

Proposal for measuring the finite-temperature Drude weight of integrable systems

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Integrable models such as the spin-1/2 Heisenberg chain, the Lieb-Liniger, or the one-dimensional Hubbard model are known to avoid thermalization, which was also demonstrated in several quantum-quench experiments. Another dramatic consequence of integrability is the zero-frequency anomaly in transport coefficients, which results in ballistic finite-temperature transport, despite the presence of strong interactions. While this aspect of nonergodic dynamics has been known for a long time, there has so far not been any unambiguous experimental realization thereof. We make a concrete proposal for the observation of ballistic transport via local quantum-quench experiments in fermionic quantum-gas microscopes. Such an experiment would also unveil the coexistence of ballistic and diffusive transport channels in one and the same system and provide a means of measuring finite-temperature Drude weights. The connection between local quenches and linear-response functions is established via time-dependent Einstein relations.

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Introduction. Nonergodic dynamics in closed many-body quantum systems is one of the most actively investigated branches of nonequilibrium physics [1–4]. The canonical examples are either Bethe-ansatz integrable one-dimensional (1D) models such as the spin-1/2 XXZ chain, the Fermi-Hubbard model [5], hard-core bosons [6], or many-body localized systems [7–9]. Integrable systems possess an extensive set of local conserved quantities that can constrain the long-time behavior in the relaxation dynamics starting from nonequilibrium initial conditions [10] induced by, e.g., quantum quenches. This leads to the failure of these systems to thermalize with respect to standard thermodynamic ensembles (for a review, see [11]), rooted in the violation of the eigenstate thermalization hypothesis [12–15].

Another prominent consequence of integrability in clean systems is the possibility of anomalous transport properties at finite temperatures and in the linear-response regime as was shown in a seminal paper by Zotos *et al.* [16]. Within the Kubo formalism, one decomposes conductivities into a regular part and a zero-frequency contribution with the Drude weight D ,

$$\text{Re } \sigma(\omega) = 2\pi D(T)\delta(\omega) + \sigma_{\text{reg}}(\omega). \quad (1)$$

The presence of a nonzero Drude weight $D(T) > 0$ signals ballistic transport, which, at finite temperatures T , is an unusual behavior in a quantum many-body system without momentum conservation. The Drude weight can be related to the long-time value of current autocorrelation functions $C(t) = \langle I(t)I(0) \rangle_{\text{eq}}/L$ (I being the extensive current operator and $\langle A \rangle_{\text{eq}} = Z^{-1} \text{tr}[Ae^{-H/T}]$) [16] and therefore, $D > 0$ directly implies nonergodic behavior in that correlation function. The connection to integrability is usually drawn via the Mazur inequality

$$D \geq \frac{1}{2T^{\nu}L} \sum_{\alpha} \frac{|\langle I Q_{\alpha} \rangle_{\text{eq}}|^2}{\langle Q_{\alpha}^2 \rangle_{\text{eq}}}, \quad (2)$$

where Q_{α} are the local and pseudolocal [17] conserved charges for the Hamiltonian H under consideration, i.e., $[H, Q_{\alpha}] = 0$, $\langle Q_{\alpha}^2 \rangle_{\text{eq}} \propto L$, $\langle I Q_{\alpha} \rangle_{\text{eq}} \propto L$, chosen such that $\langle Q_{\alpha} Q_{\alpha'} \rangle_{\text{eq}} = 0$ for $\alpha \neq \alpha'$, and $\nu = 1, 2$ for charge (spin) and energy transport, respectively.

The most famous example is energy transport in the spin-1/2 XXZ chain, whose Hamiltonian is given by $H = \sum_l h_l$ with local terms

$$h_l = J \left[\frac{1}{2} (S_l^+ S_{l+1}^- + \text{H.c.}) + \Delta S_l^z S_{l+1}^z \right], \quad (3)$$

where S_l^{μ} , $\mu = x, y, z$ are the components of a spin-1/2 operator acting on site l with $S_l^{\pm} = S_l^x \pm i S_l^y$. J is the exchange coupling and Δ denotes an exchange anisotropy. In this model, the energy-current operator I_E itself is conserved $[H, I_E] = 0$ [16,18,19], rendering spin-1/2 XXZ chains ballistic thermal conductors.

Spin transport in the same model is also ballistic in its gapless phase $|\Delta| < 1$ [20–27]. The most debated case has been zero magnetization, since until the work [28] (see also [17,29–32]) the relevant conservation laws were not known. For $\Delta > 1$, the common belief is that there is diffusion [33–36] such that there is a coexistence of ballistic thermal and diffusive spin transport. Notably, there is no final conclusion yet on the qualitative nature of spin transport at exactly $\Delta = 1$, i.e., the Heisenberg chain [22,27,37–43], with superdiffusive dynamics being one possible scenario.

A similar situation applies to the Fermi-Hubbard chain defined by

$$h_l = -t_0 \sum_{\sigma} (c_{l\sigma}^{\dagger} c_{l+1\sigma} + \text{H.c.}) + \frac{U}{2} \left[\left(n_{l\uparrow} - \frac{1}{2} \right) \left(n_{l\downarrow} - \frac{1}{2} \right) + \left(n_{l+1\uparrow} - \frac{1}{2} \right) \left(n_{l+1\downarrow} - \frac{1}{2} \right) \right], \quad (4)$$

where $c_{l\sigma}$ annihilates a fermion with spin $\sigma = \uparrow, \downarrow$ on site l , and $n_{l\sigma} = c_{l\sigma}^{\dagger} c_{l\sigma}$, with U the interaction strength and t_0 the

hopping matrix element. The Hubbard model is a ballistic thermal conductor at any $T > 0$ [16,44]. At half filling, most studies indicate diffusive charge transport [45–48], while away from half filling, charge transport is ballistic. Thus, in the Hubbard chain there is a similar coexistence between a diffusive (charge) and a ballistic (energy) transport channel [44].

While the prediction of anomalous transport in integrable models is very intriguing, its direct relevance to solid-state experiments is unclear since there, external scattering channels such as phonons or impurities will usually dominate the behavior [49–54]. Nevertheless, there has been an impressive series of experiments focusing mostly on thermal transport in quantum magnets [55–68], reporting remarkably large thermal conductivities in 1D systems [55,69]. The cleanest evidence for ballistic transport in an integrable model has so far been observed in a strongly interacting 1D Bose gas expanding in an optical lattice [70]: the density profiles of these hard-core bosons are indistinguishable from noninteracting particles and the cloud thus expands ballistically. An experimental demonstration of the aforementioned coexistence of a diffusive and a ballistic transport channel at finite temperatures thus remains open, as much as a quantitative measurement of the Drude weights, which so far were mostly considered rather academic quantities.

In our work, we propose a realistic setup to measure finite-temperature Drude weights using optical-lattice experiments with single-site resolution and addressing capabilities [71–73]. The idea relies on preparing small local perturbations in spin- or charge densities and monitoring their spreading as a function of time. The time dependence of their width yields information on whether transport is diffusive, ballistic, or something intermediate and access to diffusion constants *and* Drude weights can be gained from (generalized) Einstein relations. Einstein relations are usually quoted for the diffusion constant $\mathcal{D} = \sigma_{\text{dc}}/\chi$ relating those to the dc conductivity σ_{dc} and to the corresponding susceptibility χ . The extension of such relations to ballistic transport *and* the transient regime before the asymptotic dynamics is reached is well documented in the literature [33,74,75] (albeit perhaps little appreciated beyond the diffusive case).

In this Rapid Communication we demonstrate the validity of time-dependent generalized Einstein relations by a direct comparison of the time dependence of spatial variances in local quantum quenches to the expectation from time-dependent correlation functions for various types of transport in the spin-1/2 XXZ chain and the 1D Hubbard model, using finite-temperature density matrix renormalization group simulations with the purification approach [25,76–78]. We show that the time scales necessary to resolve ballistic dynamics are short enough to be observable in typical optical-lattice experiments. Moreover, the qualitative difference between the ballistic energy transport in both spin and Hubbard chains versus the diffusive spin or charge transport (in the respective parameter regimes) can be unveiled on time scales $\lesssim 10/J$ or $\lesssim 5/t_0$.

The 1D Hubbard model provides the most promising avenue: several fermionic quantum-gas microscopes are now operational [80–88] and the 1D regime has already been accessed [84]. Unlike the ongoing quest to observe long-range antiferromagnetic correlations in this model [84,88–92],

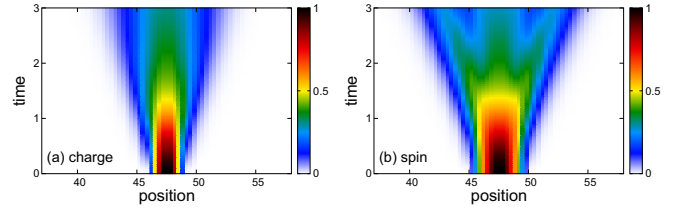


FIG. 1. Example of numerical results for the spreading of (a) a charge and (b) a spin wave packet induced by a local quench in the half-filled 1D Fermi-Hubbard model with $U/t_0 = 8$ and a spin imbalance of $\Delta n = (N_{\uparrow} - N_{\downarrow})/L = 0.267$ at a temperature $T/t_0 = 20$, where N_{σ} is the number of fermions with spin $\sigma = \uparrow, \downarrow$. In (a) and (b), we show the charge and spin density differences $\langle n_{l\uparrow} \rangle + \langle n_{l\downarrow} \rangle - 1$ and $\langle n_{l\uparrow} \rangle - \langle n_{l\downarrow} \rangle - \Delta n$ (both normalized to 1), respectively. In the initial state, we prepare two empty sites in the center of the system. Our analysis (see the discussion of Fig. 4) suggests that the charge dynamics is not ballistic, consistent with the prediction of diffusion [36,45–48,79], while spin dynamics is ballistic.

cooling to low temperatures is not necessary since the ballistic transport properties are protected by conservation laws and are thus observable at any $T > 0$. We propose to study the spin-imbalanced regime of a 1D Hubbard model at half filling, where charge transport is diffusive, but spin transport is ballistic. This effect is also of fundamental interest, being an example of a coexistence of ballistic and diffusive finite temperature transport not involving thermal transport.

Generalized time-dependent Einstein relations. The spatial variance of a density wave packet can be related to linear-response functions via [33,74,75]

$$\delta\sigma_v^2(t) = \frac{2}{L\chi_v} \int_0^t dt_1 \int_0^{t_1} dt_2 \langle I_v(t_1)I_v(t_2) \rangle_{\text{eq}}, \quad (5)$$

where $\delta\sigma_v^2(t) = \sigma_v^2(t) - \sigma_v^2(0)$, $v = \text{th}, s, c$ can be energy, spin, or charge (for definitions of I_v and of σ_v , see [93]). At infinite temperature, the susceptibility χ_v is $\chi_s = 1/4$, $\chi_{\text{th}}/J^2 = 1/8 + \Delta^2/16$ for the XXZ chain and $\chi_c = 1/2$, $\chi_{\text{th}}/t_0^2 = 1 + U^2/16$ for the Hubbard model, but is also known exactly at any finite temperature [5,94]. A typical example for an initial state used in our simulations to induce the dynamics and the real space and time dependence of such a density perturbation is shown in Fig. 1.

In the ballistic case, $\langle I_v(t)I_v \rangle_{\text{eq}}/L \rightarrow 2T^\gamma D_v$ (where $\gamma = 2$ in the thermal case and $\gamma = 1$ otherwise), and therefore, from Eq. (5), $\delta\sigma_v^2(t) \rightarrow \frac{T^\gamma D_v}{\chi_v} t^2$. For diffusive dynamics, one obtains $\delta\sigma_v^2(t) \rightarrow 2\mathcal{D}_v t$ with $\mathcal{D}_v = \sigma_{\text{dc},v}/\chi_v$.

Spin-1/2 XXZ chain. We begin by discussing results for the spin-1/2 XXZ chain with $\Delta = 0.5$ and $\Delta = 1.5$, i.e., in its massless and massive regimes. Unless stated otherwise, we consider infinite temperature $T = \infty$ (see also [36,95–97]).

Here, we provide an example of the equivalence of the left-hand side (lhs) and right-hand side (rhs) of Eq. (5) at any time $t > 0$ for local quenches with a sufficiently small amplitude (see [93] for an analysis of the influence of the amplitude). Figure 2(a) shows the results for a perturbation in the spin density that is initially prepared on $L_m = 2$ sites (see [93] for definitions). For $\Delta = 0.5$, there is a ballistic spreading with $\delta\sigma_s^2 \propto t^2$, while for $\Delta = 1.5$ and at long times,

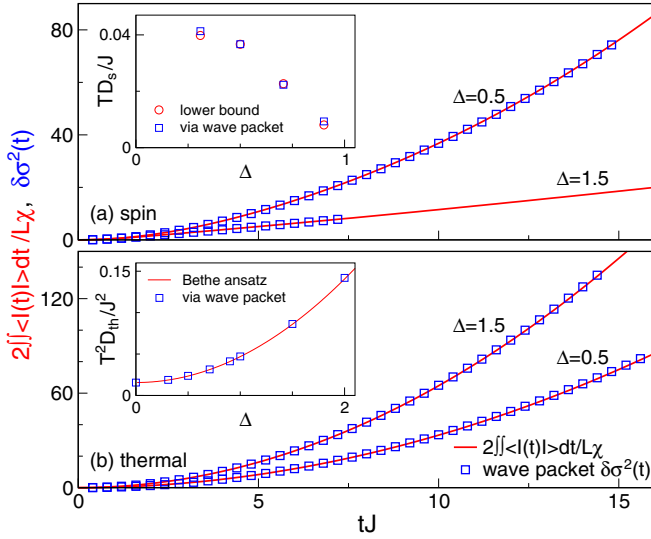


FIG. 2. Time evolution of the square variance $\sigma_v^2(t) - \sigma_v^2(0)$ of a wave packet as well as of the right-hand side of Eq. (5) for the spin-1/2 XXZ spin chain at $T = \infty$ in (a) a spin quench ($L_m = 2$) and (b) an energy quench ($T_1/J \geq 10$, $T_2 = \infty$, $L_m = 2$). Insets: Comparison of the Drude weights extracted from the wave packets with the lower bound [29] [spin case in (a)] and with the Bethe-ansatz result [18,19] for D_{th} in (b). The Drude weights D_s and D_{th} were extracted from fits to $a + bt + ct^2$ and ct^2 over the last third of the curve, respectively.

the variance approaches a diffusive behavior with $\delta\sigma_s^2 \propto t$. The agreement between $\delta\sigma_s^2(t)$ (solid lines) and the Kubo formula [double integral, rhs of Eq. (5)] is excellent at all times. While the diffusion constant was extracted already from such data for $\Delta > 1$ [36], we here present results for the spin Drude weight for $\Delta < 1$. The comparison of wave-packet dynamics versus Kubo (double integral of current autocorrelations) show that either approach gives the same quantitative results. For the comparison with the literature we will use the method that gives access to the longest times and hence more accurate estimates of the Drude weight. In the case of spin quenches, this is the time propagation of current correlations. The results for D_s computed for various values of $0 < \Delta < 1$ are shown in the inset of Fig. 2(a) and are in very good agreement with the lower bound for D_s from [29], which is believed to be exhaustive for the chosen values of Δ .

Figure 2(b) shows the corresponding data for the thermal Drude weight, where the dynamics is induced by an energy quench, in which we embed a region of higher temperature T_2 into a background with $T_1 < T_2$ (here, $T_1/J = 10$ at $\Delta = 0.5$, $T_1/J = 20$ at $\Delta = 1.5$, $T_2 = \infty$, $L_m = 2$ [93]). D_{th} is extracted from purely quadratic fits to $\delta\sigma_{th}(t)$. The agreement with the exact Bethe-ansatz results [18,19] is excellent, as shown in the inset of Fig. 2(b).

1D Hubbard model: Half filling, zero magnetization. Next we turn to the Fermi-Hubbard model with filling $n = 1$ and a vanishing magnetization ($N_\uparrow = N_\downarrow$). Figures 3(a) and 3(b) show the comparison of the lhs and rhs of Eq. (5) for a perturbation in the charge and the energy density, respectively, which agree well in both cases. The difference in the qualitative behavior of charge and energy transport is best unveiled by comparing the exponents of the variance at the longest times

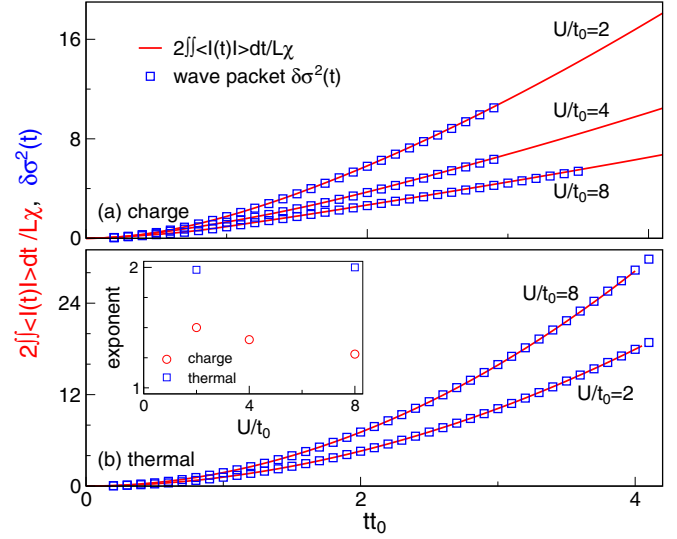


FIG. 3. The same as in Fig. 2, but for the Fermi-Hubbard model at $T = \infty$ and considering (a) a charge quench ($L_m = 2$) and (b) an energy quench ($T_1/t_0 = 20$, $T_2 = \infty$, $L_m = 2$). Inset: Exponents determined from a power-law fit to the last third of the curve.

reached. The inset in Fig. 3(b) suggests that perturbations in the energy density spread ballistically as expected from theory [16,44] with an exponent of 2, while for the charge dynamics, the exponents are consistently smaller than 2 and the smaller the exponent, the larger U/t_0 is. We expect that the asymptotic diffusive behavior [47,48] will emerge at times longer than what is accessible in our simulations. Note that one can also study the time dependence of energy and charge density perturbations in quenches that touch both densities [44]: in that case, one can visualize the coexistence of ballistic energy transport with diffusive charge transport in the same time evolution.

In principle, such an initial perturbation of both energy and charge density may thus yield the desired recipe for an experimental observation. While measuring double occupancy $d_i = \langle n_{i\uparrow}n_{i\downarrow} \rangle$, which yields the interaction energy, is standard (see, e.g., [84]), there is, unfortunately, as of now no good way of extracting energy densities in fermionic quantum-gas microscopes. We therefore turn to our most promising example, namely, spin and charge dynamics at half filling but away from zero magnetization.

1D Hubbard model: Half filling with spin imbalance. We introduce a spin imbalance $\Delta n = (N_\uparrow - N_\downarrow)/L$ while keeping the half-filling condition. In this case, the spin transport becomes ballistic (different from the case $\Delta n = 0$) since the spin current acquires a finite overlap with local conserved quantities [16], constructed in [98]. Charge transport, however, is expected to remain diffusive, since due to particle-hole symmetry and the half-filling condition the charge current operator remains orthogonal to local conserved operators [98]. It would in principle be possible that other quasilocal conserved operators existed in the Hubbard model which were not orthogonal to charge current in this case, similar as for the spin current in the massless regime of XXZ chain, but numerical computations [99] using the method of Ref. [30] indicate that this is not the case.

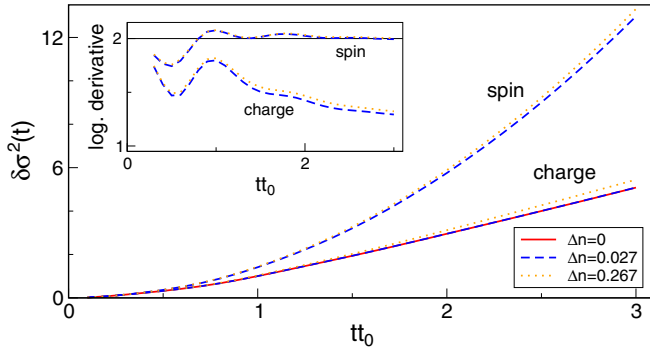


FIG. 4. Comparison of σ_c^2 and σ_s^2 for a charge quench ($L_m = 2$) as a function of the imbalance at $U/t_0 = 8$ and $T/t_0 = 20$. Inset: Logarithmic derivative.

In order to illustrate the coexistence of ballistic spin and diffusive charge dynamics at $T > 0$, we present results for a charge quench with $T/t_0 = 20$ and $L_m = 2$ empty sites, which for $\Delta n \neq 0$ also affects the spin density. Figure 4 shows the time dependence of the variances for different values of Δn , while we present the exponents in the inset. First, we observe that spin propagates much faster than charge at $T = \infty$ and moreover, $\delta\sigma_s^2 \propto t^2$ as soon as $\Delta n > 0$, clearly indicating ballistic spin transport. Second, the exponents for the charge dynamics are very close to unity for small Δn and increase slightly as Δn becomes large. This behavior suggests diffusive charge dynamics, keeping in mind that the simulations may not yet have reached the asymptotic regime. Note that this difference in the charge versus spin dynamics is not spin-charge separation, which is a low-energy phenomenon with no qualitative difference in the propagation of spin and charge.

We stress that the validity of Einstein relations (5) depends on the nature of the initial local perturbation. Namely, the wave packet of Fig. 4 is formed by an experimentally realizable quench of producing $L_m = 2$ empty sites $|00\rangle$ in the spin-imbalanced background. Unlike other initial states, it shows a curious discontinuity of the spin variance $\delta\sigma_s^2(t)$ at $T = \infty$, as for $\Delta n = 0$ there is no spin transport anymore and $\sigma_s^2(t) \equiv 0$, while at any finite $\Delta n > 0$, spin propagates ballistically with a Δn -independent velocity. However, if one considers genuine local spin quenches, e.g., preparing the central two sites in the state $|0\rangle|\uparrow\rangle + |\uparrow\rangle|0\rangle$, then $\sigma_s^2(t)$ nicely follows the linear-response prediction (5) and smoothly switches from ballistic to diffusive growth as $\Delta n \rightarrow 0$ (results not shown). (For another recent study of typical vs atypical initial states, see [97].) We stress that in order to resolve the difference between ballistic spin and diffusive charge dynamics, it is not necessary to prepare local quenches with a small amplitude.

Experimental aspects. As we argued in the Introduction, the half-filled Hubbard chain away from zero magnetization

provides the most promising route to experimentally observe the coexistence of a ballistic (here spin) and a diffusive transport channel (here charge) in one and the same model. A clear advantage of working directly with the Fermi-Hubbard chain is that one does not have to worry about defects. For instance, in simulations of the Heisenberg model using a strongly interacting two-component Bose gas [100], holes or doublons invalidate the mapping to a pure spin model. Similarly, interactions have to be large for the mapping to be valid, while the interaction strength in the Fermi-Hubbard chain plays no role for the qualitative aspects discussed here, yet may determine the onset of asymptotic diffusive behavior.

Reaching low temperatures is not important as such to see the integrability-protected ballistic transport in either the spin-1/2 XXZ or in the Hubbard chain. However, most optical-lattice experiments have a harmonic trapping potential [101]. This breaks integrability as such, yet since we are working at unit filling, a large region of the cloud will be in the Mott-insulating regime, with a constant density. This requires that temperatures should ideally be low enough to stabilize a sizable Mott plateau.

Summary. Experimentally measuring the finite-temperature Drude weight of an integrable 1D model has not been accomplished yet. We here proposed to use local quenches, i.e., perturbations in spin, charge, or energy density to induce a density wave-packet dynamics. We demonstrated that in ballistic regimes, the Drude weight can be extracted by simply monitoring the time dependence of the spatial variance and measuring the prefactor. This relies on generalized time-dependent Einstein relations [74], whose validity we verified in several examples including the spin-1/2 XXZ chain and Hubbard chains. We identified the charge and spin dynamics in the half-filled Hubbard chain with a spin imbalance as the most promising candidate for an experiment, given that many ingredients are available in several groups, such as fermionic quantum microscopes [80–88], 1D Fermi-Hubbard systems [84], and the ability to spin-selectively [84,88] monitor the time-dependent spreading of local perturbations [100,102,103]. Such an experiment would visualize the coexistence of ballistic and diffusive transport channels in the same quantum model [16,36,44,96], an extreme violation of the standard behavior of Fermi liquids or Luttinger liquids.

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