the antiferromagnetism with magnetic field. These results closely resemble what is seen for URu_{2-x}Rh_xSi₂,¹ from which we infer that charge tuning dominantly controls the ground state of URu₂Si₂, regardless of whether s/p or d electrons are replaced. Contraction of the unit-cell volume may also play a role at large x . This provides guidance for determining the specific factors that lead to hidden order versus magnetism in this family of materials and constrains possible models for hidden order.

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I. INTRODUCTION

Among the f -electron intermetallics, URu₂Si₂ continues to attract interest because it hosts an unidentified ordered state (“hidden order”) and unconventional superconductivity at temperatures below $T_0 = 17.6$ K and $T_c = 1.5$ K, respectively [1–5]. These phenomena occur within a strongly hybridized f -electron lattice that is superficially similar to that of related systems with magnetically ordered ground states [6,7]. Despite this parallel, various measurements (e.g., neutron scattering) have revealed that the ordered state does not have an intrinsic magnetic moment [8]. A multitude of theories have been proposed to describe hidden order, where a distinguishing factor is the assumed degree of f -electron localization, but no consensus has been reached regarding their applicability [4,5].

To solve this puzzle, it is important to understand what factors distinguish between the generic occurrence of magnetism in other related f -electron lattices and the singular behavior of URu₂Si₂. To some extent, the continuity of experimental information extracted from applied pressure (P), chemical substitution (x), and magnetic field ($\mu_0 H$) tuning series has been useful to address this question. For example, pressure drives a first-order phase transition from hidden order into antiferromagnetism near $P_c = 5$ kbar [9]. Chemical substitution also tends to promote magnetism, where Ru \rightarrow Fe and Os yield phase diagrams similar to those seen with pressure [10–12], Ru \rightarrow Tc and Re stabilize ferromagnetism, and Ru \rightarrow Rh and Ir eventually produce antiferromagnetism [13–15]. Particularly interesting is that large magnetic fields suppress hidden order and uncover a rich family of magnetically ordered field-induced (FI) states, where elastic neutron scattering in pulsed magnetic fields recently revealed that the lowest-in-magnetic field of them is a type of spin-density wave order [16–20].

These tuning strategies reveal rich phenomena and indicate a close relationship between hidden order and magnetism, but what is missing is both a picture that unifies the diverse behavior and simple tuning schemes to access the multitude of ordered states in clean single crystals at ambient pressure. In this context, ligand-site substitution in URu₂Si₂ is an obvious target for investigation. Thus motivated, we recently examined the chemical substitution series URu₂Si_{2-x}P_x, where the Kondo lattice behavior is preserved but the hidden order (HO) is replaced by a no-ordering (NO) heavy Fermi liquid for $0.035 \lesssim x \lesssim 0.26$ that eventually gives way to antiferromagnetism (AFM) for $x \gtrsim 0.26$ [21–23]. This phase diagram opens the opportunity to directly examine the effect of electronic shell filling, which at low x merely tunes the density of states at the Fermi energy without disturbing the underlying band structure.

Here we report magnetoresistance and magnetization measurements in pulsed magnetic fields up to $\mu_0 H = 65$ T spanning the entire T - x phase diagram of URu₂Si_{2-x}P_x. For concentrations in the HO x regime ($x \lesssim 0.035$), the critical magnetic fields of the FI phases slightly increase, even as T_0 decreases. In the NO x regime ($0.035 \lesssim x \lesssim 0.26$) a magnetic-field-induced ordered state appears for 28 T $< \mu_0 H < 43$ T which connects continuously to the low- x /large- $\mu_0 H$ ordering. Within the AFM x regime ($x \gtrsim 0.26$) magnetic fields suppress the magnetic ordering temperature towards zero, and for $\mu_0 H > 43$ T an enhanced FI state appears which connects to the FI phase seen in the NO x regime. This may suggest either that the underlying Fermi surface evolves to become more favorable to magnetism at large x or magnetic fluctuations are helpful to the high-field ordering.

II. EXPERIMENTAL

Single-crystal specimens were produced using the molten metal flux growth technique described in Refs. [21] and [22]. Samples were prepared for electrical resistance measurements by spark-welding platinum wires to their surface and then gluing them to quartz substrates, after which data was collected

¹The concentration in this paper is defined as URu_{2-x}Rh_xSi₂ while the chemical formula in the literature is given as U(Ru_{1-x}Rh_x)₂Si₂ [28–30].

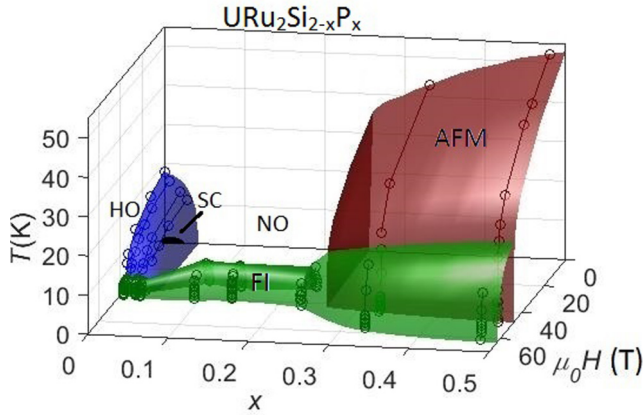


FIG. 1. Three-dimensional phase diagram for $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ single crystals constructed from magnetoresistance measurements, with temperature T , magnetic field μ_0H , and phosphorous concentration x as the three axes. μ_0H is applied parallel to the crystallographic c axis. Data for x and $H = 0$ are from Refs. [21] and [22]. Circles are our experimental data, and lines/colored regions are guides to the eye. Regions are labeled as follows: SC = superconductivity, HO = hidden x order regime, NO = no-ordered x regime, FI = field-induced order, AFM = antiferromagnetism.

using a four-point ac lock-in method with magnetic field applied parallel to the c axis. Magnetization measurements were performed using an extraction magnetometer where mosaics of 10–20 crystals ($m \approx 1$ –1.6 mg) were placed in Apiezon N grease in a cylindrical plastic capsule such that their c axis would be aligned parallel to the applied magnetic field. Due to difficulty in loading the crystals in this configuration, they were somewhat misaligned with respect to the capsule axis and magnetic field. Measurements were made with the sample/capsule both in and out of the coil, after which the two data sets were subtracted from each other to isolate the sample/capsule signal from that of the detection coil. Electrical resistance and magnetization data were collected at temperatures $0.5 \text{ K} < T < 20 \text{ K}$ and $0.6 \text{ K} < T < 20 \text{ K}$, respectively, in pulsed magnetic fields up to 65 T with pulse widths of 65 ms at the National High Magnetic Field Laboratory located at Los Alamos National Laboratory.

III. RESULTS

Figure 1 shows the three-dimensional phase diagram for single crystals of $\text{URu}_2\text{Si}_{2-x}\text{P}_x$ constructed from magnetoresistance measurements, with the axes of temperature T , phosphorous concentration x , and magnetic field μ_0H . Data for $\mu_0H = 0$ are taken from Refs. [21] and [22], where the parent compound hidden order and superconductivity are rapidly suppressed for $x \lesssim 0.035$ and are replaced by a region with persistent Kondo-lattice behavior but no low-temperature ordered state (NO x regime). Over this x range the lattice compression and strain is small and the evolution of T_0 and T_c is attributed to s -/ p -shell band filling. The NO x regime persists for $0.035 \lesssim x \lesssim 0.26$, after which antiferromagnetism emerges from the f -electron lattice for $0.26 \lesssim x \lesssim 0.5$. Starting in the middle of the NO x regime, the chemical pressure P_{ch} exceeds that needed to induce antiferromagnetism

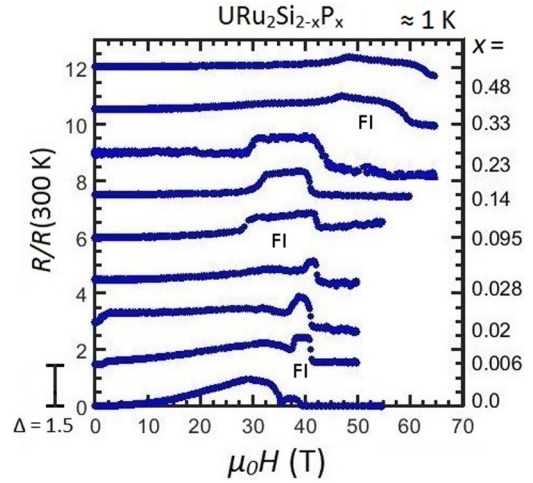


FIG. 2. Waterfall plot of normalized electrical resistance $R/R(300 \text{ K})$ vs magnetic field μ_0H for various concentrations of x at $T \approx 1 \text{ K}$. Each trace is offset by an amount $\Delta = 1.5$. The field-induced FI phase originates as a narrow region in-field in the HO x regime ($0 < x < 0.035$) and expands to a much broader range in the no-ordered x regime ($0.035 < x < 0.26$). The FI order may persist at high fields in the AFM x regime ($0.26 < x < 0.5$).

in the parent compound, and we infer that the influence of lattice compression becomes important over this range. In principle, chemical disorder might also play an important role. However, hidden order and superconductivity were previously shown to be robust even against strong disorder [22], e.g., both ordered states persist even in specimens with residual resistivity ratios $\text{RRR} = \rho_{300\text{K}}/\rho_0 \approx 10$ [24]. As shown in Fig. 4 of the Supplemental Material, the residual resistivity ratio is greater than 10 ($\text{RRR} \geq 10$) in the x regime where hidden order is destroyed and is replaced by the no-order ground state [25]. From this we infer that disorder is unlikely to be an important type of tuning. Recent NMR measurements further elucidate the behavior in these regions, where the Kondo-lattice behavior of the NO x regime is similar to that seen above T_{HO} and the antiferromagnetism occurs in the bulk and has a commensurate wave vector [23].

The response of the parent compound to an applied magnetic field is also well known [16–20]. For $T < T_0$, the magnetoresistance initially increases with μ_0H and eventually drops to a minimum near 35 T, indicating the end of the hidden-order phase (Fig. 2). Within this magnetic field range, Shubnikov–de Haas (SdH) oscillations reveal four regions with distinct oscillation frequencies, indicating a complex evolution of the Fermi surface [26,27]. At fields above 35 T a second phase (phase II) appears as a resistance minimum. At approximately 36 T a third phase (phase III) appears as a steplike increase in magnetoresistance, which extends up to ~ 39 T before giving way to a spin-polarized paramagnetic state (phase IV). Neutron-scattering experiments in pulsed magnetic fields recently showed that phase II is an incommensurate spin-density wave state with wave vector $k = (0.6, 0, 0)$ [20].

To compare with $x = 0$, low-temperature field sweeps of $R/R(300 \text{ K})$ up to $\mu_0H \lesssim 65 \text{ T}$ for $x \lesssim 0.48$ are shown in Fig. 2. An important feature that is seen in these data

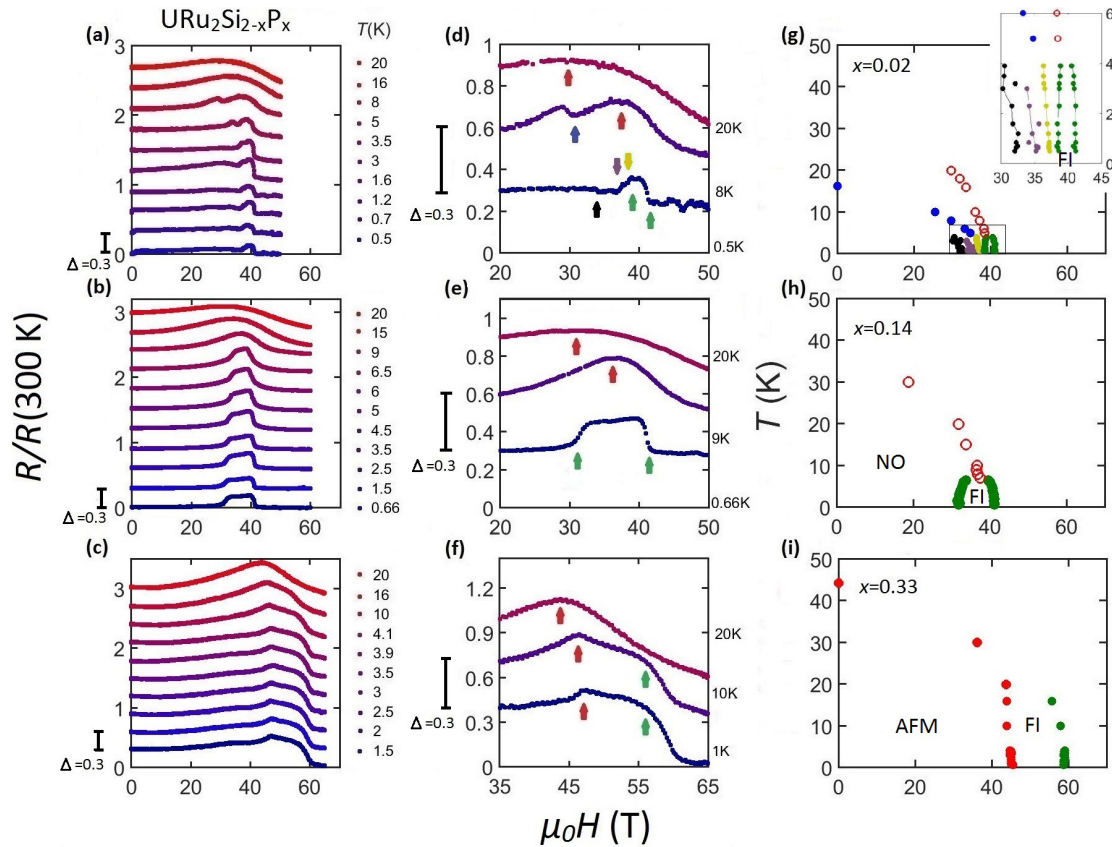


FIG. 3. Representative data from the three regimes: hidden-order x regime ($x = 0.02$), no-order x regime ($x = 0.14$), and the antiferromagnetic x regime ($x = 0.33$). Normalized resistance $R/R(300\text{ K})$ vs field $\mu_0 H$ plots are shown in panels (a)–(c), and the data are offset vertically by a constant amount Δ indicated in each panel. Panels (d)–(f) highlight the FI ordering from panels in (a)–(c). Colored arrows indicate phase transitions. Panels (g)–(i) show the $T - \mu_0 H$ phase diagrams, where phase boundaries are determined following the conventions from Ref. [19]. The inset region in panel (g) is the region of the cascade of phase transitions defined previously.

is that some form of high-field ordering persists for all x . Figure 3 details the high-field ordering and summarizes resulting $T - \mu_0 H$ phase diagrams for three concentrations spanning the $T - x$ phase diagram. Waterfall plots of all substitutions studied can be seen in the Supplemental Material [25]. For $0 < x \lesssim 0.035$, we first see a suppression of the magnetoresistance hump with increasing x , which may be due to increasing charge-carrier scattering due to chemical disorder. Over this x range there is a slight increase in the onset field of the FI phases and an enhancement of phase III, even as T_0 is suppressed [Figs. 3(a), 3(d), and 3(g)]. In the NO x regime ($0.035 \lesssim x \lesssim 0.26$) we unexpectedly observe a nearly square step FI feature between 30 and 45 T [Figs. 3(b), 3(e), and 3(h)] that appears to be a continuation of the low- x field-induced phases. This feature resembles that seen in URu_{1.92}Rh_{0.08}Si₂, which shows zero-field behavior similar to that seen in the NO x regime of our series [28–30]. Within the AFM ordered regime ($0.26 \lesssim x \lesssim 0.5$), an applied field suppresses the antiferromagnetism and produces a step in the magnetoresistance similar to that seen for the FI phase [Figs. 3(c), 3(f), and 3(i)]. We note that this phase is enhanced in both temperature and magnetic field range by comparison to the lower- x field-induced phase. These results might suggest that it is a distinct phase, and measurements such as neutron scattering are needed to clarify this question.

Figure 4 shows waterfall plots of magnetization M vs $\mu_0 H$ for concentrations in the different regions of the phase diagram: $x = 0$ (HO x regime with FI phase), $x = 0.1$ (NO x regime with FI phase), and $x = 0.33$ (AFM x regime with FI phase). The $x = 0$ data, taken from Ref. [28], reveal a linear in-field magnetization up to 35.8 T, where a jump

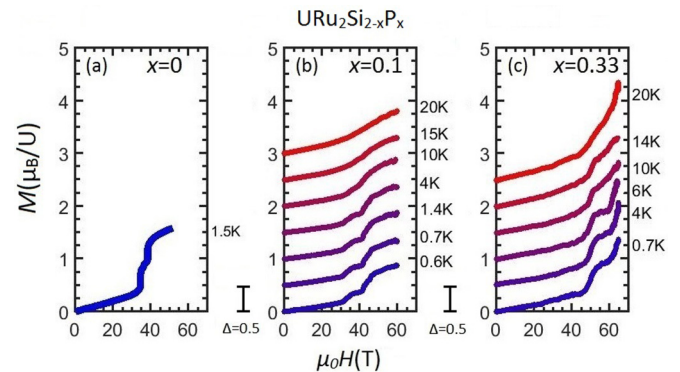


FIG. 4. Waterfall plots of the magnetization M vs magnetic field $\mu_0 H$ data. Data in panel (a) are taken from Ref. [28]. Panels (b) and (c) summarize data for two substitutions, representing materials in the no-ordered ($x = 0.1$) and antiferromagnetic ($x = 0.33$) x regimes, respectively. Data are offset by $\Delta = 0.5$ for clarity.

to 1/3 of the saturation value occurs, followed by a series of more subtle features before reaching a saturation value. Qualitatively similar behavior is seen for $x = 0.1$, where the data displays a 1/3 step feature at $\mu_0 H = 35$ T. The double step feature that is seen at 35.6 T in the parent compound is absent, and the second jump to the saturation moment occurs near 40 T. This single-plateau region matches with the FI phase seen in the magnetoresistance data at similar P concentration (see Fig. 2). The $x = 0.33$ data also reveal a step in M , characterized by a broadened transition width starting at around 47.5 T. After the plateau the magnetization rises again and does not reach a saturation value. The magnetization plateau occurs on the same field range as the FI phase, which appears at high field past the AFM phase in magnetoresistance measurements and is most likely due to magnetic ordering.

IV. DISCUSSION

From these measurements we construct the $T - x - \mu_0 H$ phase diagram for $\text{URu}_2\text{Si}_{2-x}\text{P}_x$, which features significant complexity with high-field ordering persisting across the entire substitution series (Fig. 1). In the HO x regime, the FI phases retain many of the characteristic features of the parent compound. This is even as the hidden-order temperature is suppressed, suggesting that the FI behavior is not solely tied to the field-driven collapse of HO. Within the NO x regime, the onset of the FI state expands to lower $\mu_0 H$ but terminates at high $\mu_0 H$ values, similar to what is seen in the HO x regime. Furthermore, the qualitative features of the magnetization field dependence are similar to what is seen in the HO x regime: there are plateau regions in the magnetization with similar step sizes for both $x = 0$ and 0.1. An attractive explanation for this is that related types of FI order emerge in the HO and NO x regimes when a non-symmetry-breaking electronic crossover originating from the hybridization between the f and conduction electron states is suppressed towards zero temperature, where additional small features are seen for $x < 0.03$ that relate to the suppression of HO. Here the crossover is represented by the magnetoresistance maximum ρ_{max} . Similar behavior is seen in several other strongly correlated metals without zero-field ordered ground states, including CeRu_2Si_2 , UPt_3 , and $\text{Sr}_2\text{Ru}_3\text{O}_7$ [24,31].

Measurements that target the order parameters are needed to distinguish between the ordered states in this phase diagram. Nonetheless, some insight is gained by considering the similarities between the s/p (Si \rightarrow P) and d -shell (Ru \rightarrow Rh) chemical substitution series [29,30]. In Fig. 5 we compare the in-field magnetization of 4%-Rh-substituted to 5%-P-substituted specimens, both of which are in the NO-order x region of their respective phase diagrams. For both examples, metamagnetic jumps appear at similar fields and consist of a 1/3 jump to the first plateau and then a second 2/3 jump to the full saturation value. Earlier work shows that the Rh-substituted material orders in a ferrimagnetic up-up-down state, which is seen in Fig. 5 as the first 1/3 magnetization jump [28]. We suggest that a similar type of ordering may occur for the $x = 0.1$ P-substituted example. Another intriguing feature is that the strength of the field-induced phase is enhanced by the suppression of AFM for $0.26 \lesssim x \lesssim 0.5$. The discontinuous evolution of the field and temperature extent

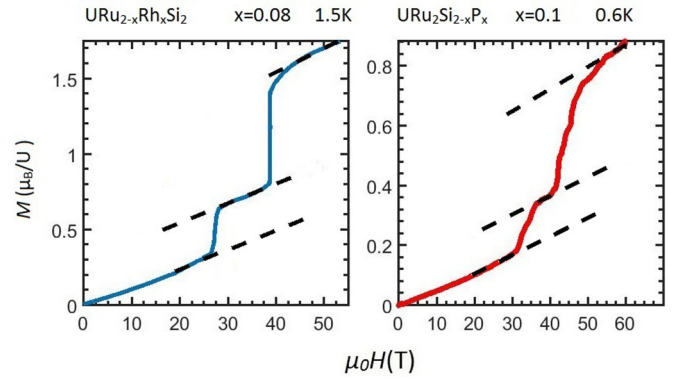


FIG. 5. Magnetization versus magnetic field $\mu_0 H$ for Rh-substituted [28] and P-substituted URu_2Si_2 at low temperatures in the no-order x regime. Dashed lines are guides to the eye to illustrate the 1/3 and 2/3 magnetization jumps seen in each material.

of the FI phase indicates that magnetic fluctuations resulting from the field-suppressed antiferromagnetism are involved in stabilizing this phase. Similar trends are also seen at high magnetic fields in the antiferromagnetic regions of the $T - P$ and Ru \rightarrow Fe phase diagrams, indicating a connection between these different parts of electronic phase space [32–37].

From these measurements, we conclude that hidden order is acutely unstable against simple electronic shell filling, regardless of whether it is done through chemical substitution on the d - or s/p electron sites. This is highlighted by considering that Si \rightarrow P and Ru \rightarrow Rh substitution might reasonably be expected to have distinct influences. For instance, they have different impact on (1) the spin-orbit coupling, (2) the lattice contraction and strain, (3) the local crystal electric field, and (4) in principle could alter different parts of the Fermi surface. Furthermore, given the complexity that is seen in other chemical substitution series, under applied pressure and in high magnetic fields [9–20], *a priori* it seems unlikely that Si \rightarrow P and Ru \rightarrow Rh substitution would be equivalent. Despite this, we find phase diagrams with similar features along both the electronic shell filling and applied magnetic field tuning axes. Given that this happens on the few-percent chemical substitution level and that the foundational Kondo lattice is unchanged by such small changes, it may now be possible to systematically uncover which factors underpin hidden order and thereby constrain possible theoretical models. Future measurements to probe the electronic state using advanced techniques such as angle-resolved photoelectron spectroscopy and electronic Raman spectroscopy will be useful to do this.

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- [1] T. T. M. Palstra, A. A. Menovsky, J. van den Berg, A. J. Dirkmaat, P. H. Kes, G. J. Nieuwenhuys, and J. A. Mydosh, *Phys. Rev. Lett.* **55**, 2727 (1985).
- [2] W. Schlabit, J. Baumann, B. Pollit, U. Rauchschalbe, H. M. Mayer, U. Ahlheim, and C. D. Bredl, *Z. Phys. B* **62**, 171 (1986).
- [3] M. B. Maple, J. W. Chen, Y. Dalichaouch, T. Kohara, C. Rossel, M. S. Torikachvili, M. W. McElfresh, and J. D. Thompson, *Phys. Rev. Lett.* **56**, 185 (1986).
- [4] J. A. Mydosh and P. M. Oppeneer, *Rev. Mod. Phys.* **83**, 1301 (2011).
- [5] J. A. Mydosh and P. M. Oppeneer, *Philos. Mag.* **94**, 3642 (2014).
- [6] T. T. M. Palstra, A. A. Menovsky, G. J. Nieuwenhuys, and J. A. Mydosh, *J. Magn. Magn. Mater.* **54**, 435 (1986).
- [7] T. Endstra, G. J. Nieuwenhuys, and J. A. Mydosh, *Phys. Rev. B* **48**, 9595 (1993).
- [8] P. Das, R. E. Baumbach, K. Huang, M. B. Maple, Y. Zhao, J. S. Helton, J. W. Lynn, E. D. Bauer, and M. Janoschek, *New J. Phys.* **15**, 053031 (2013).
- [9] M. W. McElfresh, J. D. Thompson, J. O. Willis, M. B. Maple, T. Kohara, and M. S. Torikachvili, *Phys. Rev. B* **35**, 43 (1987).
- [10] N. Kanchanavatee, M. Janoschek, R. E. Baumbach, J. J. Hamlin, D. A. Zocco, K. Huang, and M. B. Maple, *Phys. Rev. B* **84**, 245122 (2011).
- [11] N. Kanchanavatee, B. D. White, V. W. Burnett, and M. B. Maple, *Philos. Mag.* **94**, 3681 (2014).
- [12] P. Das, N. Kanchanavatee, J. S. Helton, K. Huang, R. E. Baumbach, E. D. Bauer, B. D. White, V. W. Burnett, M. B. Maple, J. W. Lynn, and M. Janoschek, *Phys. Rev. B* **91**, 085122 (2015).
- [13] Y. Dalichaouch, M. B. Maple, J. W. Chen, T. Kohara, C. Rossel, M. S. Torikachvili, and A. L. Giorgi, *Phys. Rev. B* **41**, 1829 (1990).
- [14] N. P. Butch and M. B. Maple, *J. Phys.: Condens. Matter* **22**, 164204 (2010).
- [15] Y. Dalichaouch, M. B. Maple, M. S. Torikachvili, and A. L. Giorgi, *Phys. Rev. B* **39**, 2423 (1989).
- [16] M. Jaime, K. H. Kim, G. Jorge, S. McCall, and J. A. Mydosh, *Phys. Rev. Lett.* **89**, 287201 (2002).
- [17] J. S. Kim, D. Hall, P. Kumar, and G. R. Stewart, *Phys. Rev. B* **67**, 014404 (2003).
- [18] K. H. Kim, N. Harrison, M. Jaime, G. S. Boebinger, and J. A. Mydosh, *Phys. Rev. Lett.* **91**, 256401 (2003).
- [19] G. W. Scheerer, W. Knafo, D. Aoki, G. Ballon, A. Mari, D. Vignolles, and J. Flouquet, *Phys. Rev. B* **85**, 094402 (2012).
- [20] W. Knafo, F. Duc, F. Bourdarot, K. Kuwahara, H. Nojiri, D. Aoki, J. Billette, P. Frings, X. Tonon, E. Lelièvre-Berna, J. Flouquet, and L.-P. Regnault, *Nat. Commun.* **7**, 13075 (2016).
- [21] A. Gallagher, K.-W. Chen, C. M. Moir, S. K. Cary, F. Kametani, N. Kikugawa, D. Graf, T. E. Albrecht-Schmitt, S. C. Riggs, A. Shekhter, and R. E. Baumbach, *Nat. Commun.* **7**, 10712 (2016).
- [22] A. Gallagher, K.-W. Chen, S. K. Cary, F. Kametani, D. Graf, T. E. Albrecht-Schmitt, A. Shekhter, and R. E. Baumbach, *J. Phys.: Condens. Matter* **29**, 024004 (2017).
- [23] K. R. Shirer, M. Lawson, T. Kissikov, B. T. Bush, A. Gallagher, K.-W. Chen, R. E. Baumbach, and N. J. Curro, *Phys. Rev. B* **95**, 041107(R) (2017).
- [24] S. A. Grigera, R. S. Perry, A. J. Schofield, M. Chiao, S. R. Julian, G. G. Lonzarich, S. I. Ikeda, Y. Maeno, A. J. Mills, and A. P. Mackenzie, *Science* **294**, 329 (2001).
- [25] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevB.96.085141> for data across full range of temperature and magnetic field, phase diagrams, and a plot of the c/a ratio and the RRR of the material.
- [26] M. M. Altarawneh, N. Harrison, S. E. Sebastian, L. Balicas, P. H. Tobash, J. D. Thompson, F. Ronning, and E. D. Bauer, *Phys. Rev. Lett.* **106**, 146403 (2011).
- [27] E. Hassinger, G. Knebel, T. D. Matsuda, D. Aoki, V. Taufour, and J. Flouquet, *Phys. Rev. Lett.* **105**, 216409 (2010).
- [28] K. Kuwahara, S. Yoshii, H. Nojiri, D. Aoki, W. Knafo, F. Duc, X. Fabrèges, G. W. Scheerer, P. Frings, G. L. J. A. Rikken, F. Bourdarot, L. P. Regnault, and J. Flouquet, *Phys. Rev. Lett.* **110**, 216406 (2013).
- [29] K. H. Kim, N. Harrison, H. Amitsuka, G. A. Jorge, M. Jaime, and J. A. Mydosh, *Phys. Rev. Lett.* **93**, 206402 (2004).
- [30] K. H. Kim, N. Harrison, H. Amitsuka, G. A. Jorge, M. Jaime, and J. A. Mydosh, [arXiv:cond-mat/0411068v1](https://arxiv.org/abs/cond-mat/0411068v1).
- [31] F. Honda, T. Takeuchi, S. Yasui, Y. Taga, S. Yoshiuchi, Y. Hirose, Y. Tomooka, K. Sugiyama, M. Hagiwara, K. Kindo, R. Settai, and Y. Ōnuki, *J. Phys. Soc. Jpn.* **82**, 084705 (2013).
- [32] D. Aoki, F. Bourdarot, E. Hassinger, G. Knebel, A. Miyake, S. Raymond, V. Taufour, and J. Flouquet, *J. Phys. Soc. Jpn.* **78**, 053701 (2009).
- [33] S. Ran, I. Jeon, N. Kanchanavatee, K. Huang, M. B. Maple, A. Gallagher, K. Chen, D. Graf, R. Baumbach, and J. Singelton Bulletin Am. Phys. Soc. **62** (2017).
- [34] K. T. Moore and G. van Laan, *Rev. Mod. Phys.* **81**, 235 (2009).
- [35] R. E. Baumbach, Z. Fisk, F. Ronning, R. Movshovich, J. D. Thompson, and E. D. Bauer, *Philos. Mag.* **94**, 3663 (2014).
- [36] E. Hassinger, G. Knebel, K. Izawa, P. Lejay, B. Salce, and J. Flouquet, *Phys. Rev. B* **77**, 115117 (2008).
- [37] S. Ran, C. T. Wolowiec, I. Jeon, N. Pouse, N. Kanchanavatee, K. Huang, D. Martien, T. DaPron, D. Snow, M. Williamsen, S. Spagna, and M. B. Maple, *Proc. Natl. Acad. Sci. USA* **113**, 13348 (2016).